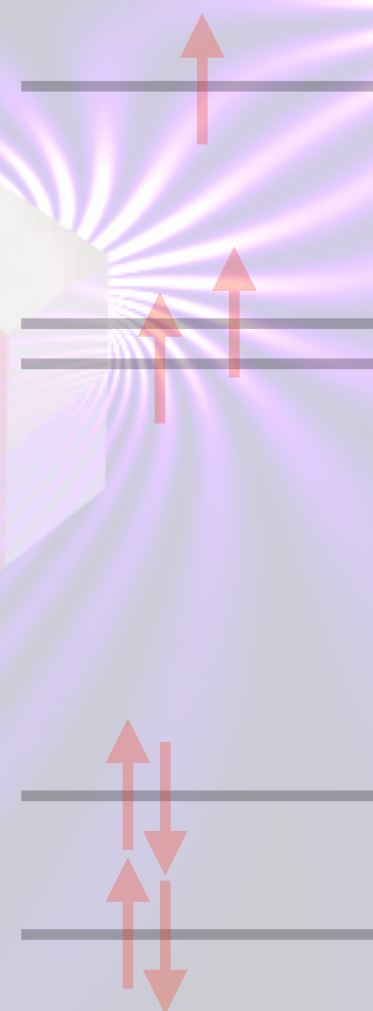
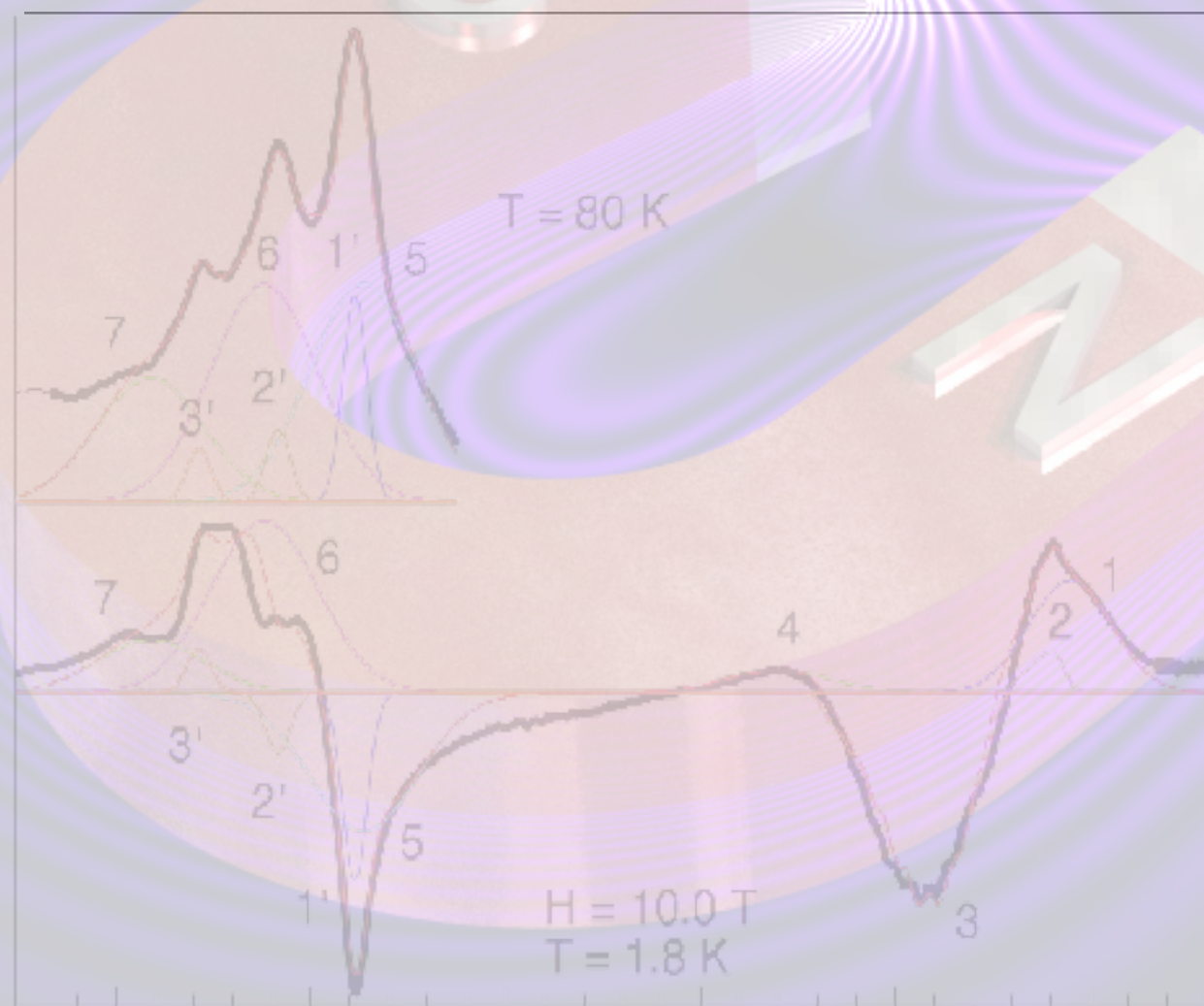


The Bioinorganic Workshop/Summer Symposium
in Molecular Biology 2018

Introduction to Coordination Chemistry and Ligand Field Theory



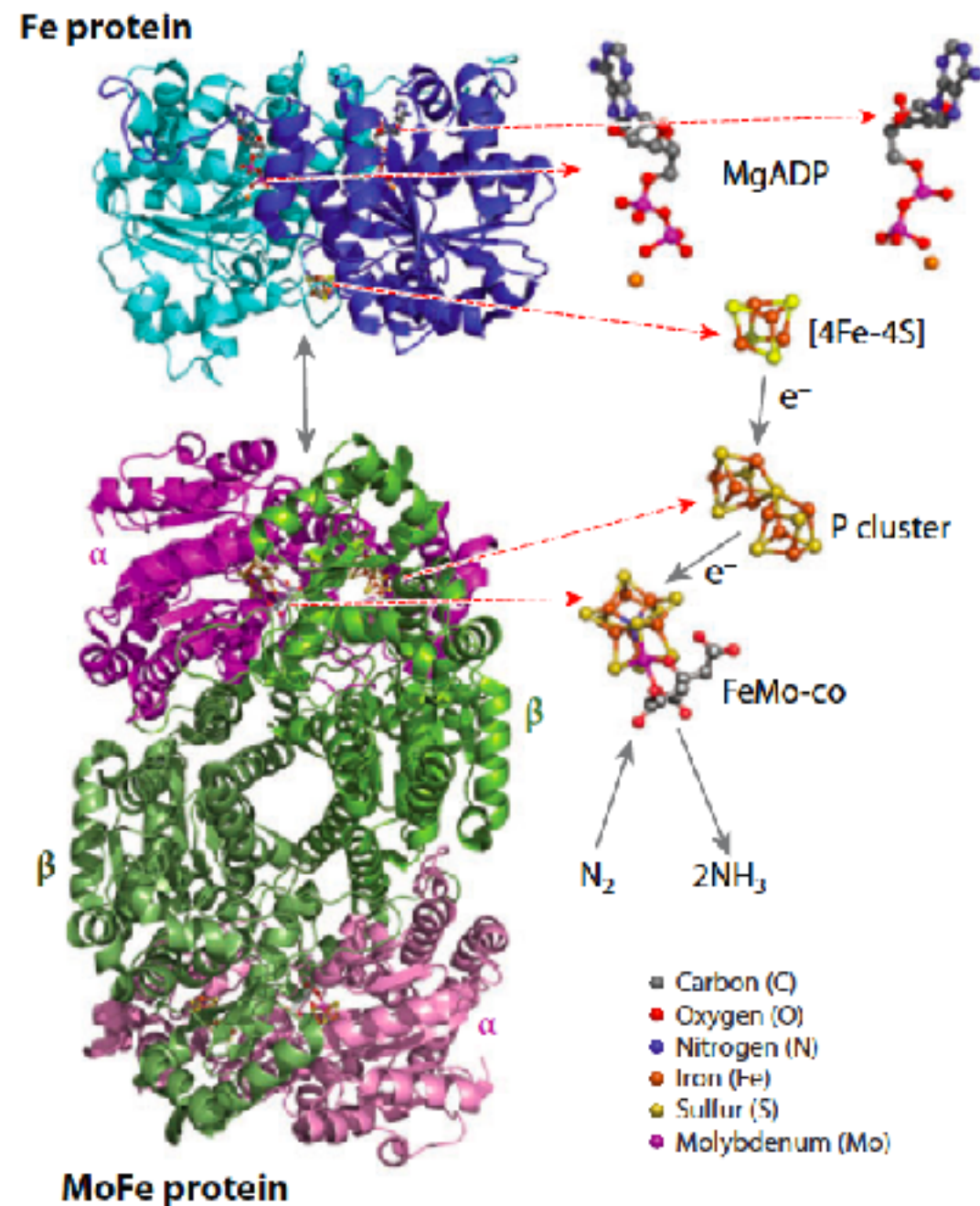
Frank Neese



MAX-PLANCK-GESELLSCHAFT

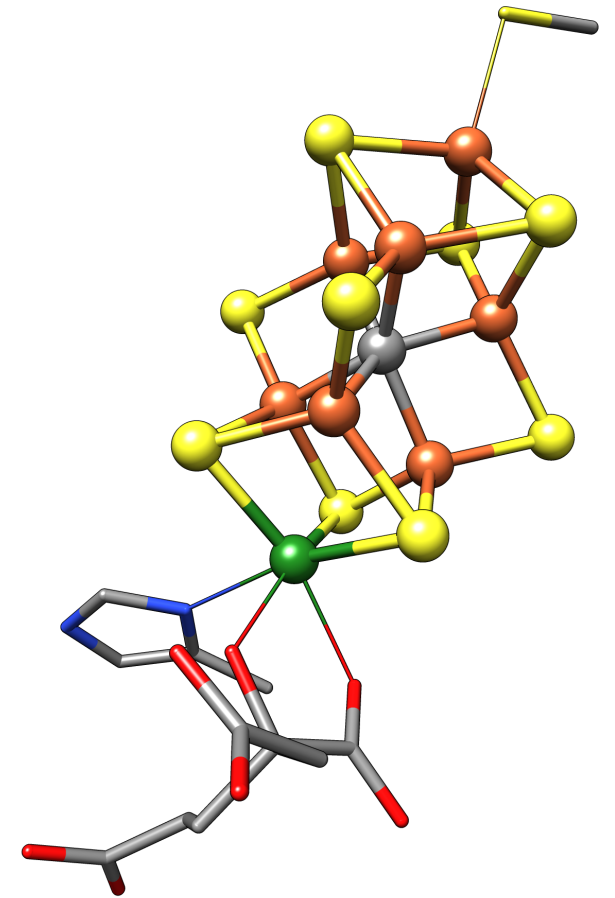
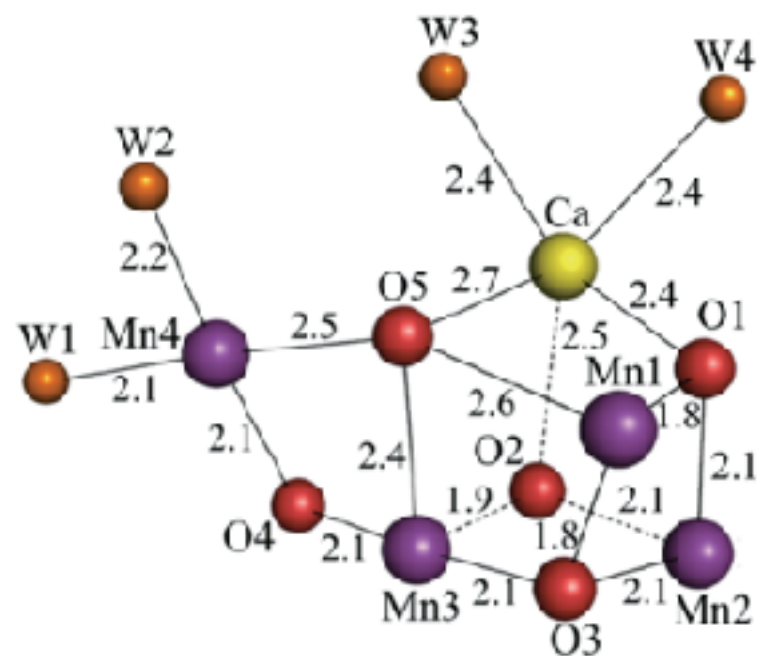
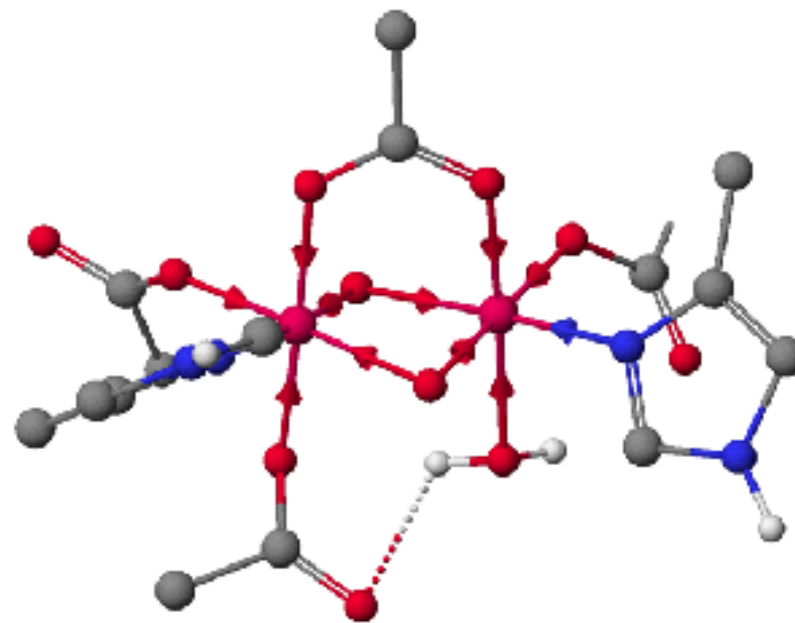
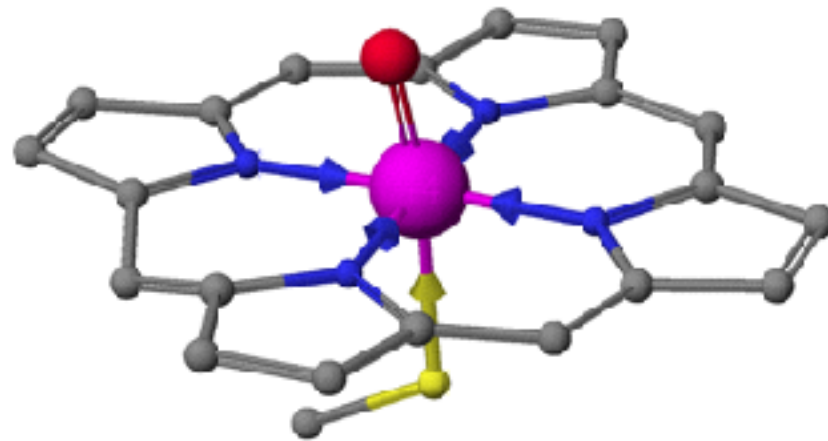
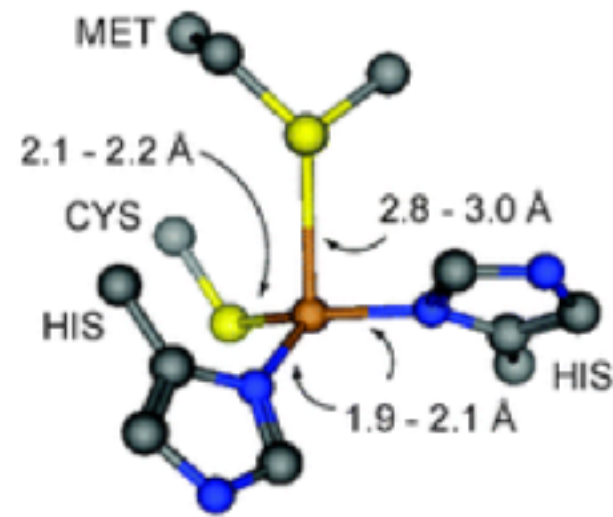
MPI für Kohlenforschung
Kaiser-Wilhelm Platz 1
Mülheim an der Ruhr

Transition Metal Sites in Enzymes



- ▶ ~1/3 of all enzymes contain metal cofactors
- ▶ These cofactors are the site of redox processes, substrate binding, reactivity
- ▶ A crystal structure does not tell the whole story...
- ▶ Spectroscopy provides a route to understanding changes that occur at the active sites during electron transfer and catalytic processes

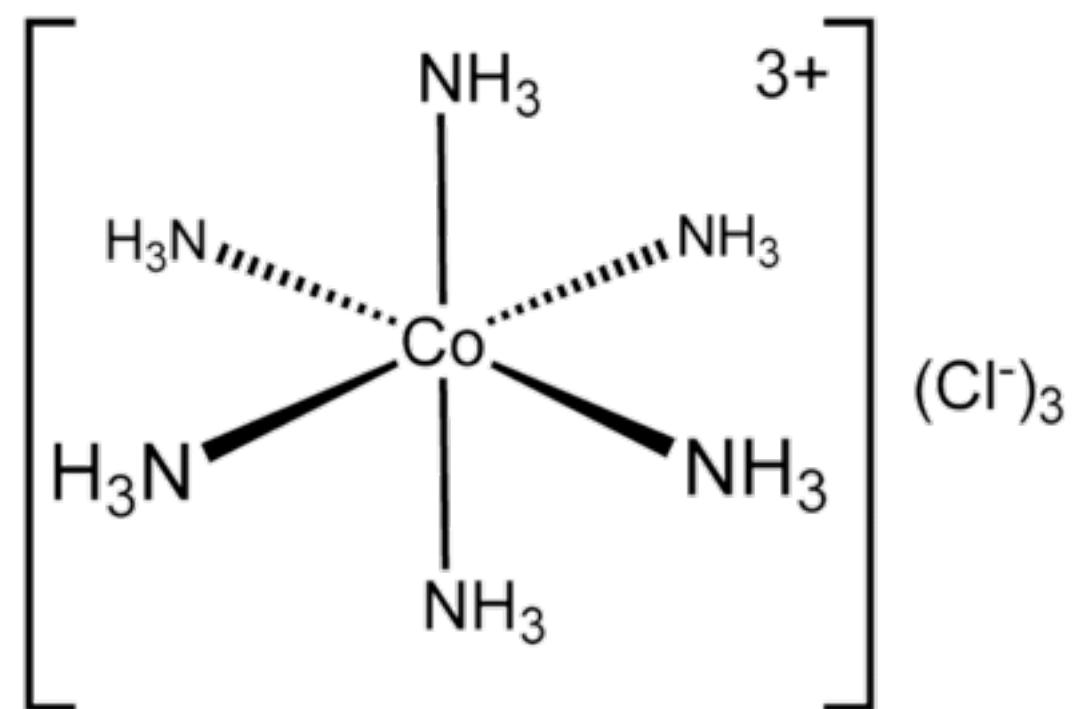
Active Sites in Biology - A Chemist's View



how do these unique coordination environments enable different functions?

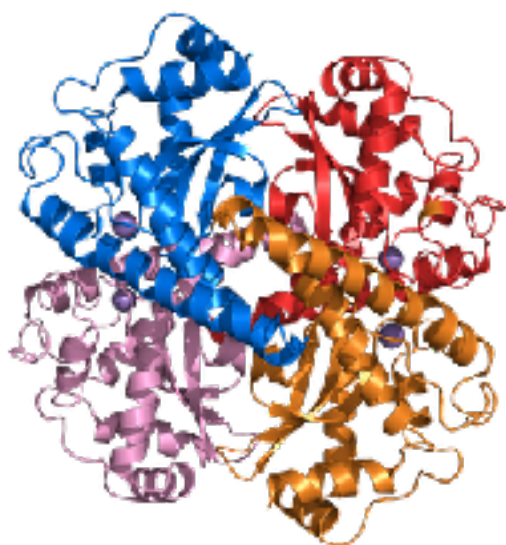


Alfred Werner (1866-1919)

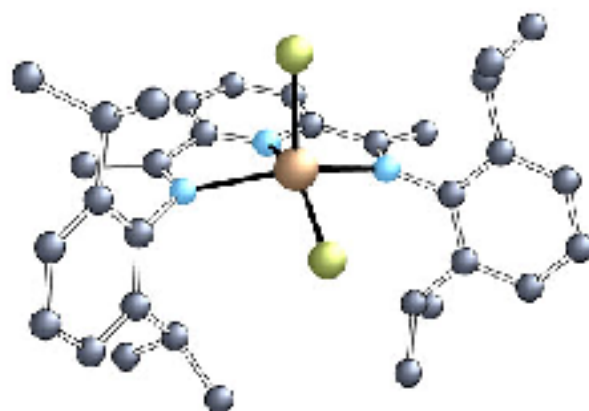


The Magic of d- and f-Block Elements

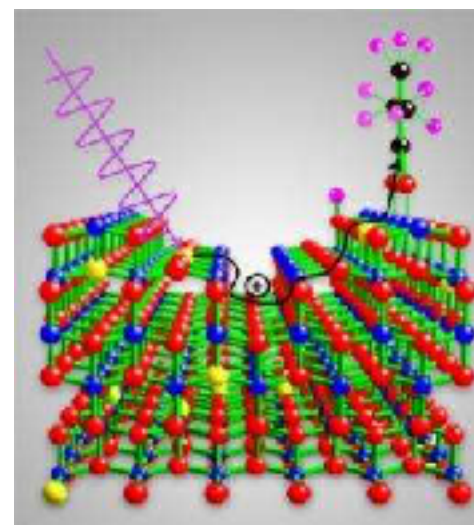
Metalloenzymes



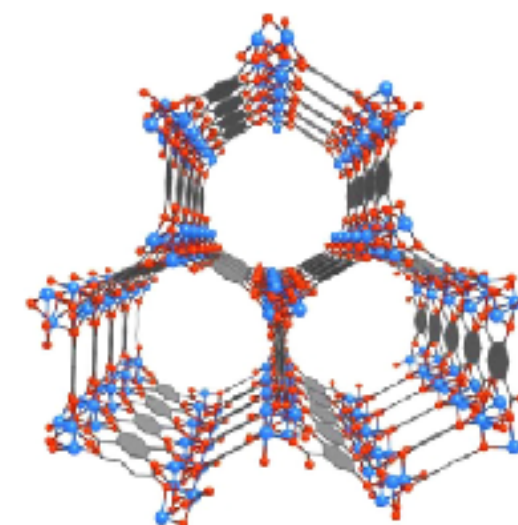
Homogeneous catalysts



Trans. Metal Oxides



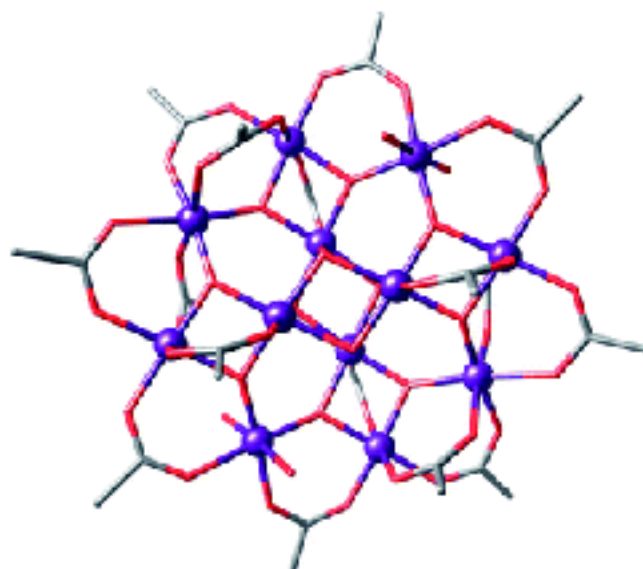
MOFs



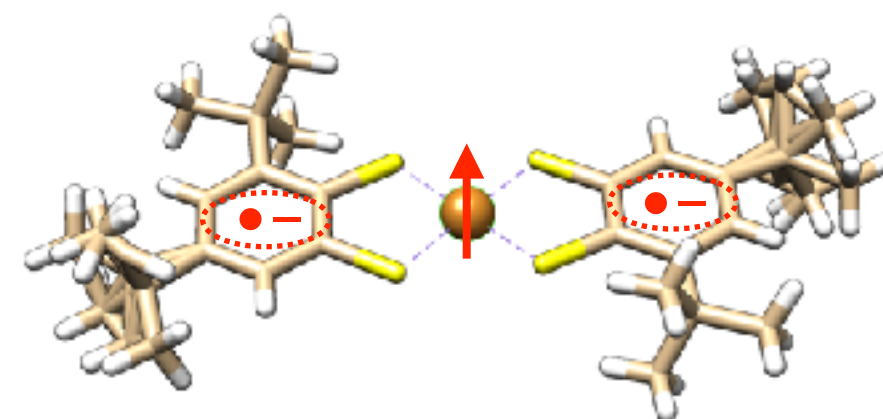
Rare-Earth Luminescence



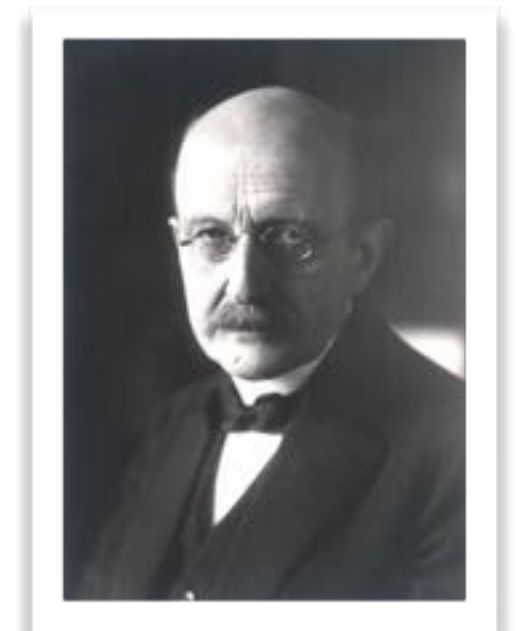
Single-Molecule Magnets



Metal-Radicals

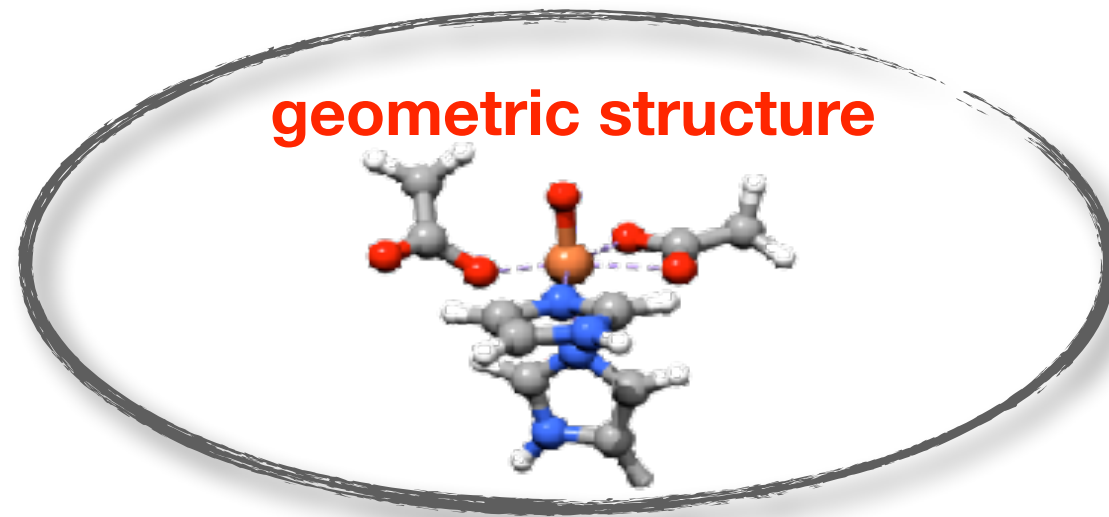


„The only source of knowledge is experiment. The rest is poetry, imagination“



(Max Planck)

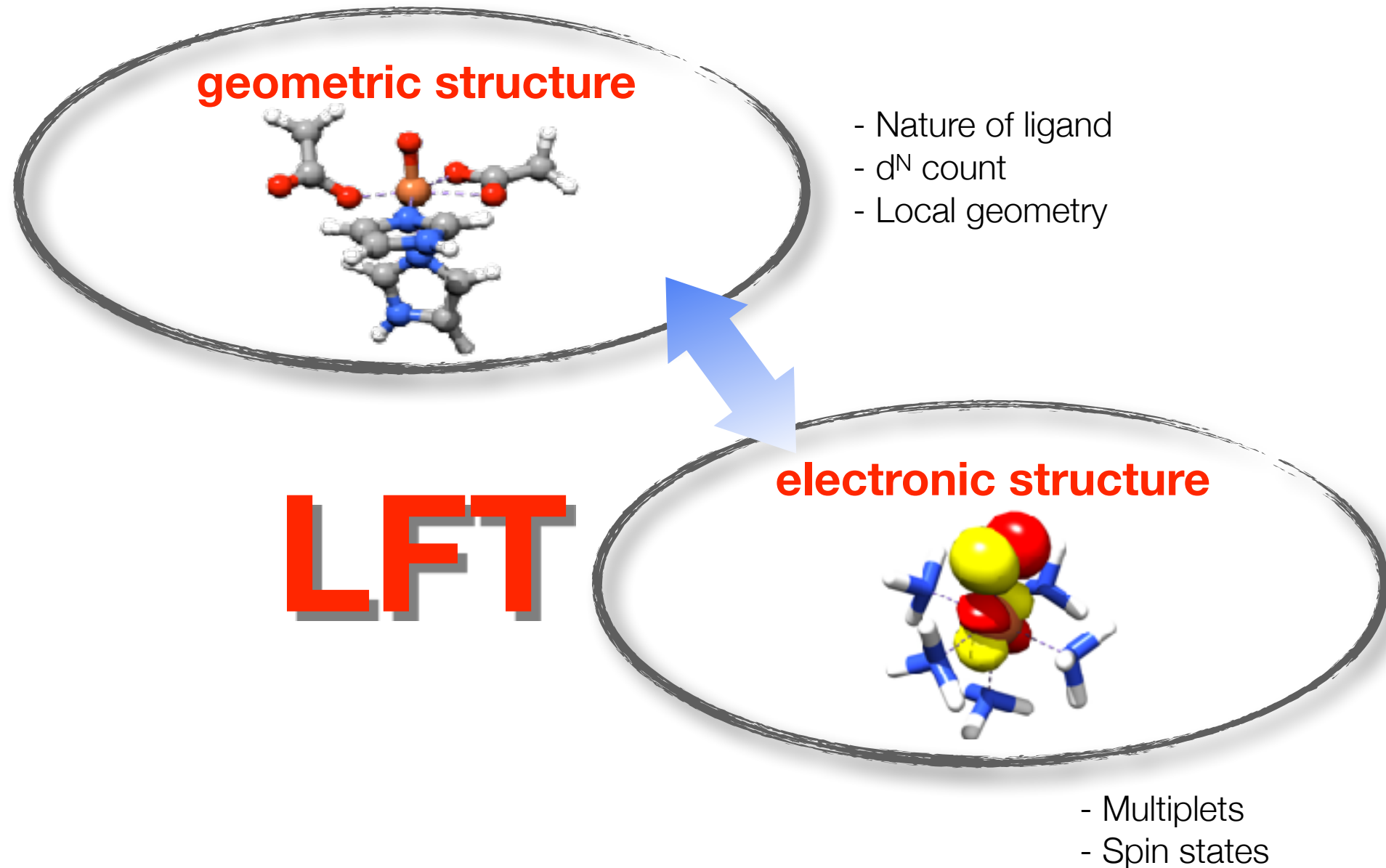
Ligand Field Theory as the Guide



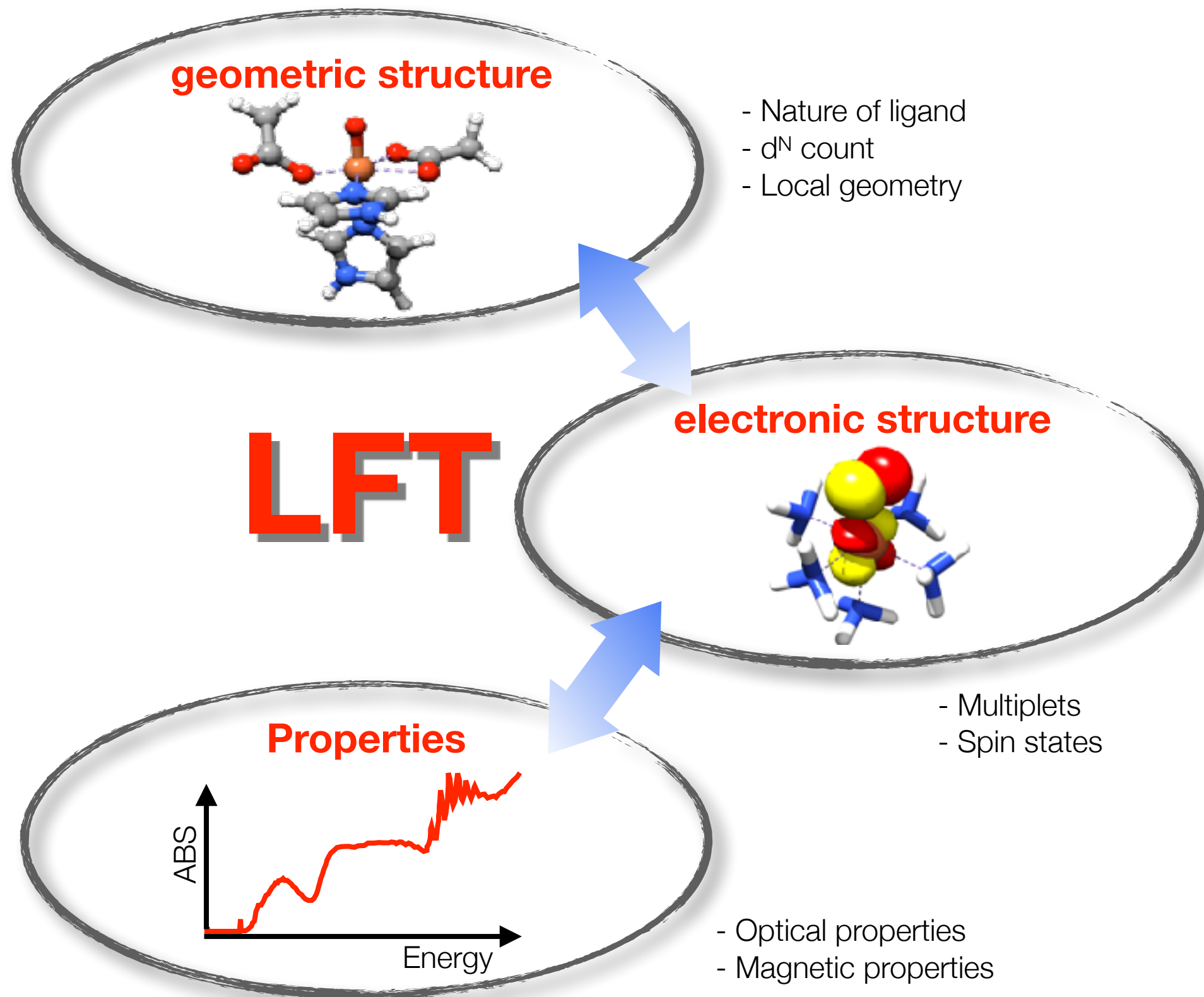
- Nature of ligand
- d^N count
- Local geometry

LFT

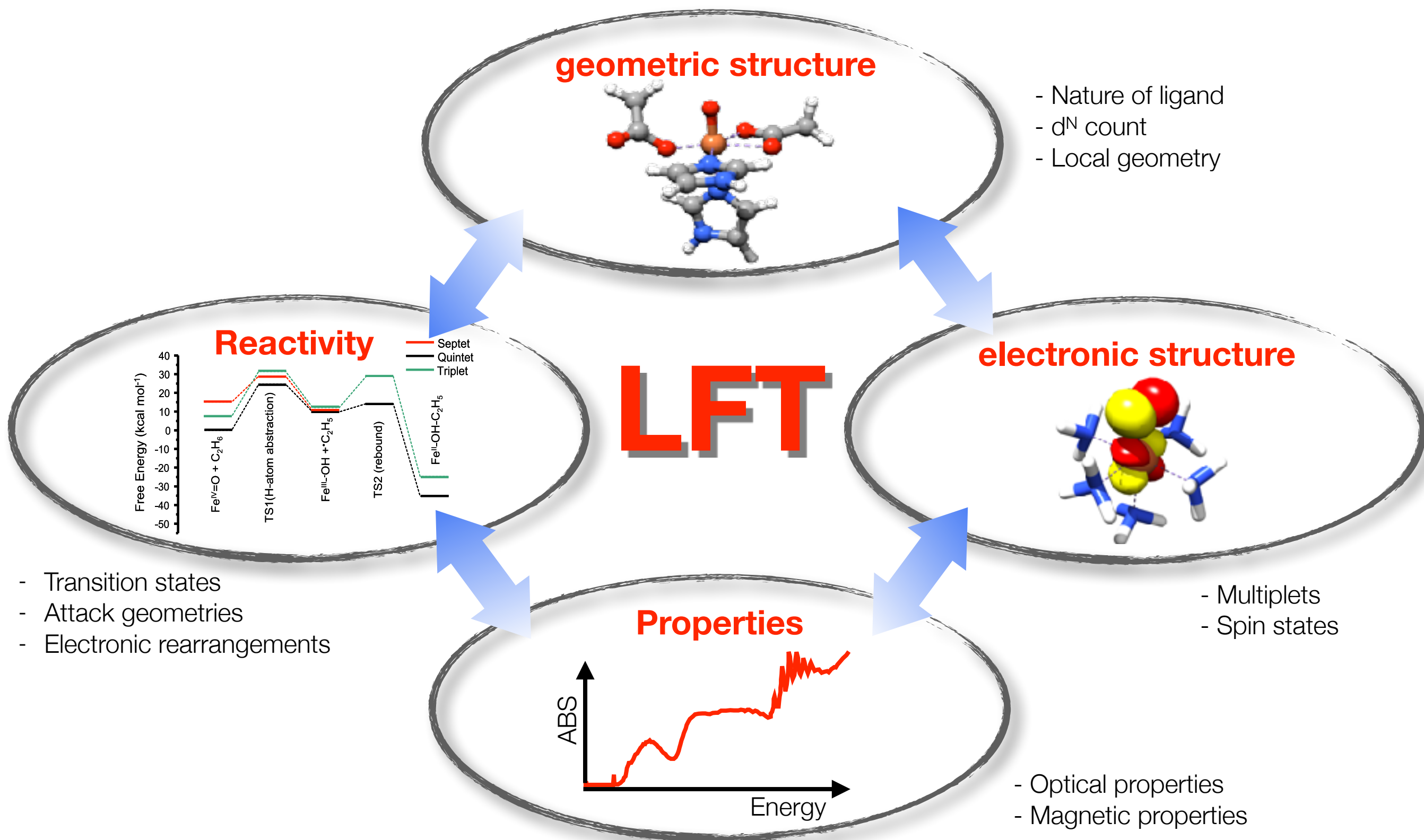
Ligand Field Theory as the Guide



Ligand Field Theory as the Guide



Ligand Field Theory as the Guide



What is Ligand Field Theory ?

★ Ligand Field Theory is:

- ▶ A **semi-empirical theory** that applies to a **CLASS** of substances (transition metal complexes).
- ▶ A **LANGUAGE** in which a vast number of experimental facts can be rationalized and discussed.
- ▶ A **MODEL** that applies only to a restricted part of reality.



Hans Primas
(1928-2014)

★ Ligand Field Theory is NOT:

- ▶ An **ab initio theory** that lets one predict the properties of a compound 'from scratch'
- ▶ A physically rigorous treatment of transition metal electronic structure

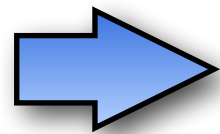
Many Electron Atoms and the ‚Aufbau‘ Principle

H																	He
Li	Be											B	C	N	O	F	Ne
Na	Mg											Al	Si	P	S	Cl	Ar
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr
					Mo												
					W												



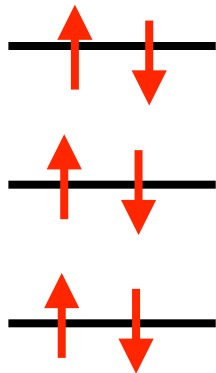
States of Atoms and Molecules

- ★ Atoms and Molecule exist in electronic **STATES**
- ★ A **STATE** of an atom or molecule may be characterized by four criteria:
 1. The distribution of the electrons among the available orbitals (the electron **CONFIGURATION**) (A set of occupation numbers)
 2. The overall **SYMMETRY** of the STATE (Γ Quantum Number)
 3. The **TOTAL SPIN** of the STATE (S-Quantum Number)
 4. The **PROJECTION** of the Spin onto the Z-axis (M_S Quantum Number)

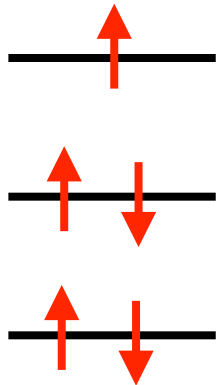
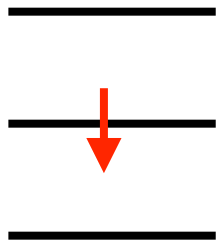


$$|n(SM)(\Gamma M_\Gamma)\rangle$$

An Example

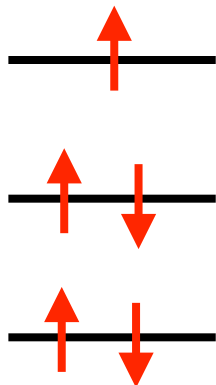
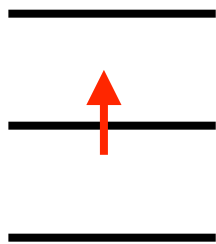


An Example



One Singlet

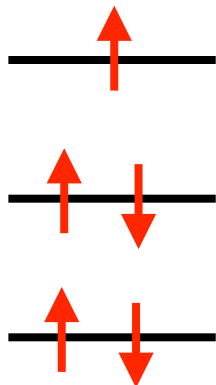
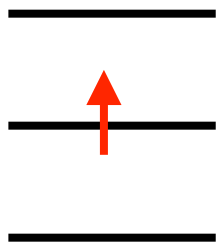
An Example



One Singlet

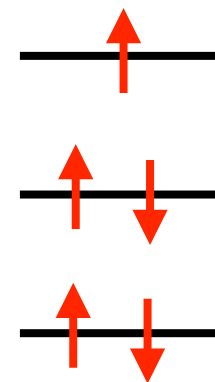
One Triplet

An Example

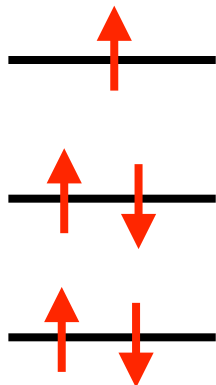
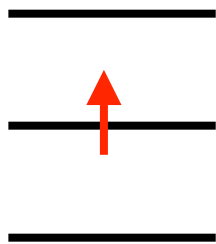


One Singlet

One Triplet

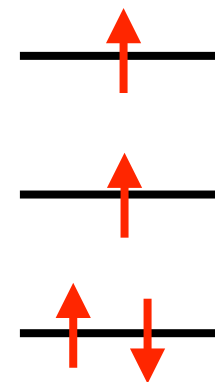
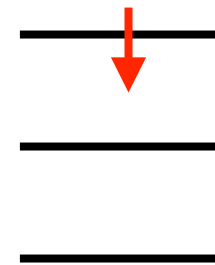


An Example



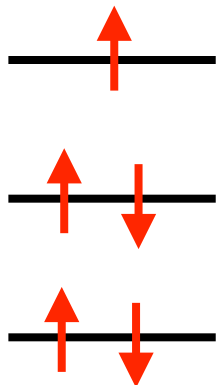
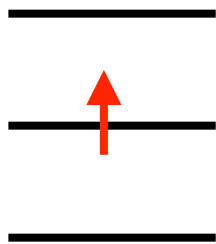
One Singlet

One Triplet



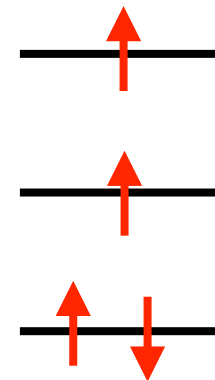
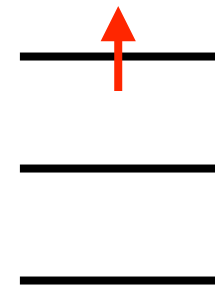
One Doublet

An Example



One Singlet

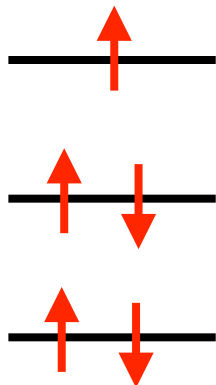
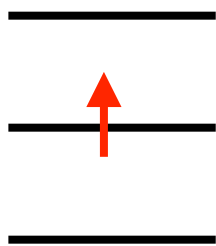
One Triplet



One Doublet

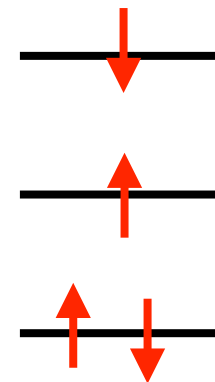
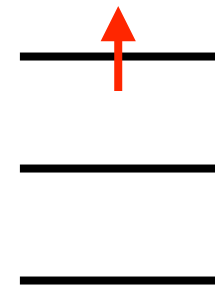
One Quartet

An Example



One Singlet

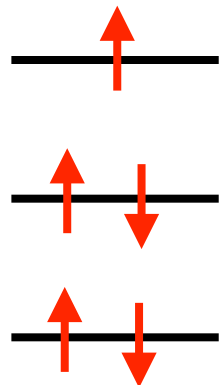
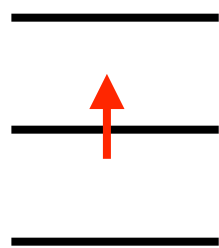
One Triplet



TWO Doublets

One Quartet

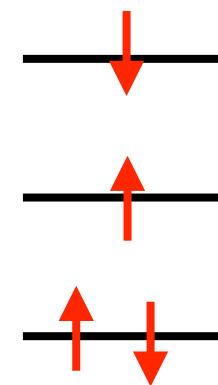
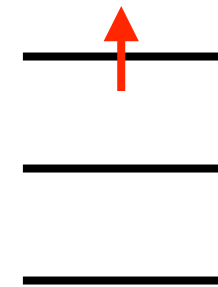
An Example



One Singlet

One Triplet

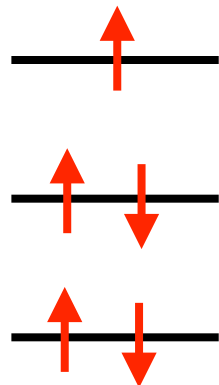
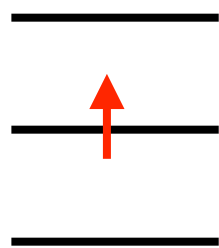
Configurations with
several unpaired
electrons generate
several states



TWO Doublets

One Quartet

An Example

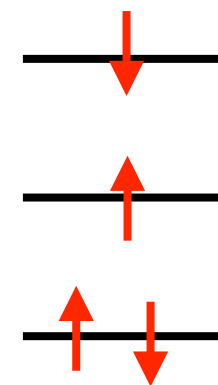
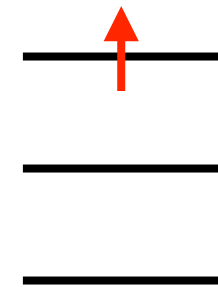


One Singlet

One Triplet

Configurations with
several unpaired
electrons generate
several states

Only STATES are
observable,
ORBITALS are NOT!



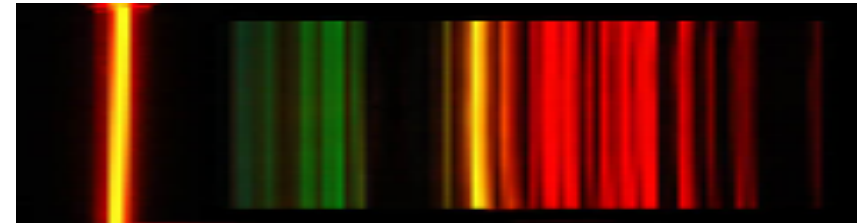
TWO Doublets

One Quartet

Atoms: Atomic „Russel-Saunders“ Terms

Atomic Term Symbol:

$$2S+1L$$



Rules:

- ▶ A L-Term is $2L+1$ fold **orbitally degenerate** and $(2S+1)(2L+1)$ fold degenerate in total
- ▶ $L=0,1,2,3,4\dots$ terms are given the symbols **S,P,D,F,G,...**
- ▶ Terms of a given configuration with higher S are lower in energy (**Hund I**)
- ▶ Terms with given configuration and equal spin have the higher L lower in energy (**Hund II**)

Examples for d^N Configurations:



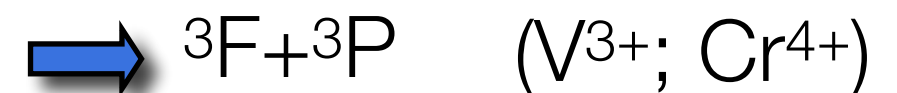
$2S+1=2;$
5 equivalent ways to put one e^-
into five degenerate orbitals



$2S+1=6;$
1 equivalent ways to put five e^-
with parallel spin in five orbitals



$2S+1=3;$
10 Ways to put two e^- with parallel
spin in five orbitals



Note: „holes“ create the same terms as „electrons“

Molecules: Symmetry and Group Theory

- ★ A Molecule can be classified according to the operations that turn the molecule into itself (=symmetry operations), i.e. rotations, improper rotations, inversion, reflection.
- ★ The precise mathematical formulation is part of „group theory“
- ★ The results is that states can be classified according to their „irreducible representation“ („symmetry quantum number“)

Rules for naming „irreducible representations“:

- ▶ Small Letters : Reserved for **orbitals** (One-electron level)
- ▶ Capital Letters : Reserved for **states** (Many electron level)
- ▶ T,t : Triply degenerate level (Mulliken notation)
- ▶ E,e : Doubly degenerate level
- ▶ A,B : Non-degenerate Levels

Term-Symbol:

$2S+1\Gamma$

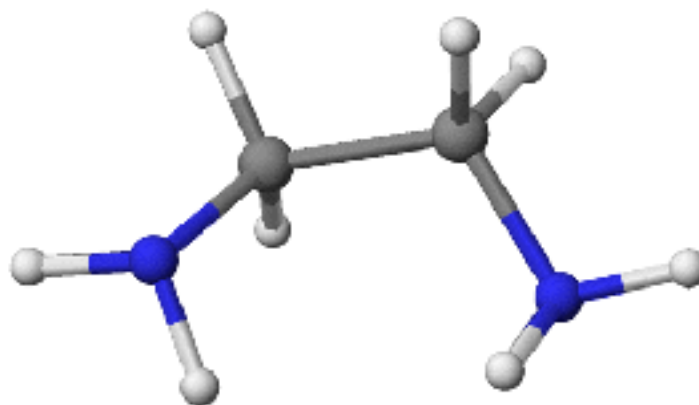
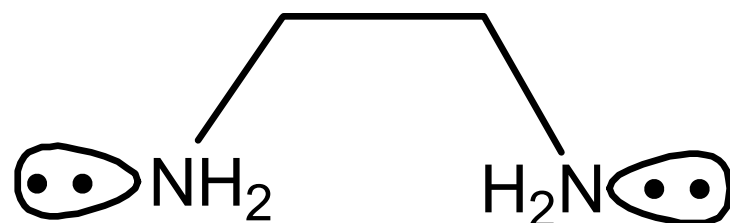
$2S+1$: „Multiplicity“ = Spin Degeneracy

Γ : „Irreducible Representation“

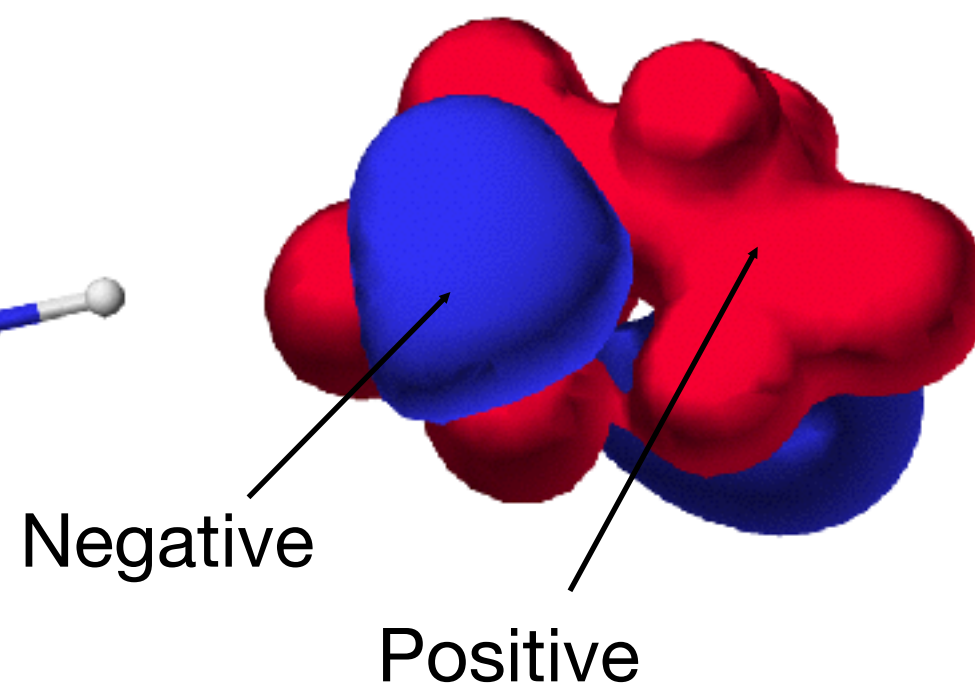
Principles of Ligand Field Theory



Example:

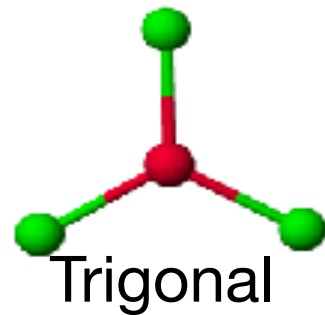


Electrostatic Potential

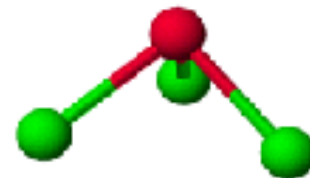


Complex Geometries

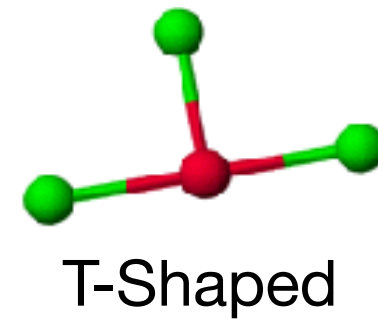
3



D_{3h}

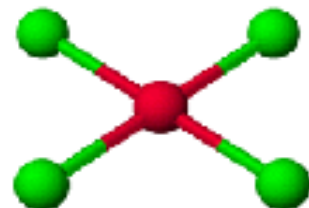


D_{3h}

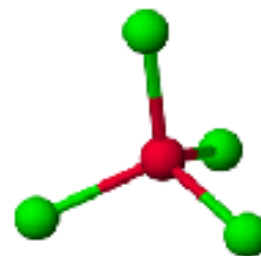


C_{2v}

4



D_{4h}



T_d

5



C_{4v}



D_{3h}

6

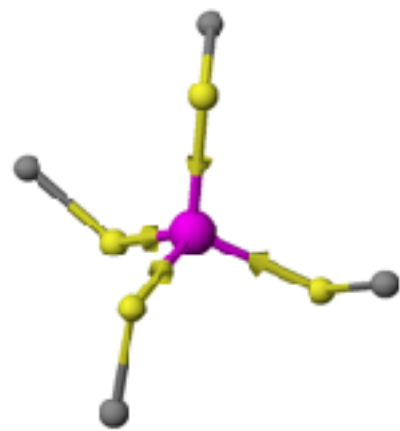
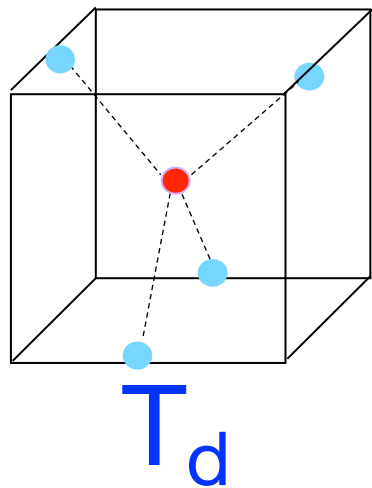


O_h

Coordination Geometries

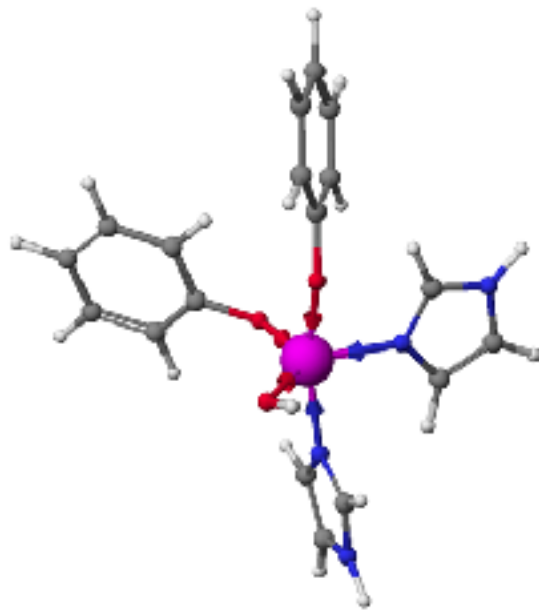
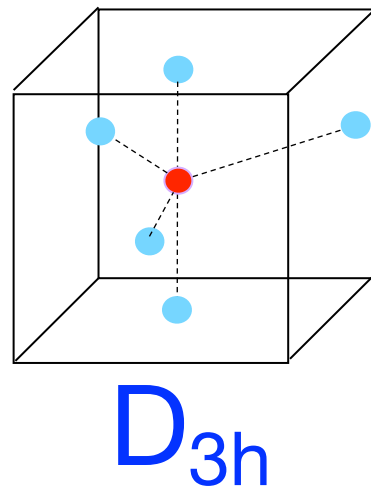
- Approximate Symmetries Observed in Enzyme Active Sites -

Tetrahedral



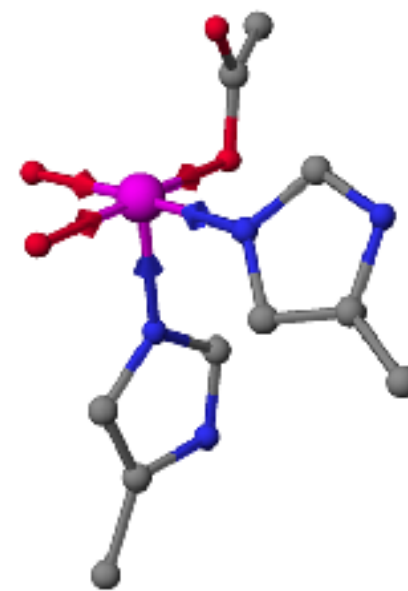
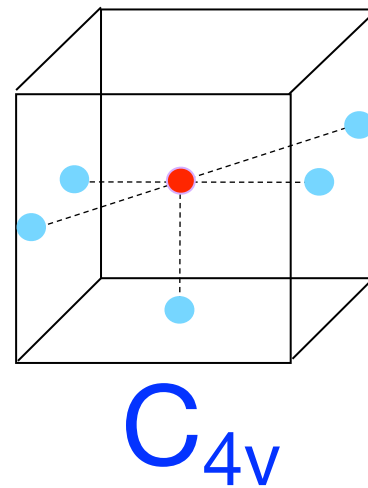
Rubredoxin

Trigonal Bipyramidal



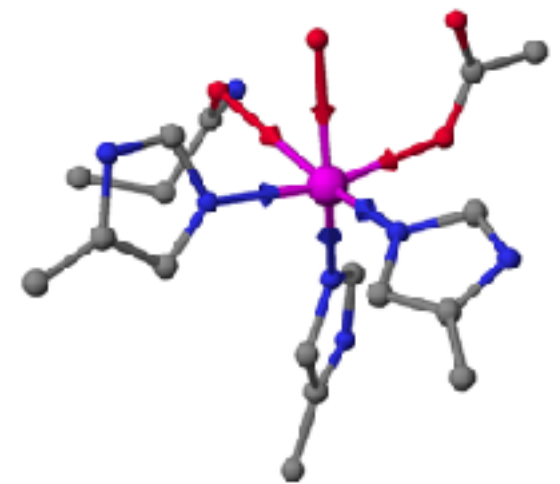
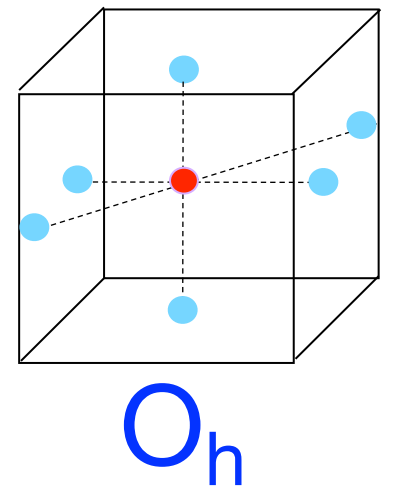
3,4-PCD

Tetragonal Pyramidal



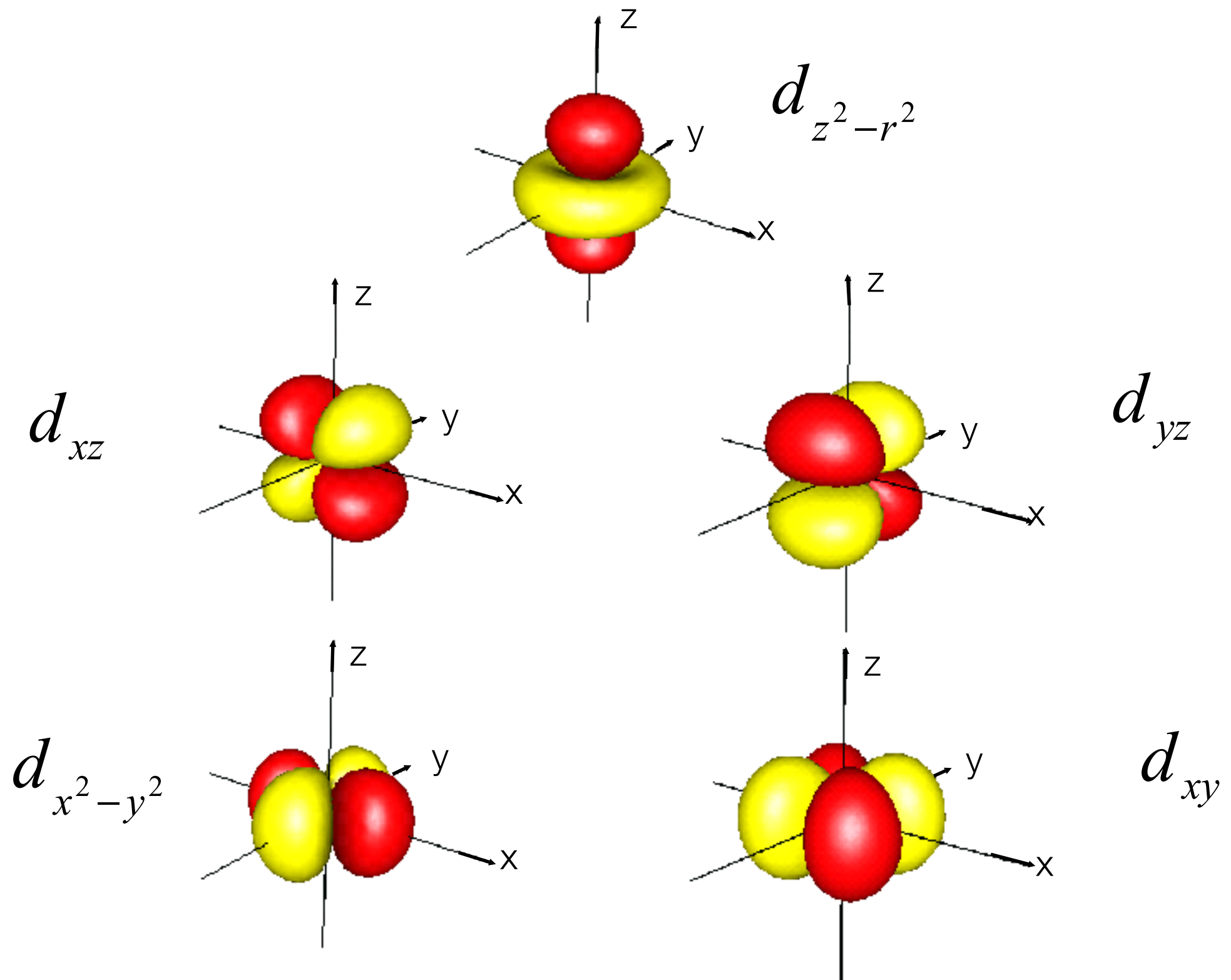
Tyrosine Hydroxylase

Octahedral



Lipxygenase

The Shape of Orbitals



A Single d-Electron in an Octahedral Field

The Negatively Charged Ligands Produce an Electric Field+Potential



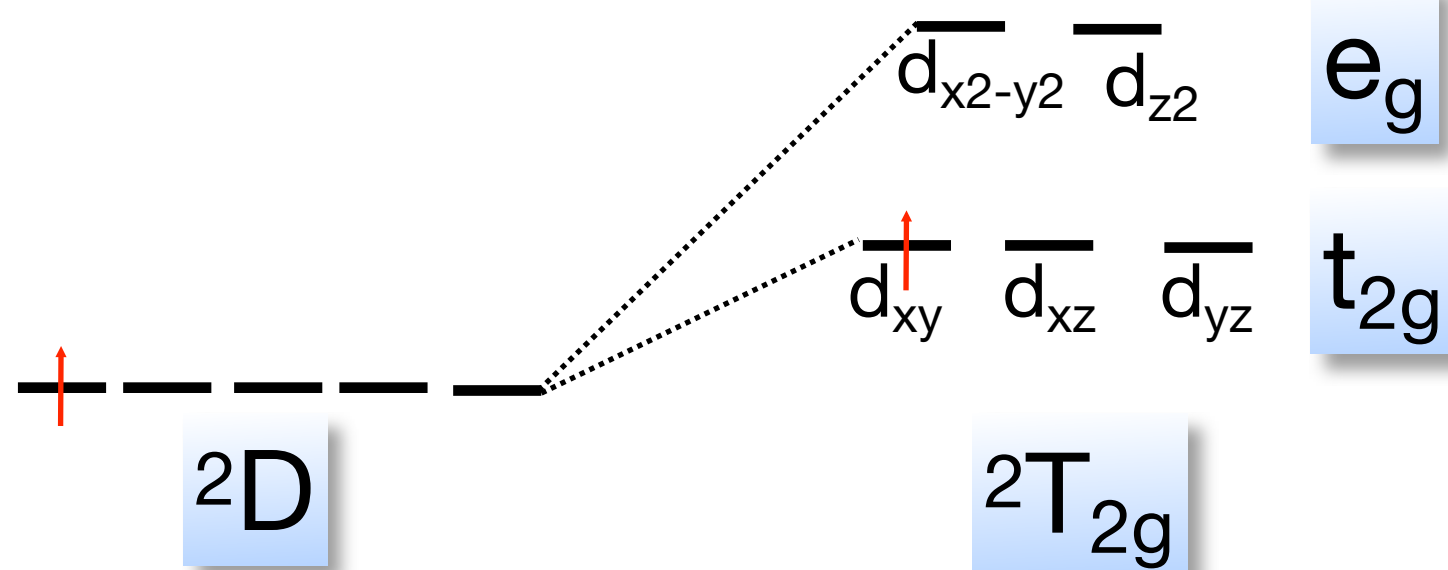
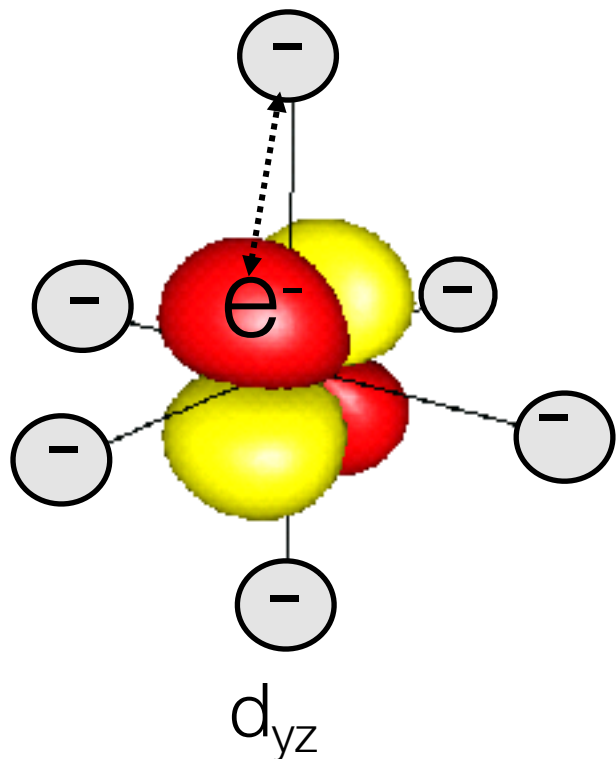
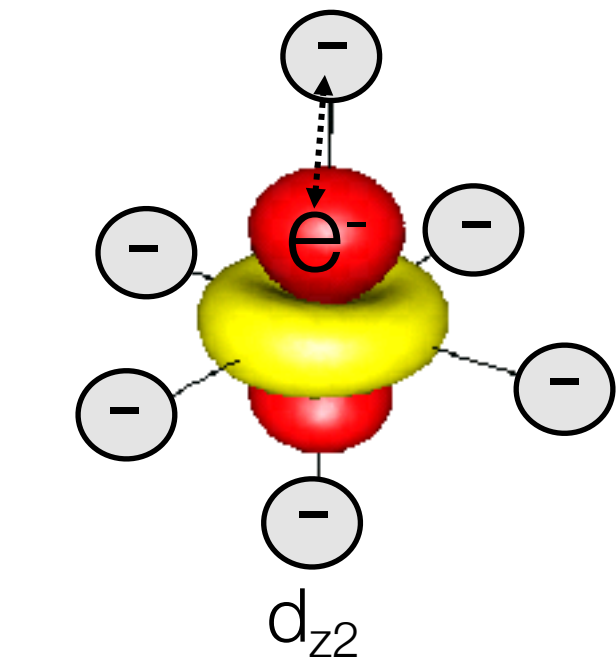
The Field Interacts with the d-Electrons on the Metal (Repulsion)



The Interaction is NOT Equal for All Five d-Orbitals



1. The Spherical Symmetry of the Free Ions is Lifted
2. The d-Orbitals Split in Energy
3. The Splitting Pattern Depends on the Arrangement of the Ligands



Making Ligand Field Theory Quantitative?



Hans Bethe
1906-2005

Charge Distribution of Ligand Charges: $\rho(\mathbf{R}) = \sum_{i=1}^{N_L} q_i \delta(\mathbf{R} - \mathbf{R}_{L_i})$

Ligand field potential: $V_{LF}(\mathbf{r}) = \int \frac{\rho(\mathbf{R})}{|\mathbf{R} - \mathbf{r}|} d\mathbf{R}$ $q_i = \text{charges}$

Expansion of inverse distance: $\frac{1}{|\mathbf{R} - \mathbf{r}|} = \sum_{l=0}^{\infty} \frac{4\pi}{2l+1} \frac{r_{<}^l}{r_{>}^{l+1}} \sum_{m=-l}^l S_{lm}(\mathbf{R}) S_{lm}(\mathbf{r})$ $S_{lm} = \text{real spherical harmonics}$
 $r_{<,>}$ Smaller/Larger of r and R

Insertion into the potential: $V_{LF}(\mathbf{r}) = \sum_{l=0}^{\infty} r^l \sum_{m=-l}^l S_{lm}(\mathbf{r}) A_{lm}$

„Geometry factors“: $A_{lm} = \sum_{i=1}^{N_L} \frac{4\pi}{2l+1} \frac{q_i}{R_{L_i}^{l+1}} S_{lm}(\mathbf{R}_{L_i})$

Ligand-field matrix elements: $\langle d_i | V_{LF} | d_j \rangle = - \sum_{l=0}^{\infty} \langle r^l \rangle \sum_{m=-l}^l A_{lm} \begin{pmatrix} l_i & l_j & l \\ m_i & m_j & m \end{pmatrix}$ $\begin{pmatrix} l_i & l_j & l \\ m_i & m_j & m \end{pmatrix}$

Ligand-field splitting in O_h : $\langle d_{e_g} | V_{LF} | d_{e_g} \rangle - \langle d_{t_{2g}} | V_{LF} | d_{t_{2g}} \rangle \equiv 10Dq = \frac{5}{3} \frac{q}{R^5} \langle d | r^4 | d \rangle$ =Gaunt Integral (tabulated)

➡ Don't evaluate these integrals analytically, plug in and compare to experiment! LFT is not an ab initio theory (the numbers that you will get are ultimately absurd!). What we want is a parameterized model and thus we want to leave $10Dq$ as a fit parameter. The ligand field model just tells us how many and which parameters we need what their relationship is

Making Ligand Field Theory Quantitative?

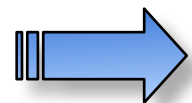


Hans Bethe
1966-2005

Charge Distribution of Ligand Charges: $\rho(\mathbf{R}) = \sum_{i=1}^{N_L} q_i \delta(\mathbf{R} - \mathbf{R}_{L_i})$

THUS

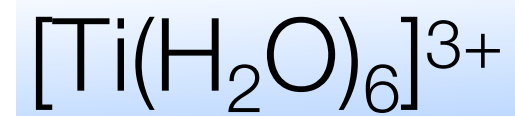
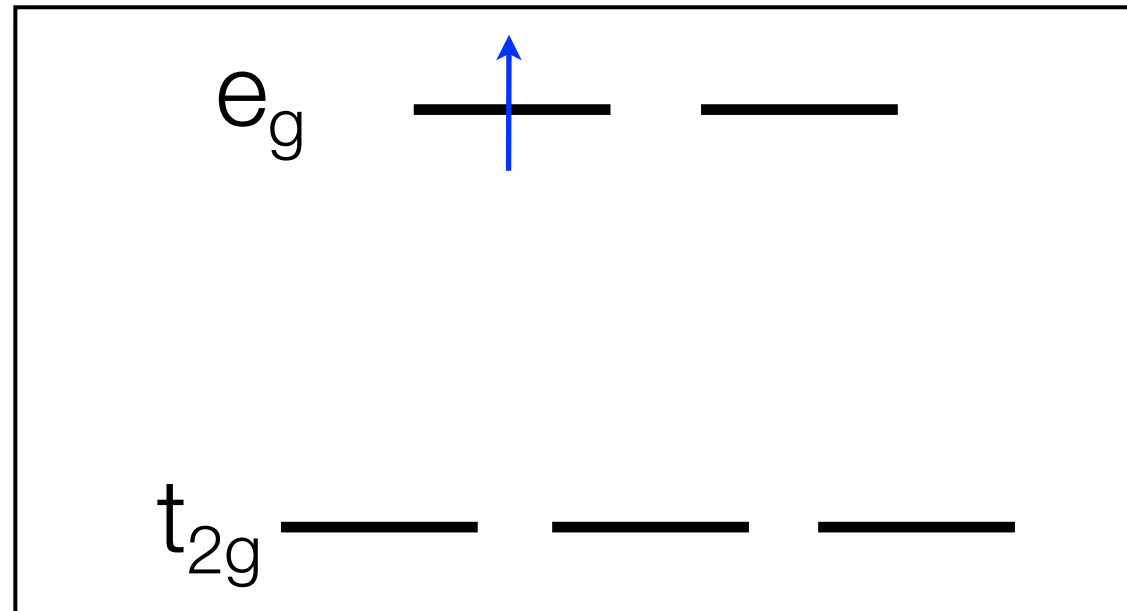
- LFT taken literally is unphysical
- LFT Parameters are not theoretically defined
- Fitting to experimental data is mandatory to determine them



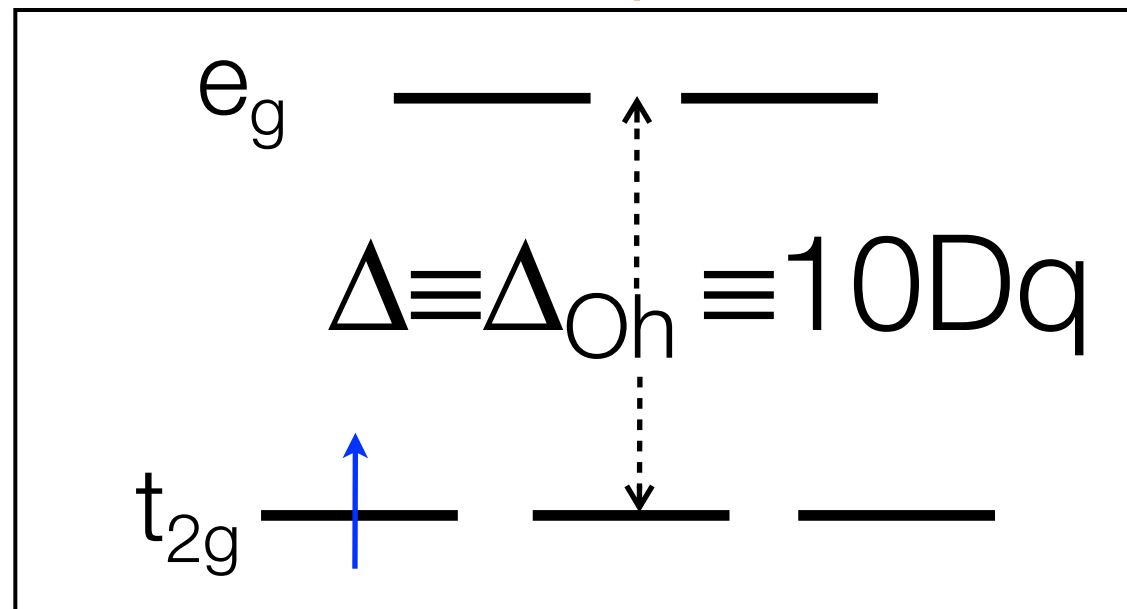
Don't evaluate these integrals analytically, plug in and compare to experiment! LFT is not an ab initio theory (the numbers that you will get are ultimately absurd!). What we want is a parameterized model and thus we want to leave $10Dq$ as a fit parameter. The ligand field model just tells us how many and which parameters we need what their relationship is

Optical Measurement of Δ : d-d Transitions

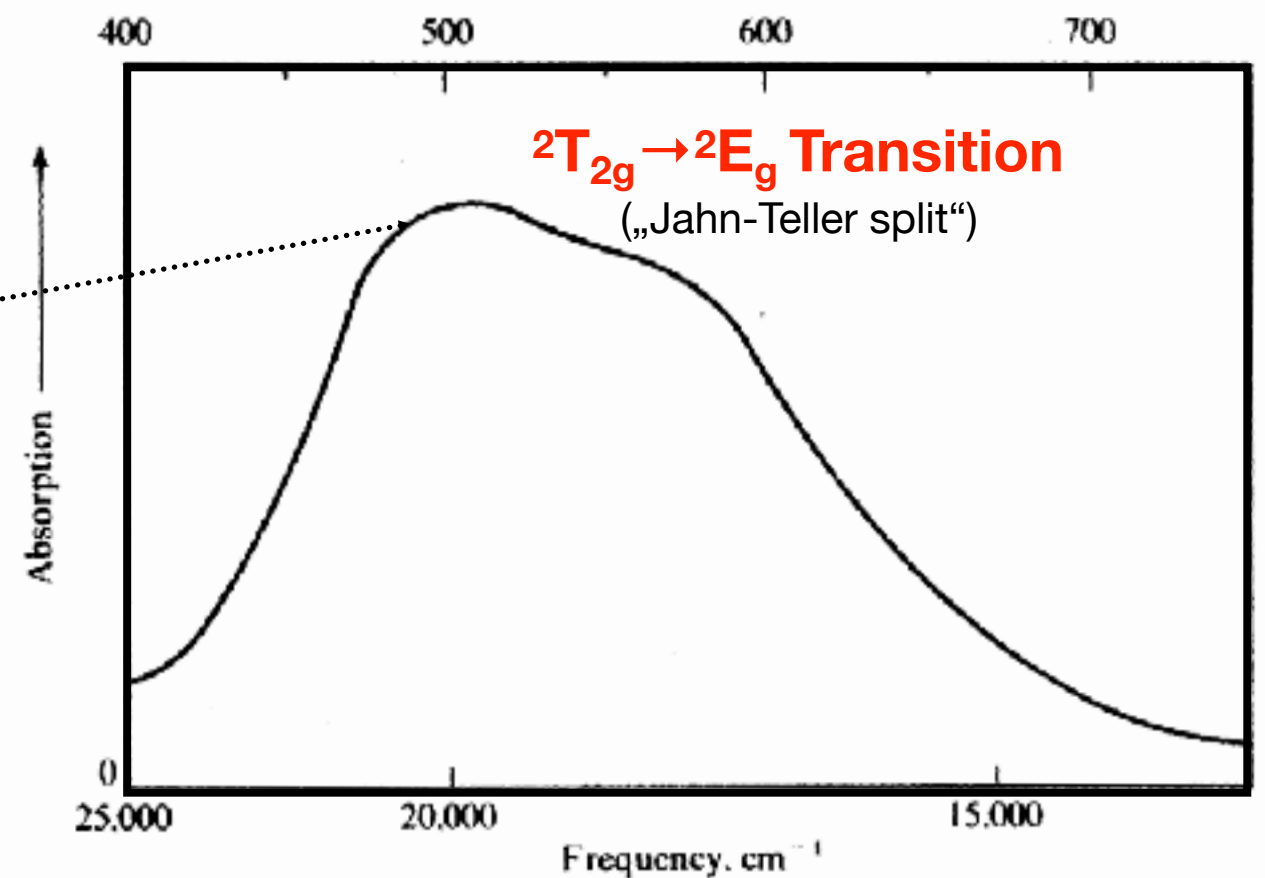
Excited State



$h\nu = \Delta_{\text{Oh}}$

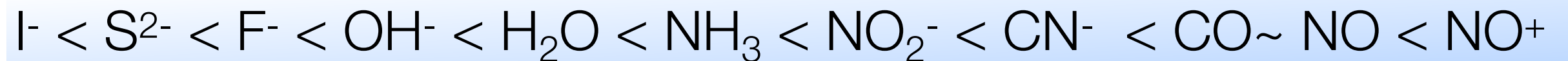


Ground State



The Spectrochemical Series

A „Chemical“ Spectrochemical Series



Δ SMALL

Δ LARGE

A „Biochemical“ Spectrochemical Series (A. Thomson)



Δ SMALL

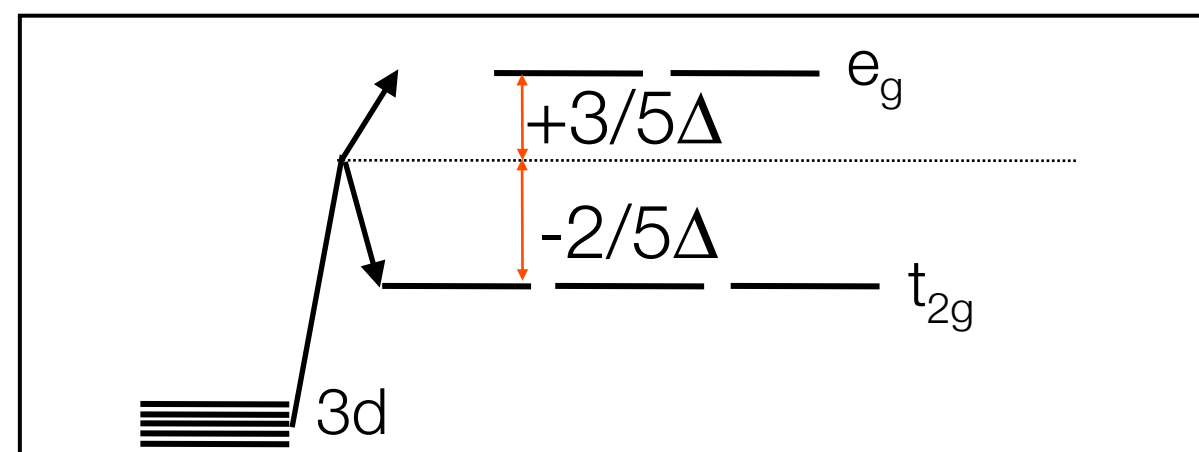
Δ LARGE

Ligand Field Stabilization Energies

Central Idea:

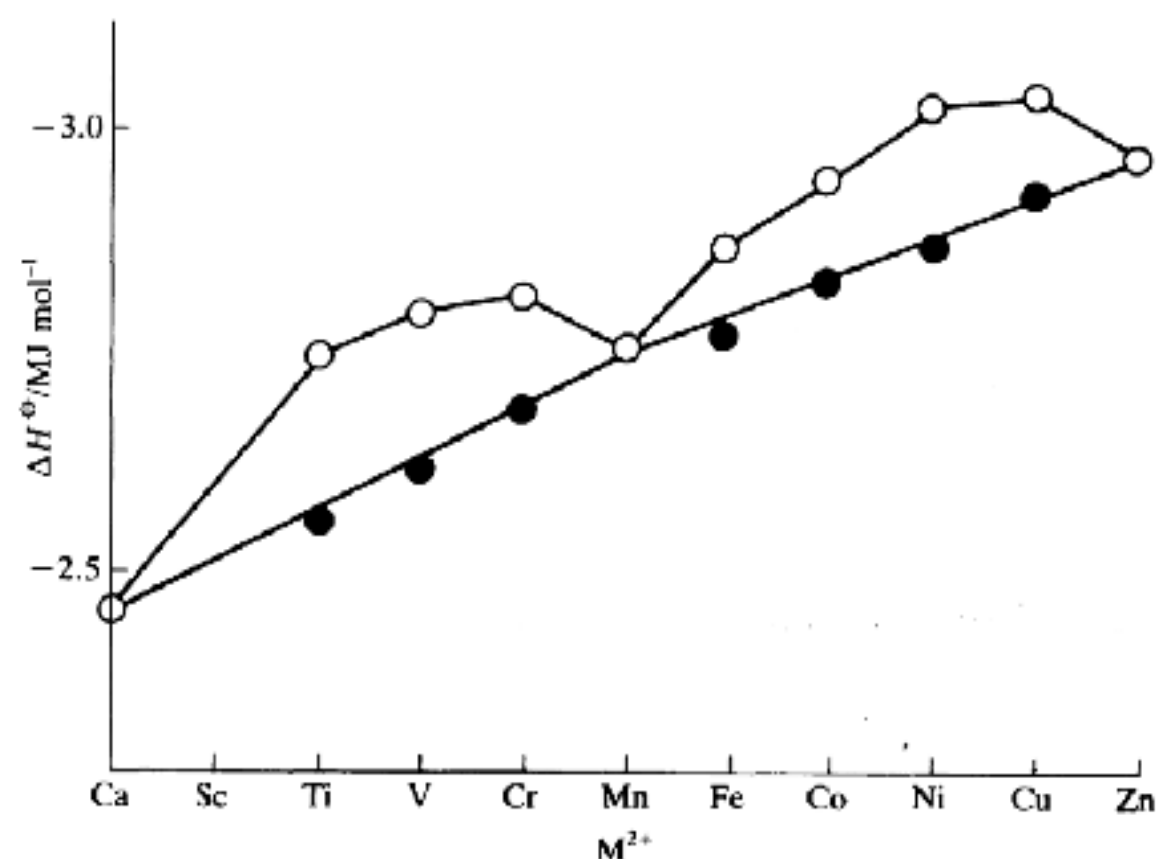
- ➔ Occupation of t_{2g} orbitals stabilizes the complex while occupation of e_g orbitals destabilizes it.
- ➔ Ligand Field Stabilization Energy (LSFE)

d ^N	LSFE
1	$-2/5\Delta$
2	$-4/5\Delta$
3	$-6/5\Delta$
4	$-3/5\Delta$
5	0
6	$-2/5\Delta$
7	$-4/5\Delta$
8	$-6/5\Delta$
9	$-3/5\Delta$
10	0



Experimental Test:

- ➔ Hydration energies of hexaquo M^{2+}



Many Electrons in a Ligand Field: Electron Repulsion

BASIC TRUTH: Electrons REPEL Each Other



Rules:

- ▶ Electrons in the **SAME orbital** repel each other most strongly.
- ▶ Electrons of **opposite spin** repel each other more strongly than electrons of the same spin.

Consequences:

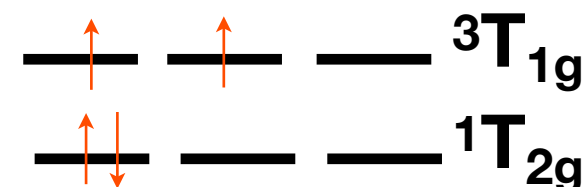
- ▶ In degenerate orbitals electrons enter first with the same spin in different orbitals (→ **Hund's Rules** in atoms!)
- ▶ A given configuration produces **several states** with different energies

Ligand Field Theory:

- ▶ Electron repulsion can be taken care of by **2 PARAMETERS: B (~1000 cm⁻¹) and C**
 - ▶ **C/B ~ 4**

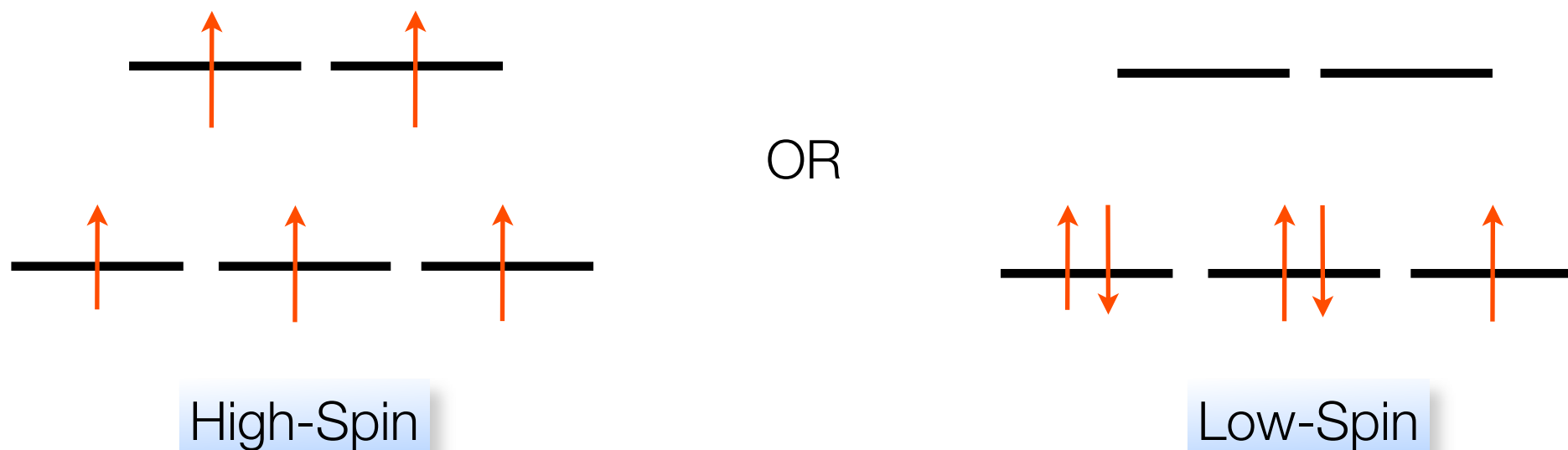
Example:

- ▶ d²-Configuration: $\Delta E(^1T_{2g} - ^3T_{1g}) \sim 6B + 2C \sim 14,000 \text{ cm}^{-1}$

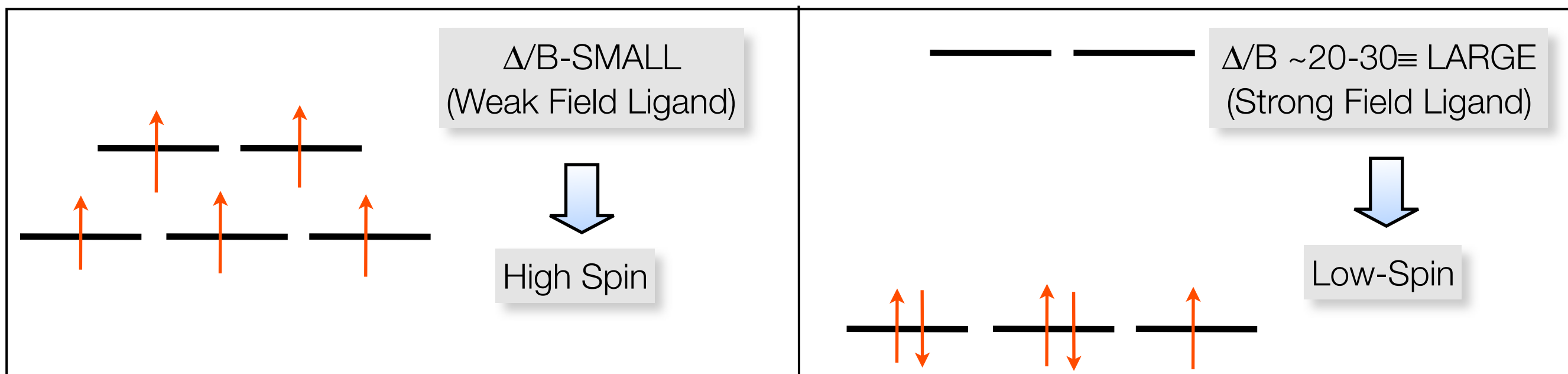


High-Spin and Low-Spin Complexes

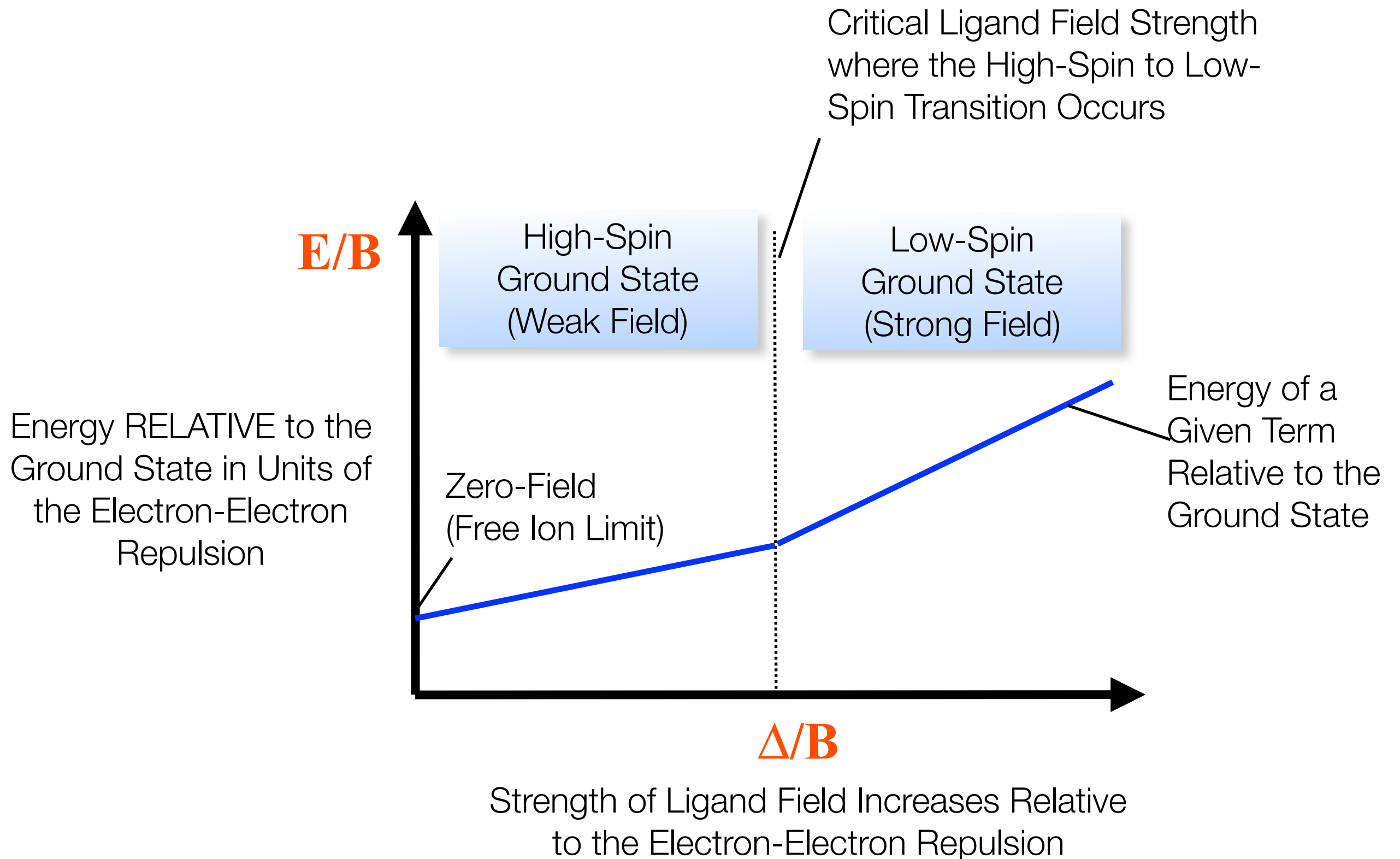
QUESTION: What determines the electron configuration?



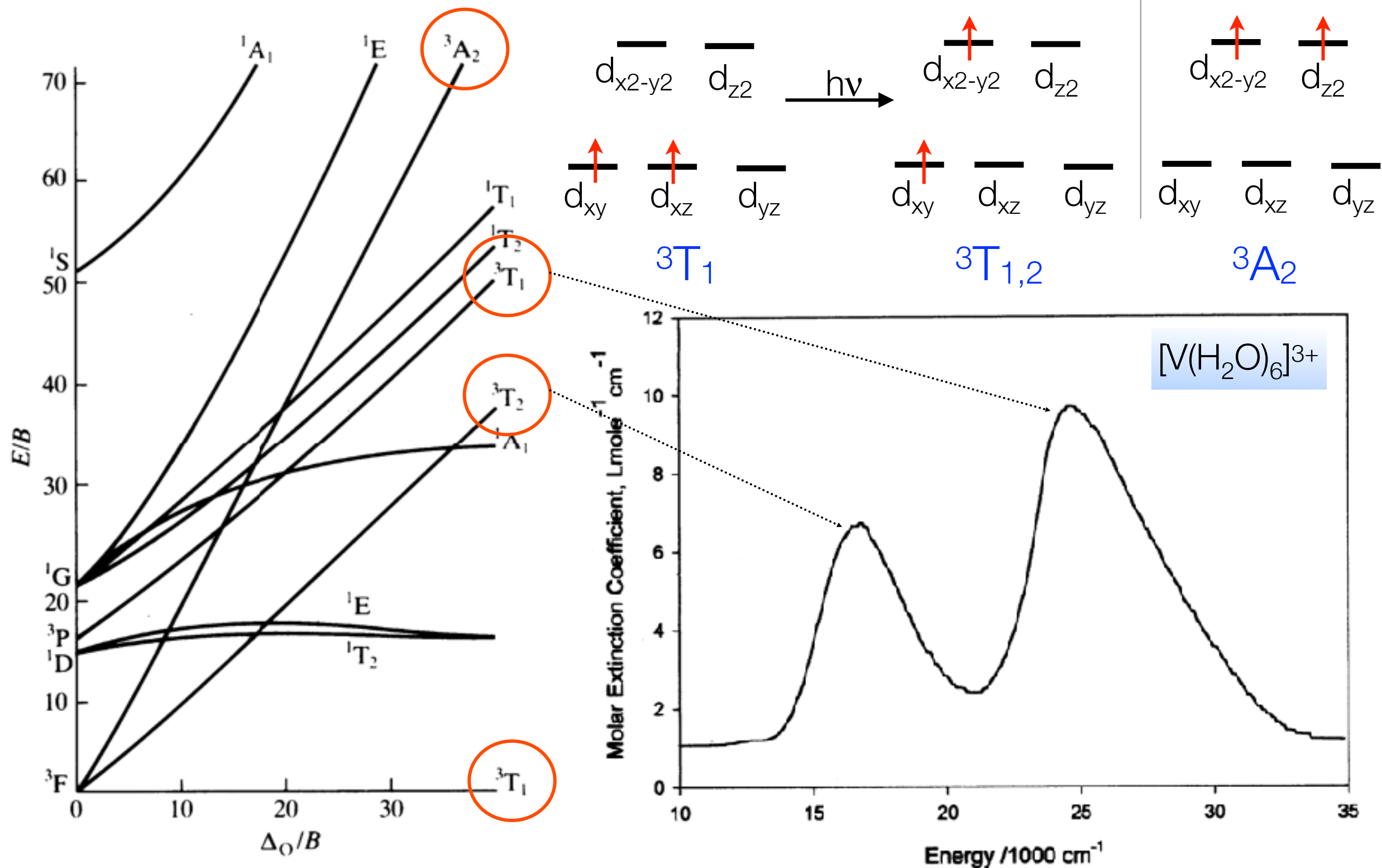
ANSWER: The balance of ligand field splitting and electron repulsion (**Spin-Pairing Energy** $P=f(B)$)



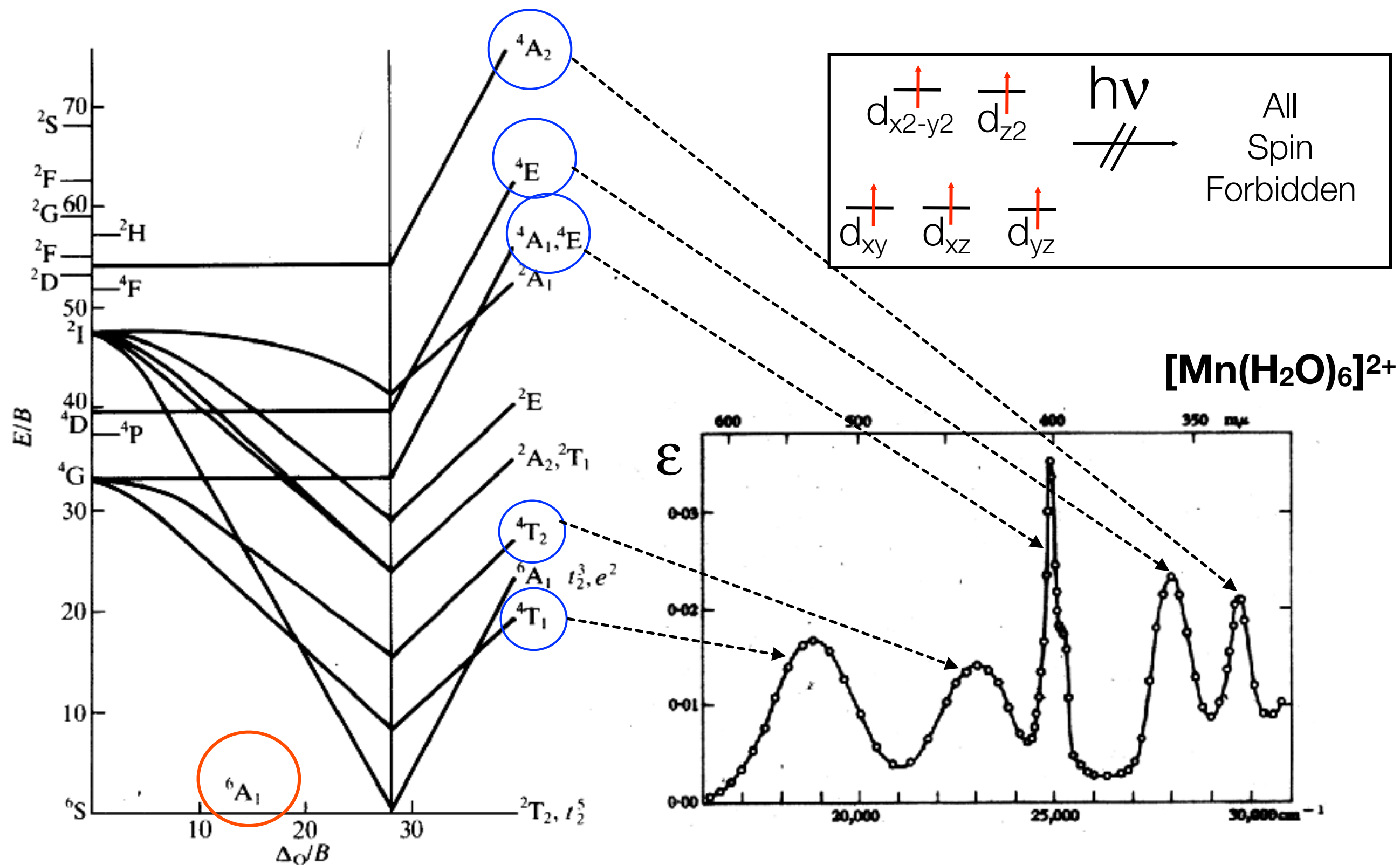
Inside Ligand Field Theory: Tanabe-Sugano Diagrams



Optical Properties:d-d Spectra of d² Ions

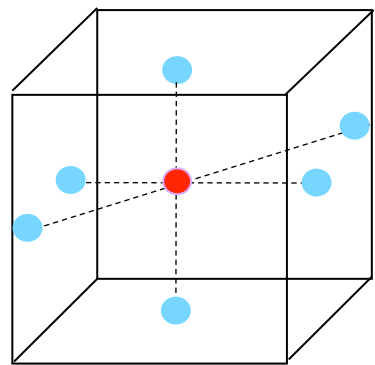


d-d Spectra of d⁵ Ions (Fe^{III}, Mn^{II})



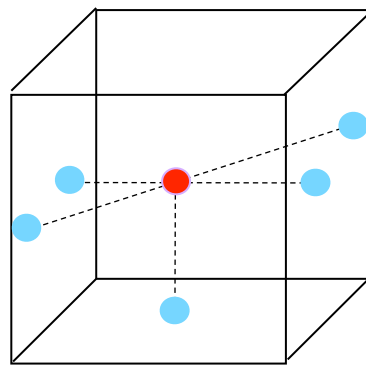
Ligand Field Splittings in Different Coordination Geometries

Octahedral



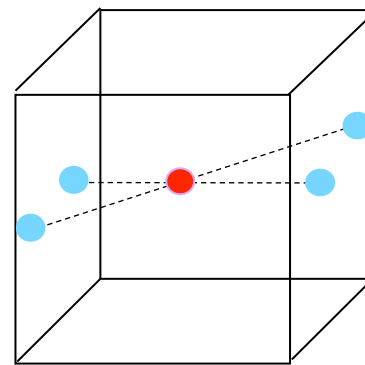
O_h

Tetragonal
Pyramidal



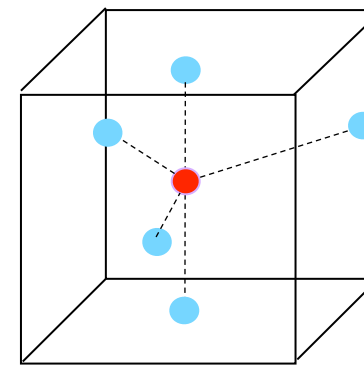
C_{4v}

Square
Planar



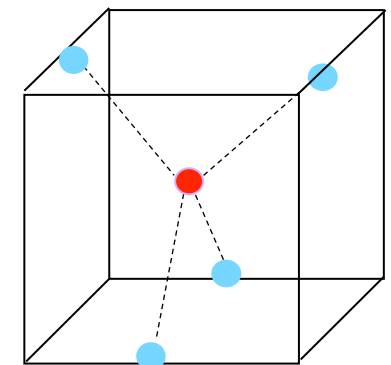
D_{4h}

Trigonal
Bipyramidal

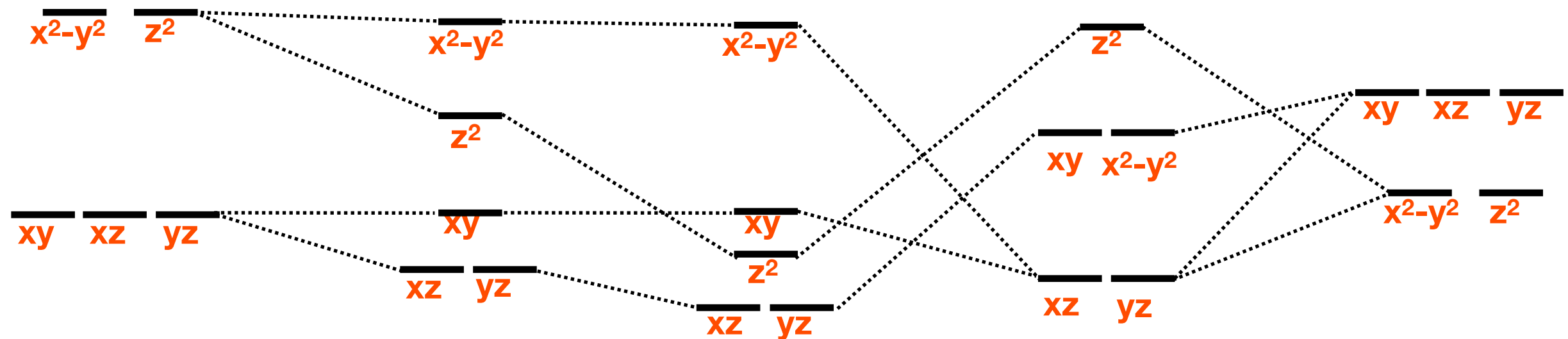


D_{3h}

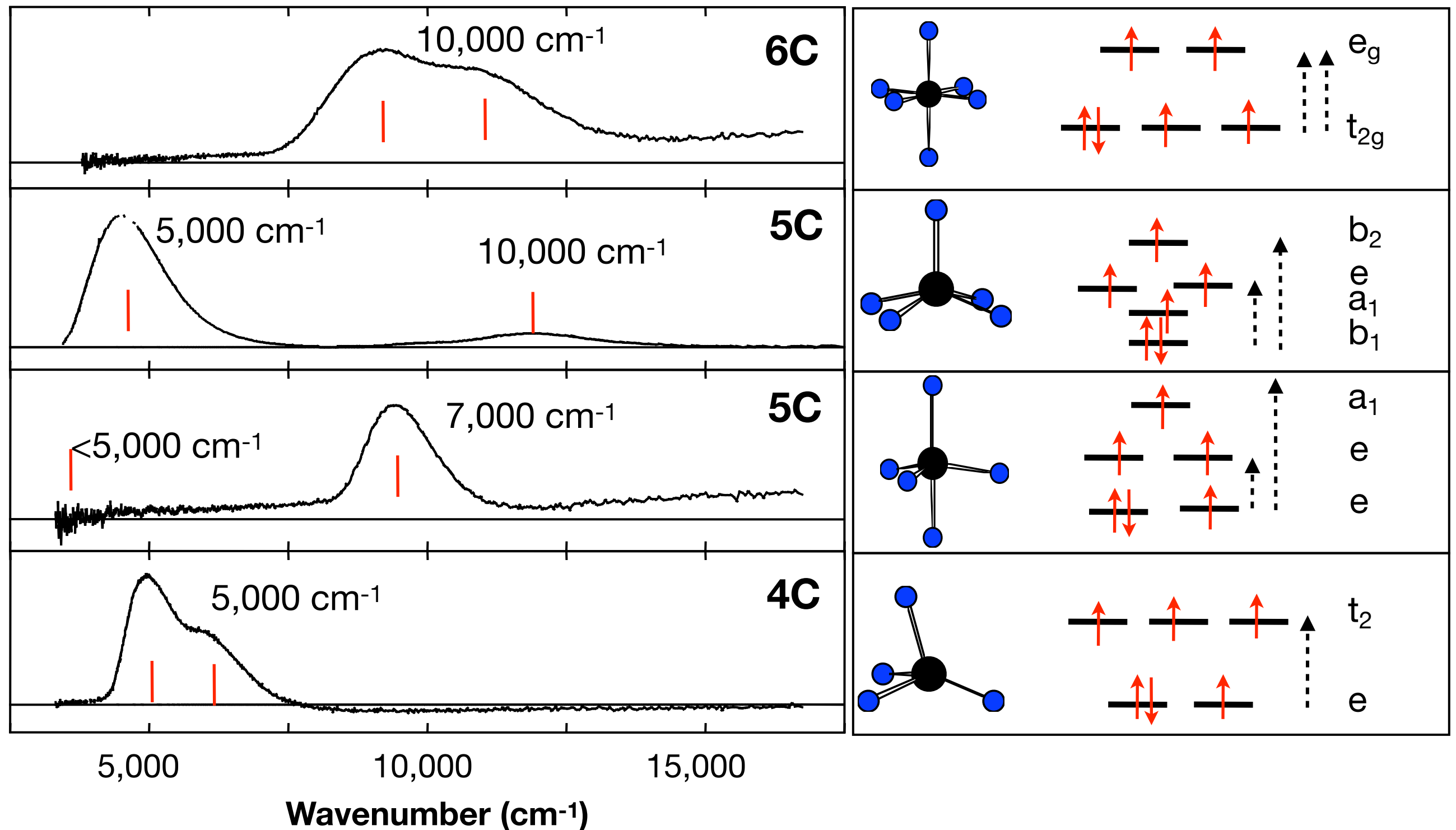
Tetrahedral



T_d

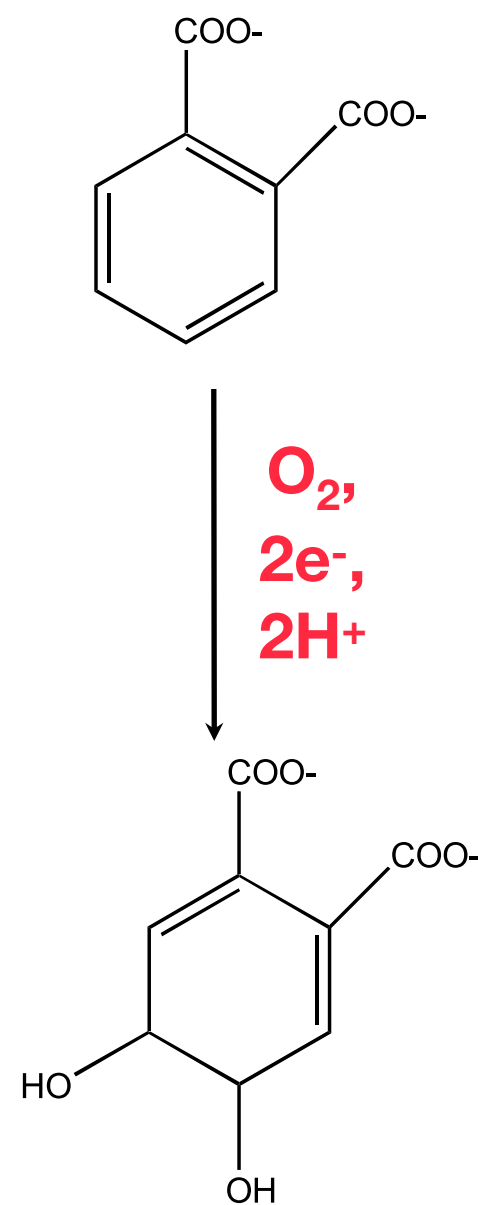
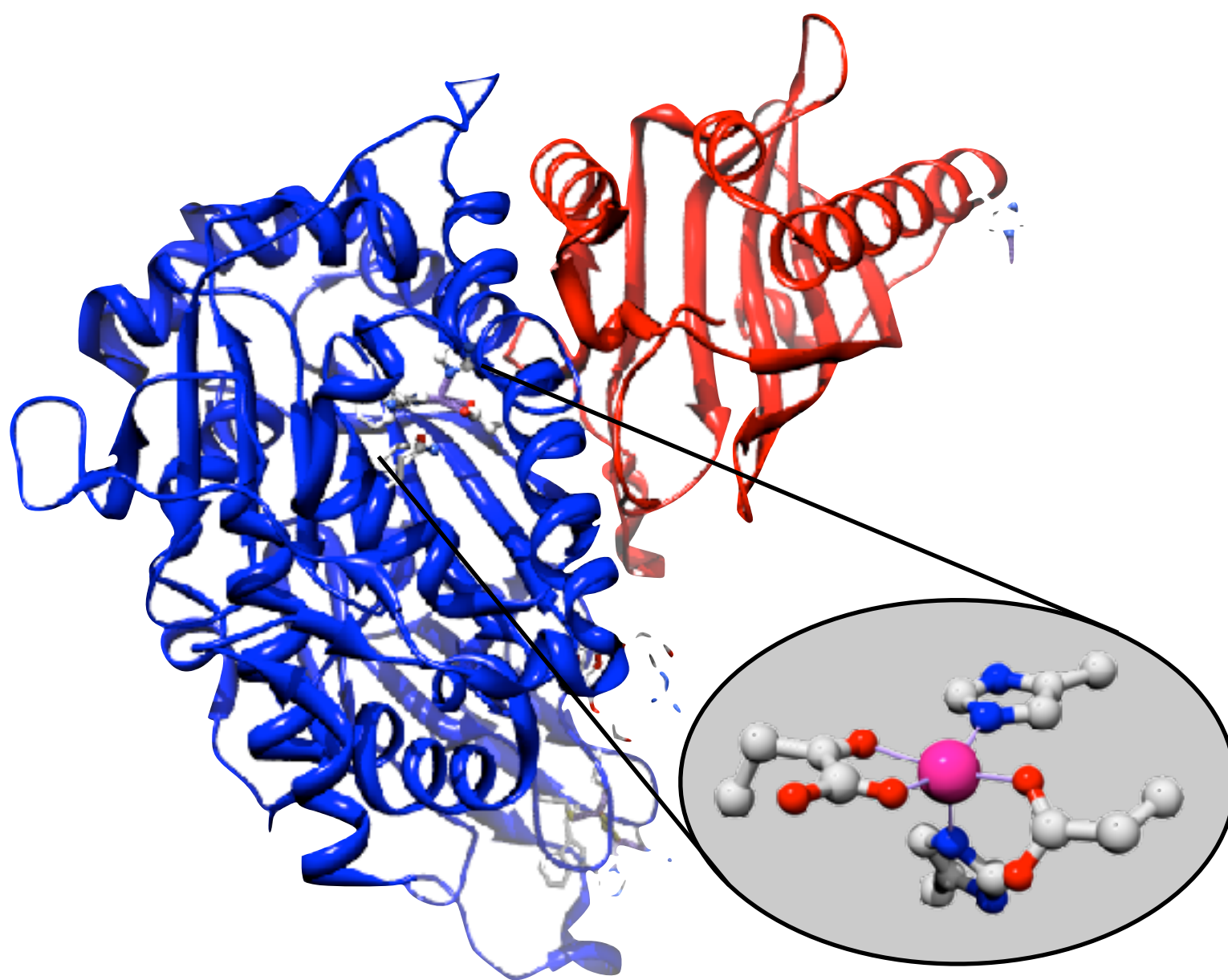


Coordination Geometry and d-d Spectra: HS-Fe(II)



Studying Enzyme Mechanisms

Rieske-Dioxygenases



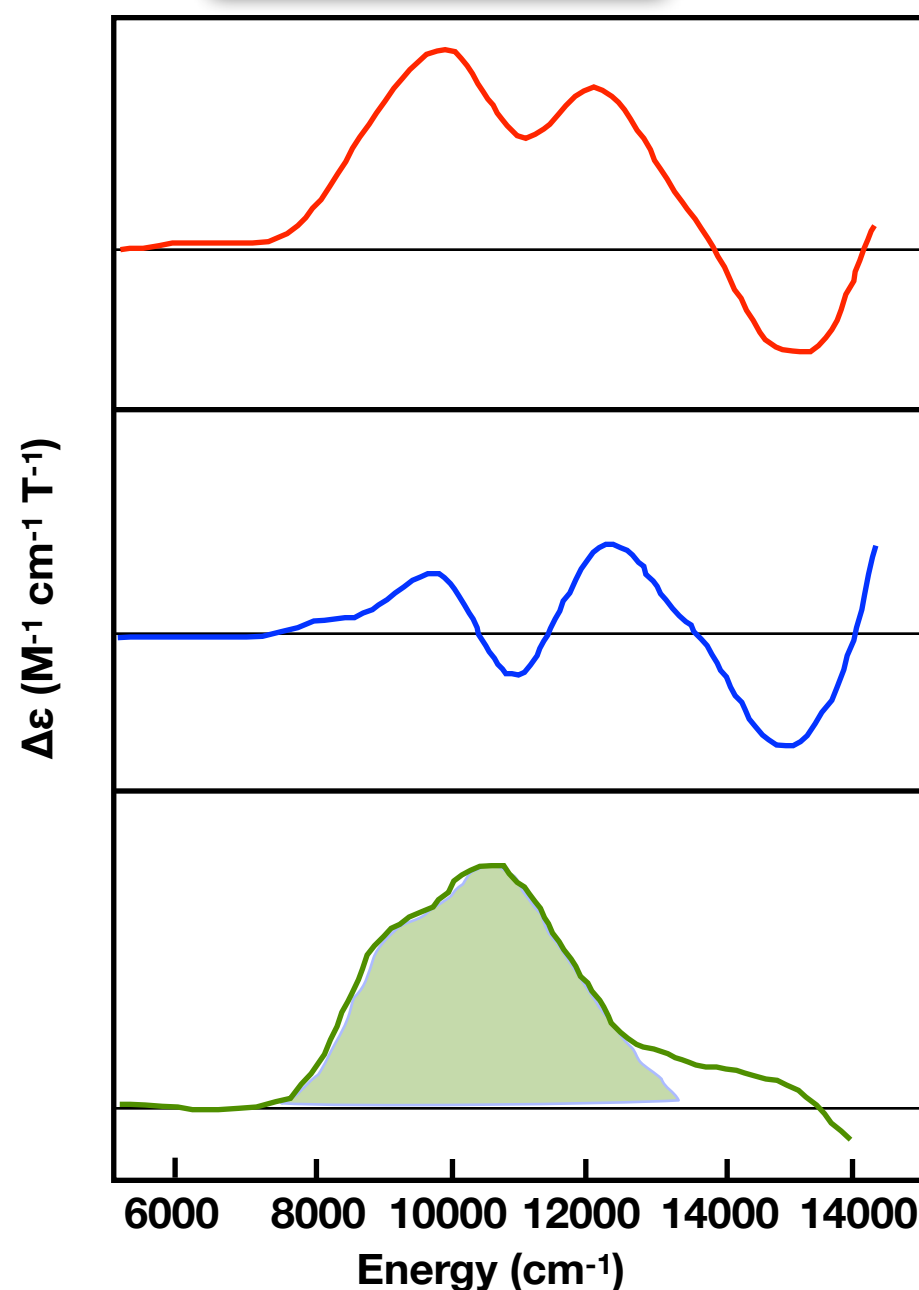
Active Site Geometry from d-d Spectra

Holoenzyme

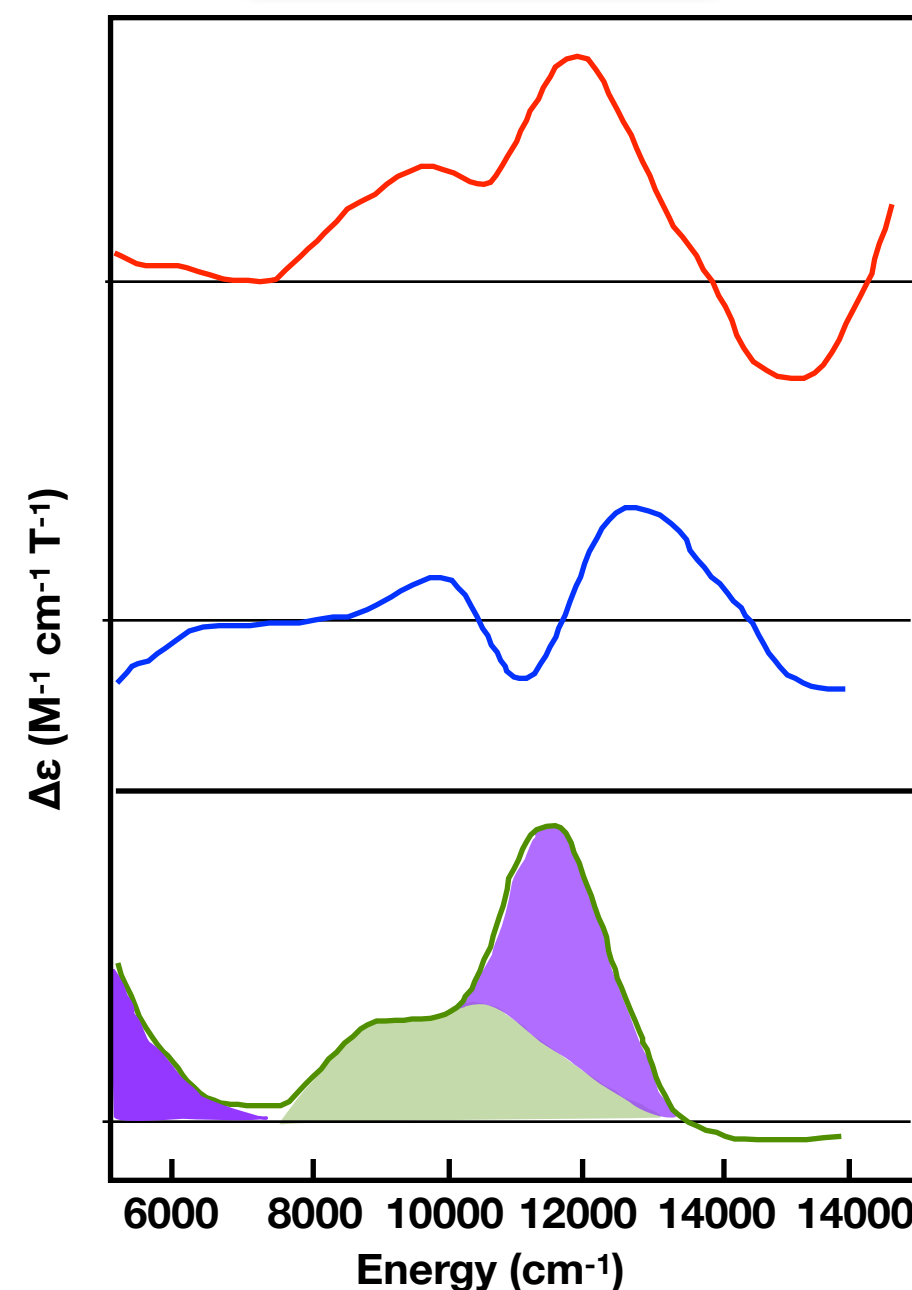
Rieske only

Difference

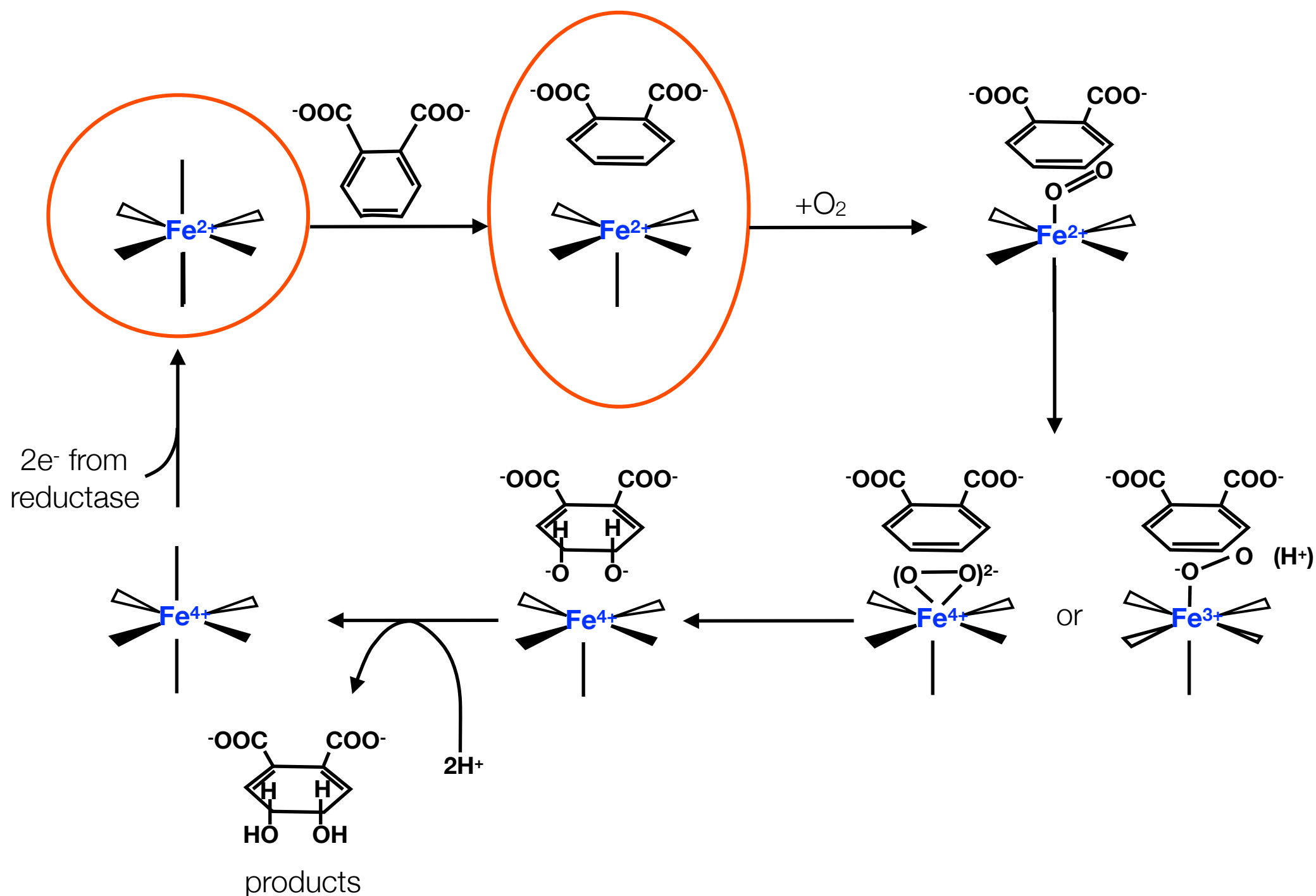
-Substrate



+Substrate



Mechanistic Ideas from Ligand Field Studies





Christian Klixbüll Jörgensen
(1931-2001)



Christian Klixbüll Jörgensen
(1931-2001)

„Personally, I do not believe much of the electrostatic
romantics, many of my colleagues talked about“

(C.K. Jörgensen, **1966** *Recent Progress in Ligand Field Theory*)

Failures of CFT...

J. OWEN
K. W. H. STEVENS

Crystal of $[\text{Ir(IV)Cl}_6]^{4-}$ d^5 $S=1/2$ exhibited an EPR spectrum with complex hyperfine (no spectrum actually shown in orig. publication)

For Ir ($I=3/2$) 4 lines are expected, but many more observed. Must be due to additional hyperfine splitting due to $I=3/2$ chlorines ligands....

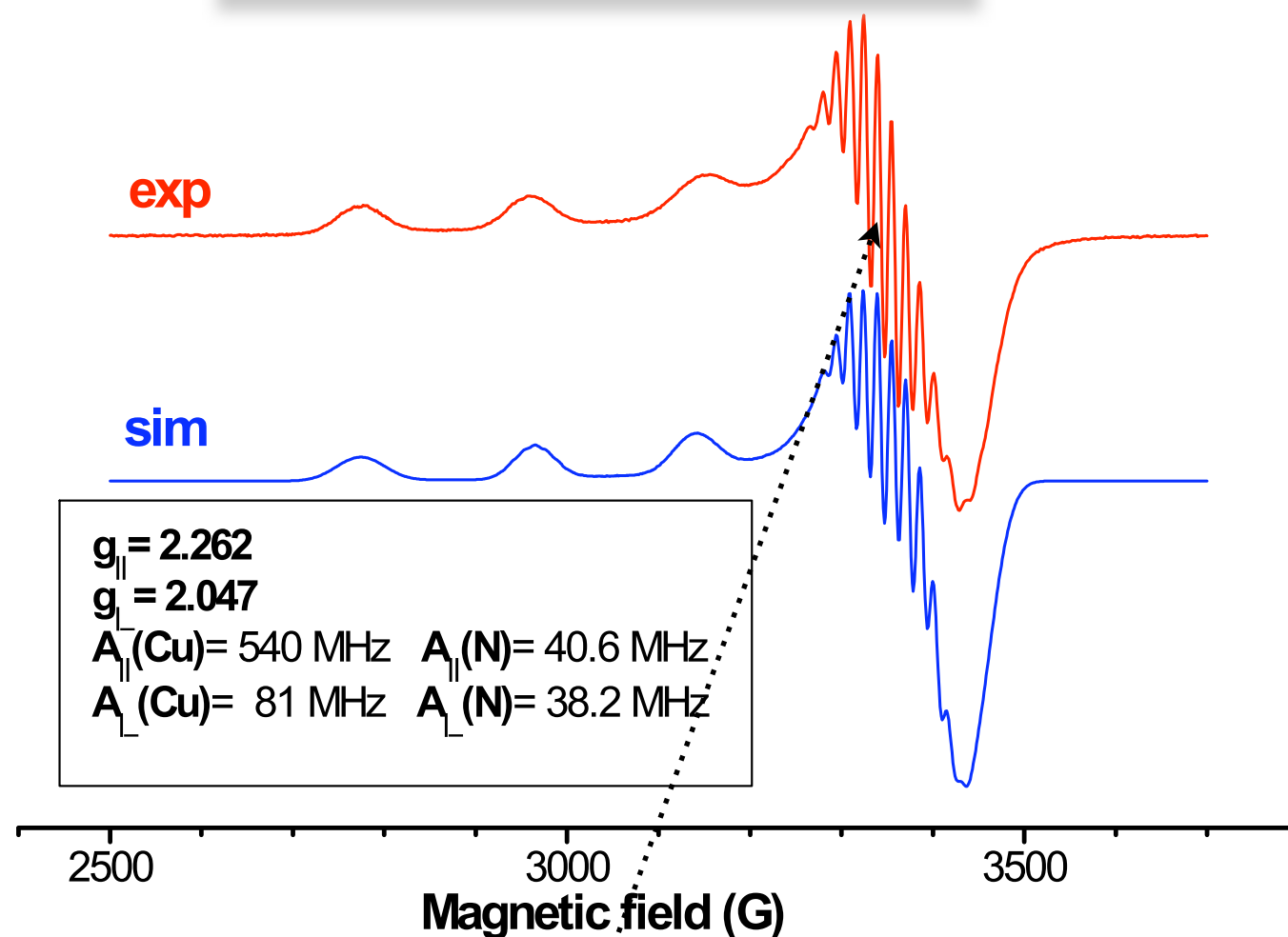
N A T U R E

May 9, 1953 VOL. 171

The experimental evidence in ammonium chloroiridate suggests that the magnetic hole is described by admixtures of $3p_\pi$ (Cl) orbitals with $5d_{xy}$, d_{yz} , d_{zx} (Ir) orbitals. This can also be pictured as a mechanism whereby for part of the time an electron is transferred from Cl^- to Ir^{4+} , cancelling the unpaired spin on the iridium, and leaving an unpaired spin on a chlorine atom.

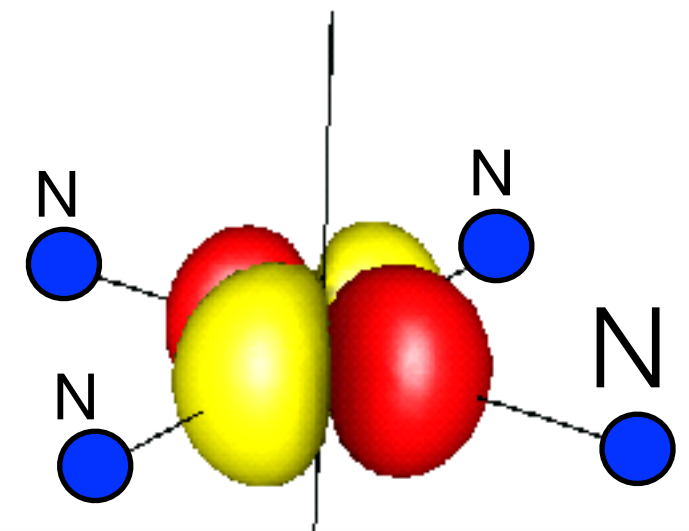
Van Vleck's theory. From the measured size of the hyperfine structure, and from the observed value, $g = 1.8$, we have tentatively estimated that the electron hole is on a particular chlorine atom for about 3 per cent of the time, and on some chlorine atom for about 18 per cent. The paramagnetic resonance method thus gives a direct and quantitative measure of π -bonding, or, crudely, of the whereabouts of the electron hole.

Experimental Proof of the Inadequacy of LFT



Clearly Observed Coupling Between
The Unpaired Electron and the Nuclear Spin
of Four ^{14}N Nitrogens ($I=1$)

Ligand Field Picture



Wavefunction of the Unpaired
Electron Exclusively Localized on
the Metal

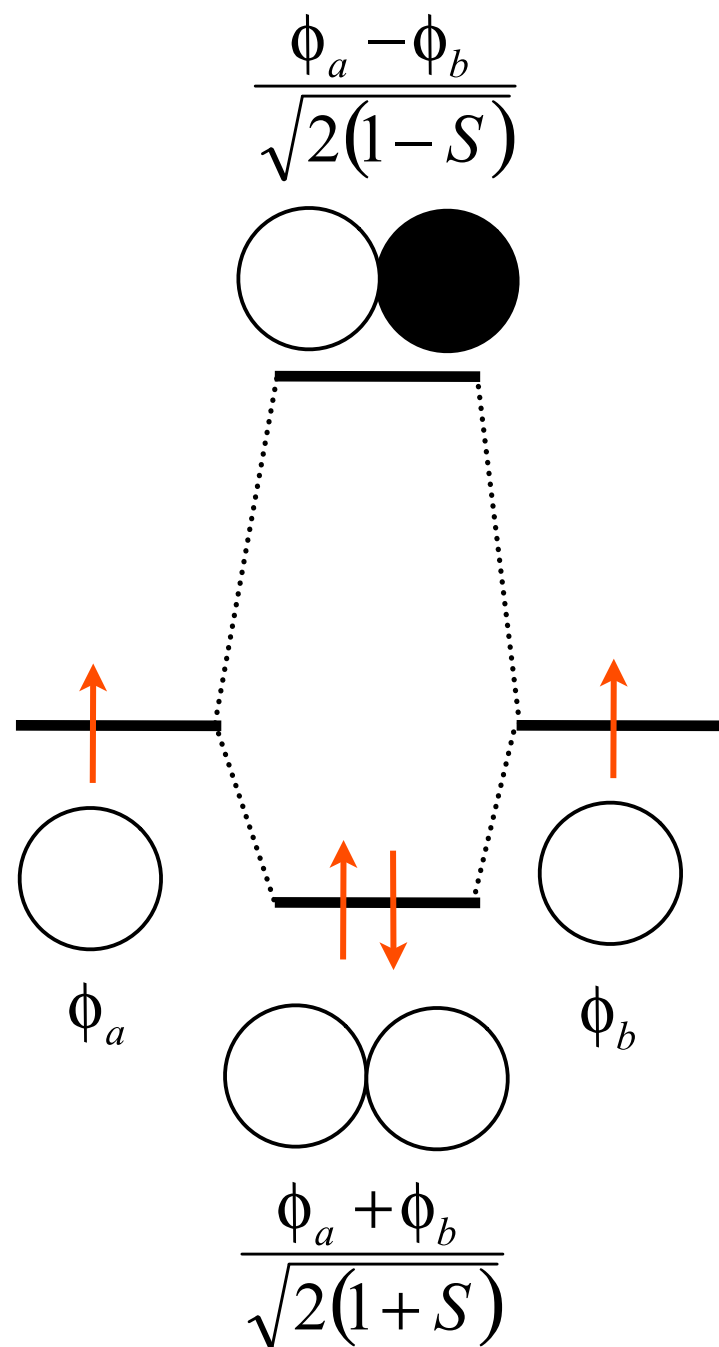


No Coupling Expected

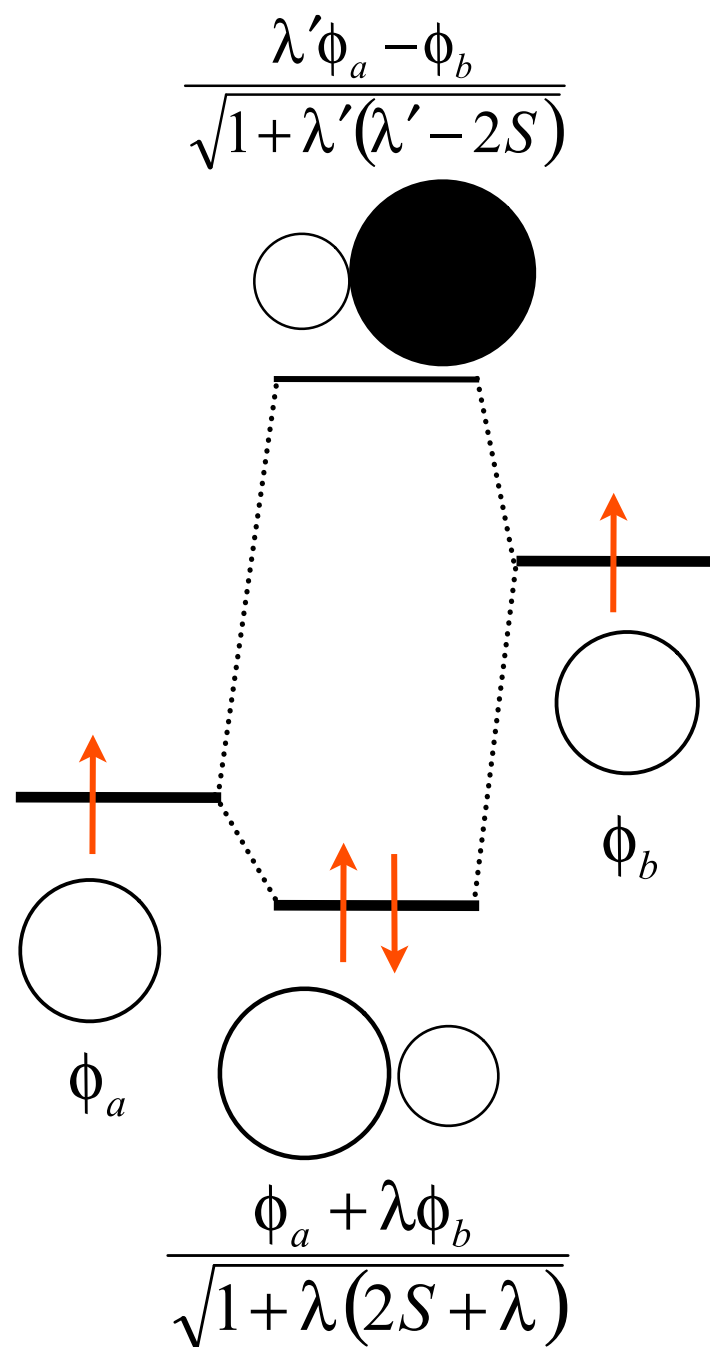


Need a Refined Theory that
Includes the Ligands Explicitly

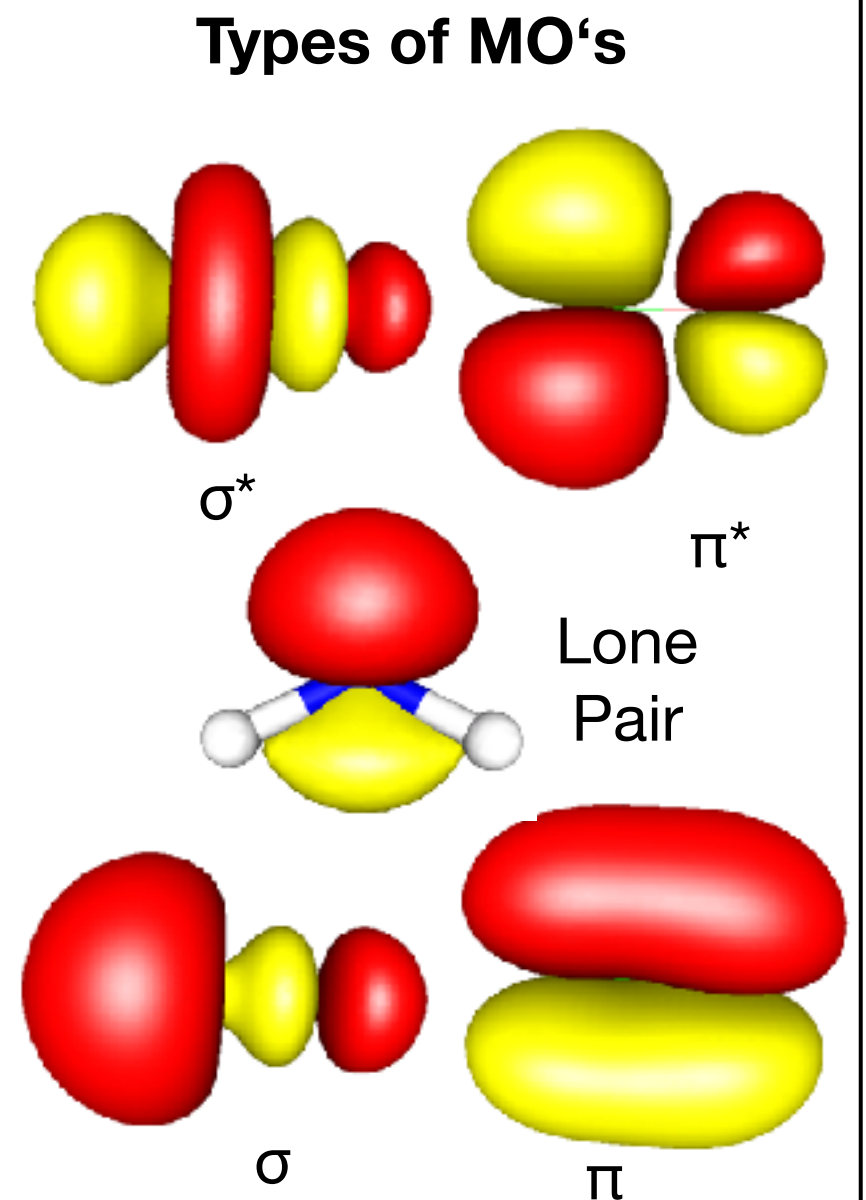
Description of Bonds in MO Theory



Homopolar Bond



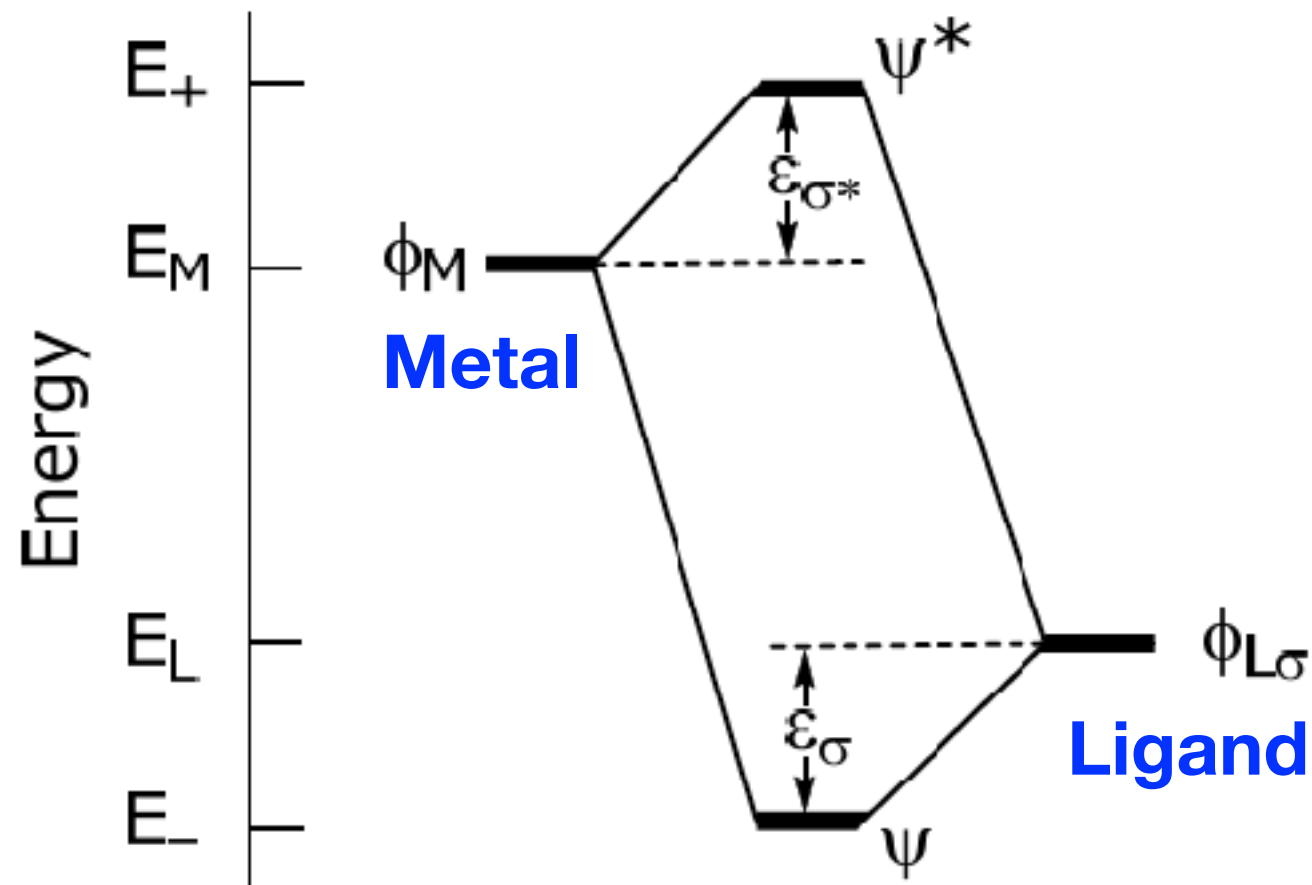
Heteropolar Bond



Bond-Order:

$$B = \frac{1}{2} (N_B - N_A)$$

Rules for constructing MO diagrams...



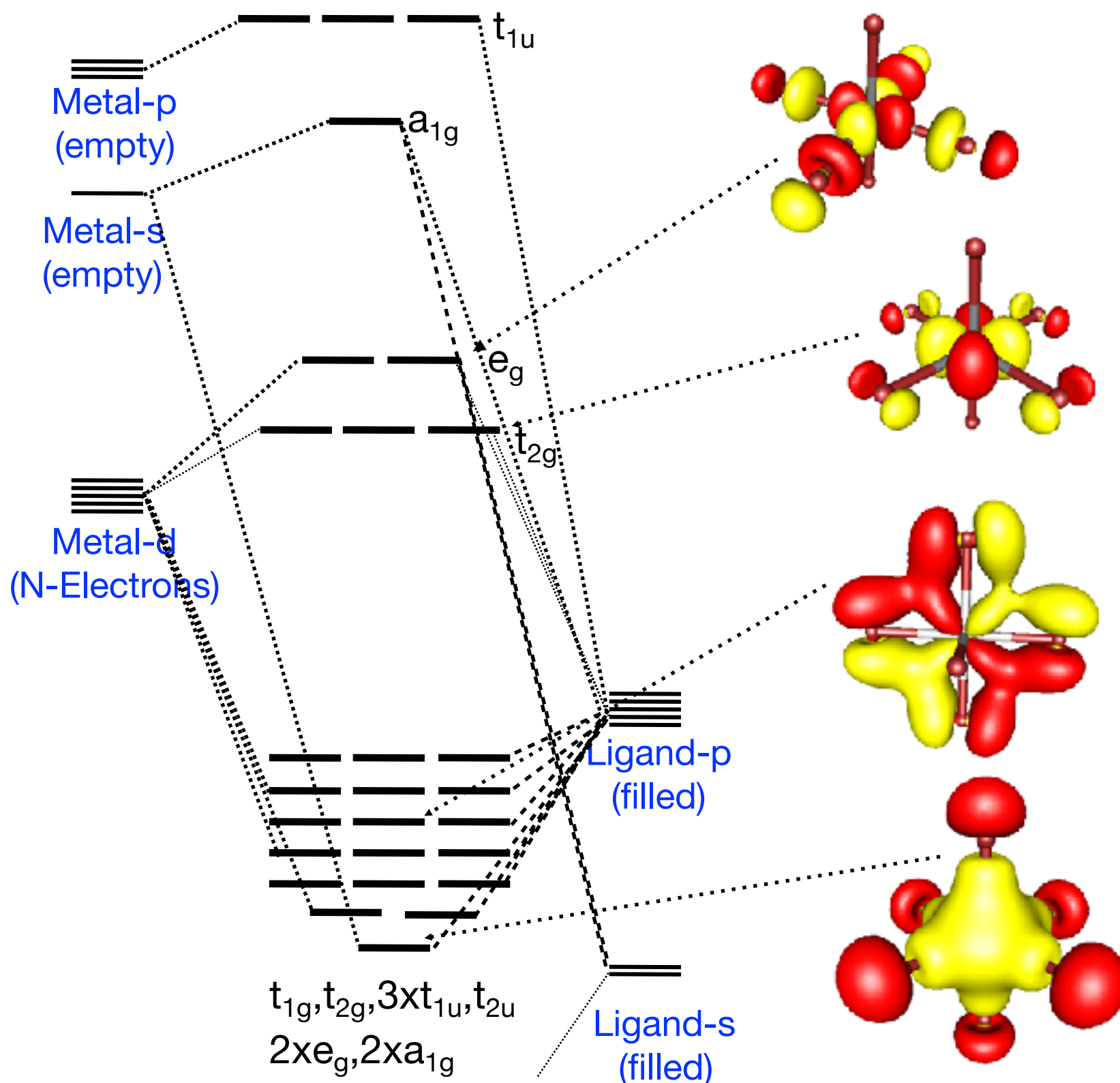
$$\epsilon_{\sigma^*} \approx \frac{\beta^2}{\Delta E_{ML}} - \frac{\beta S_{ML}}{1 - S_{ML}^2}$$

$$\epsilon_\sigma \approx -\frac{\beta^2}{\Delta E_{ML}}$$

$$\beta \propto k S_{ML} \quad \text{„resonance integral“}$$

- ✓ Metal-ligand orbital mixing is proportional to the **overlap** of the metal and ligand orbital (S_{ML})
- ✓ Metal-ligand orbital mixing is inversely proportional to **energy difference** of mixing orbitals (i.e. ΔE_{ML})

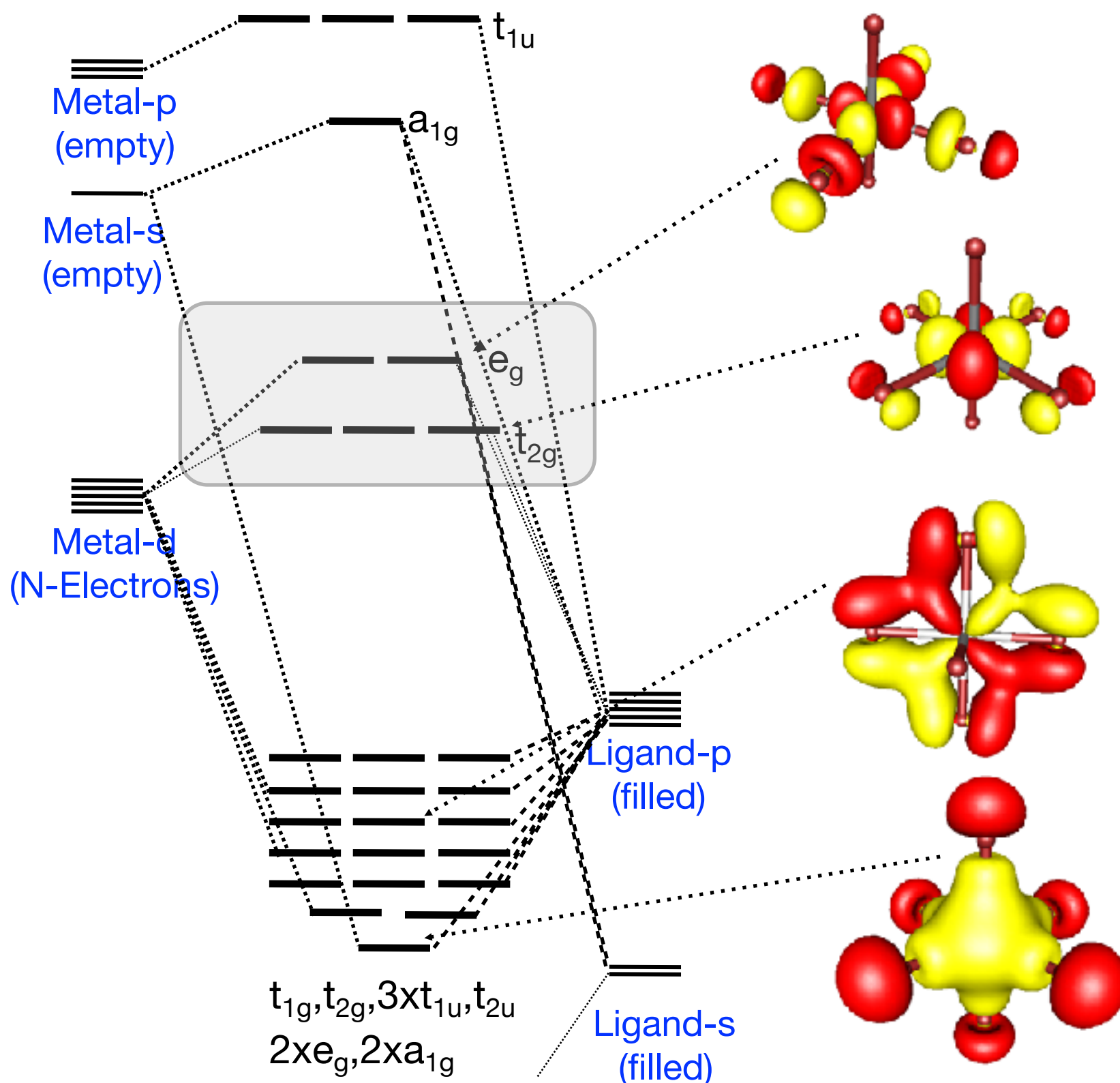
MO Theory of ML_6 Complexes



Key Points:

- ▶ Filled ligand orbitals are **lower in energy** than metal d-orbitals
- ▶ **The orbitals that are treated in LFT correspond to the anti-bonding metal-based orbitals in MO Theory**
- ▶ Through bonding some **electron density** is transferred from the ligand to the metal
- ▶ The extent to which this takes place defines the **covalency** of the M-L bond

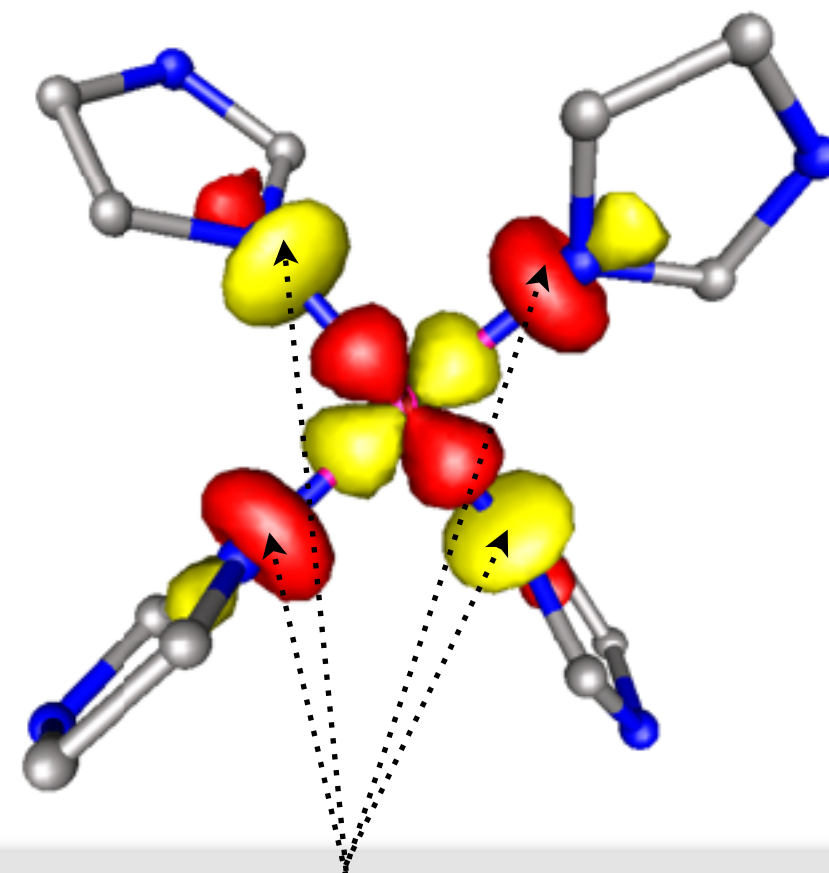
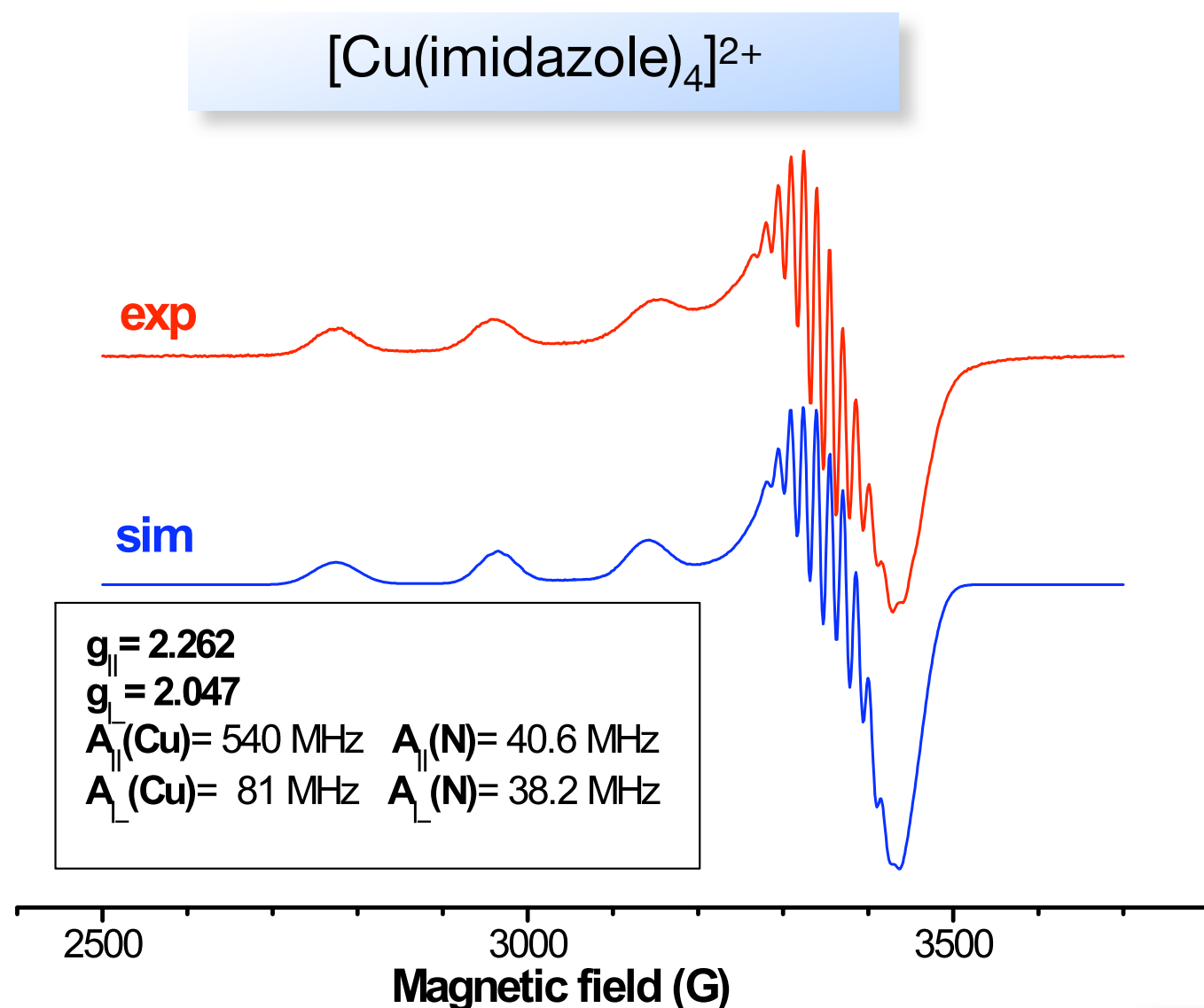
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MO Theory and Covalency



The Unpaired Electron is Partly
Delocalized Onto the Ligands

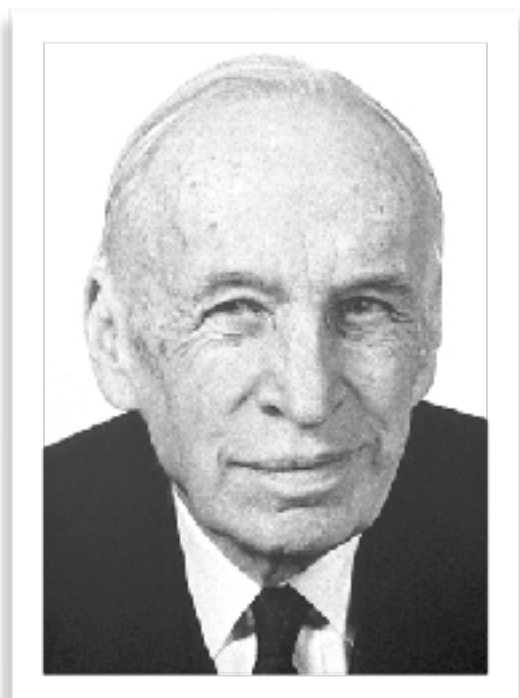


Magnetic Coupling Expected



1. Theory Accounts for Experimental Facts
2. Can Make Semi-Quantitative Estimate of the Ligand Character in the SOMO

Ligand Field Theory (LFT)



John H. van Vleck
(1935)

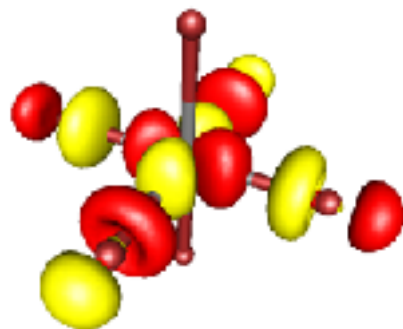
- *LFT combines principles laid out in CFT with molecular orbital theory
- * Accounts for the *nature* of the ligand donor properties
- *Relies on **symmetry** and **covalency** to form sigma, pi and delta bonds

“In 1935 I published a paper in which I amplified and generalized...the primitive crystal field theory... I showed that Bethe’s grouping of energy levels according to symmetry type was still valid even if one allowed the electrons in the unclosed shells to wander away sometimes from the central paramagnetic ion and take a look at the diamagnetic atoms clustered around it.” - J.H. van Vleck’s, Nobel Prize Address (1977)

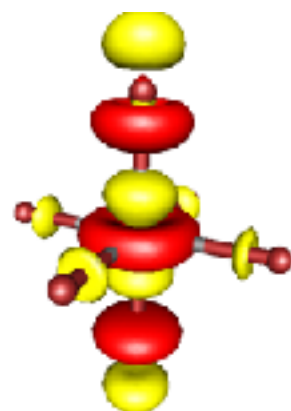
π -Bonding and π -Backbonding



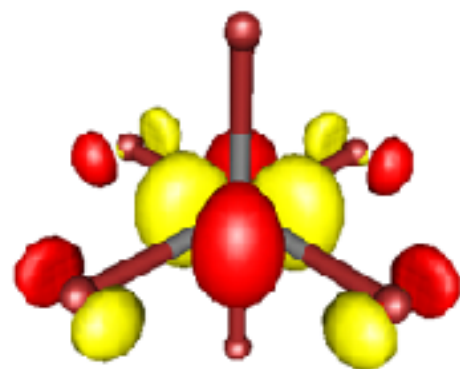
$d_{x^2-y^2}$



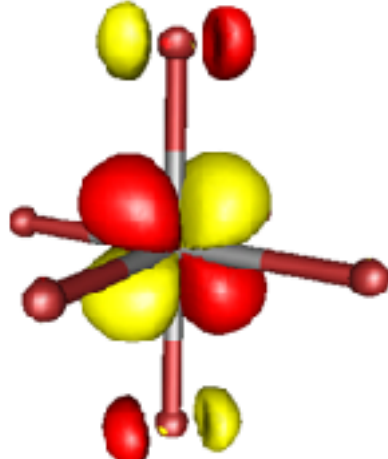
d_{z^2}



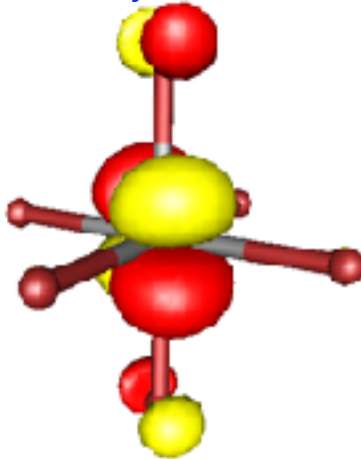
d_{xy}



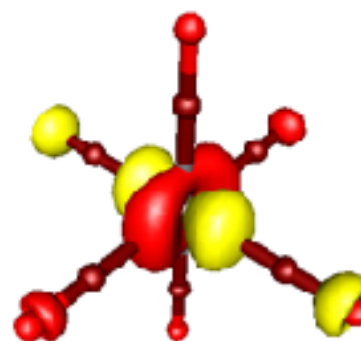
d_{xz}



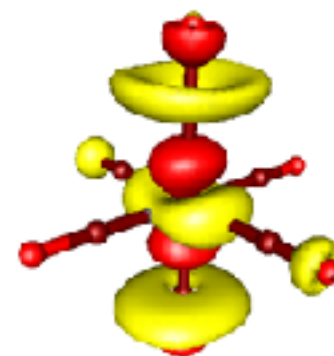
d_{yz}



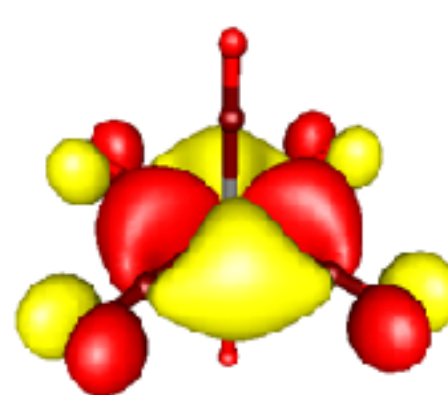
$d_{x^2-y^2}$



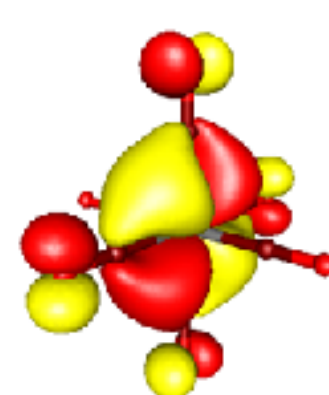
d_{z^2}



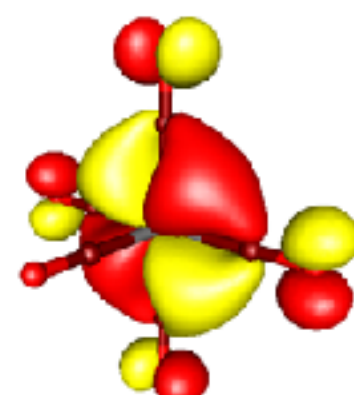
d_{xy}



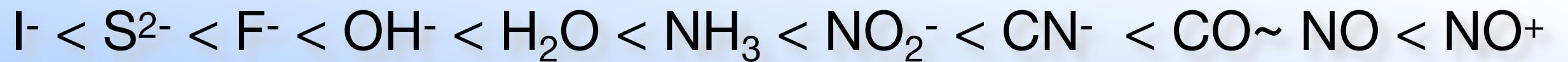
d_{xz}



d_{yz}

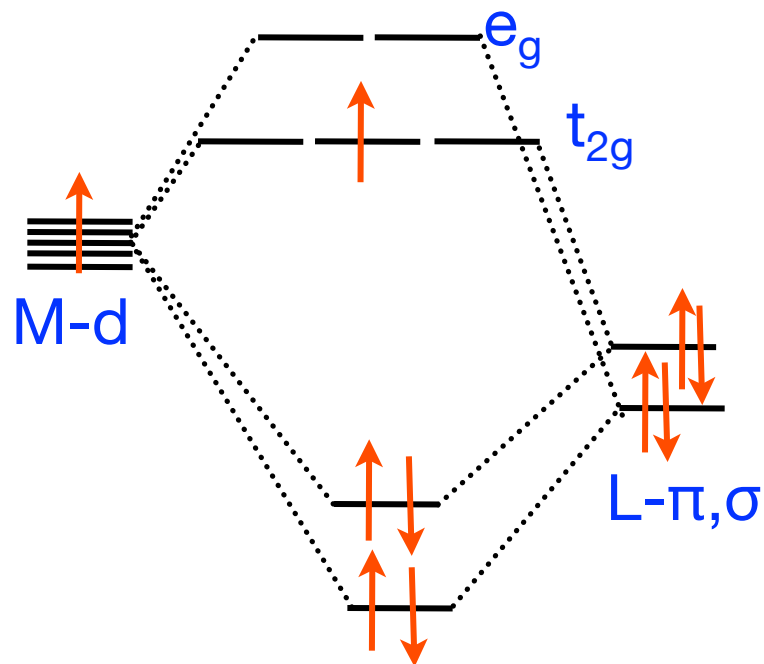


Interpretation of the Spectrochemical Series



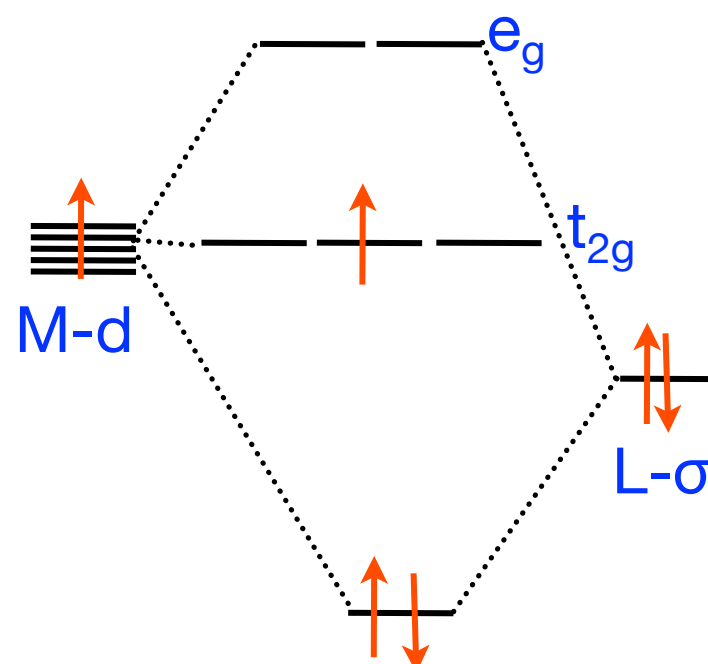
Δ SMALL

π -DONOR

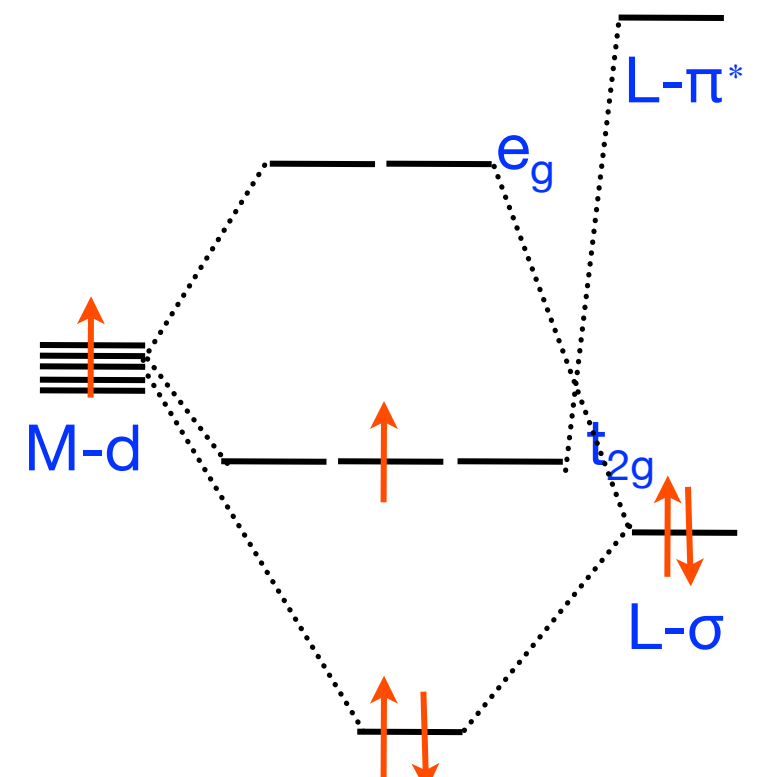


Δ LARGE

π -'NEUTRAL'



π -ACCEPTOR



Oxidation States, Calculations and Chemical Thinking

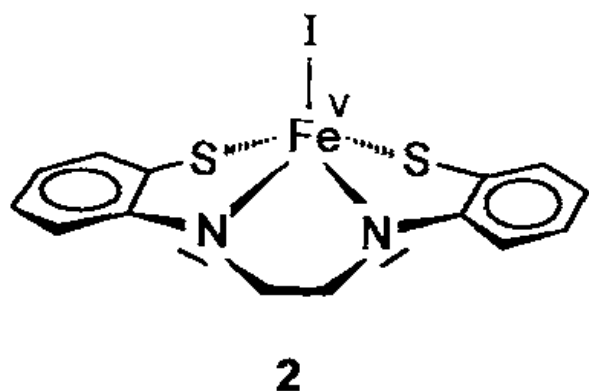
Physical versus Formal Oxidation States

The **formal oxidation state** of a metal ion in a complex is the d^N configuration that arises upon dissociating all ligands in their closed shell „standard“ states taking into account the total charge of the complex

The **physical oxidation state** of a metal ion in a complex is the d^N configuration that arises from an analysis of its electronic structure by means of spectroscopic measurements and molecular orbital calculations

Chaudhuri, P.; Verani, C.N.; Bill, E.; Bothe, E.; Weyhermüller, T.; Wieghardt, K. *J. Am. Chem. Soc.*, **2001**, 123, 2213 „*The Art of Establishing Physical Oxidation States in Transition-Metal Complexes Containing Radical Ligands* “. However, The concept goes back to CK Jörgensen

- ★ The two often coincide but may well be different! In the example before, the formal oxidation state is Ni(IV) but the physical oxidation state is Ni(II)

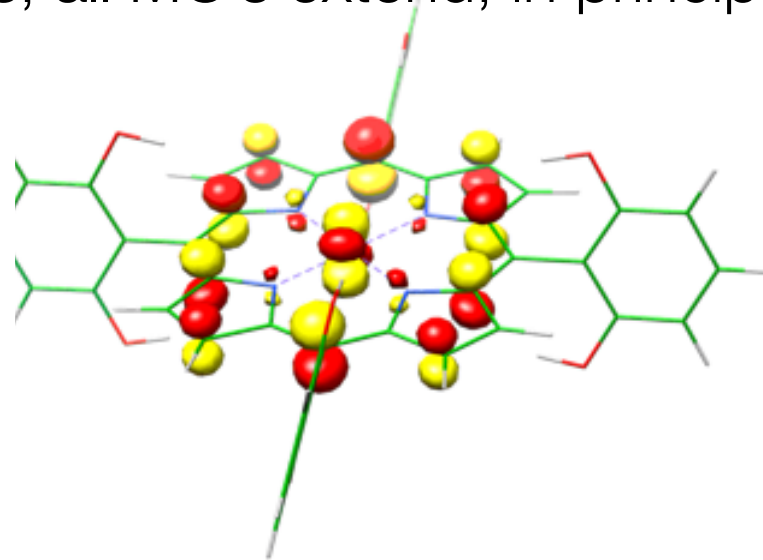


This complex has first been described by its formal oxidation state of Fe(V) but has a physical oxidation of Fe(III)

Chlopek, C. *et al. Chem. Eur. J.* **2007**, 13, 8390

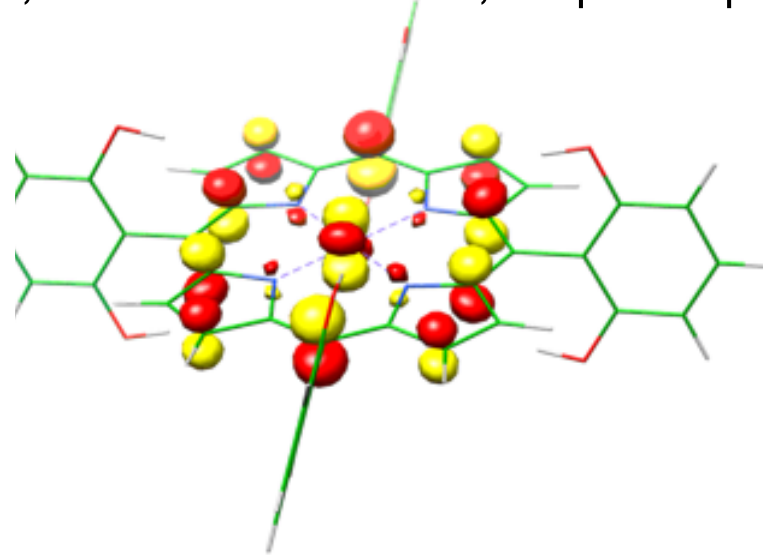
Oxidation States and MO Theory

In Molecular Orbital theories, all MO's extend, in principle over the whole molecule



Oxidation States and MO Theory

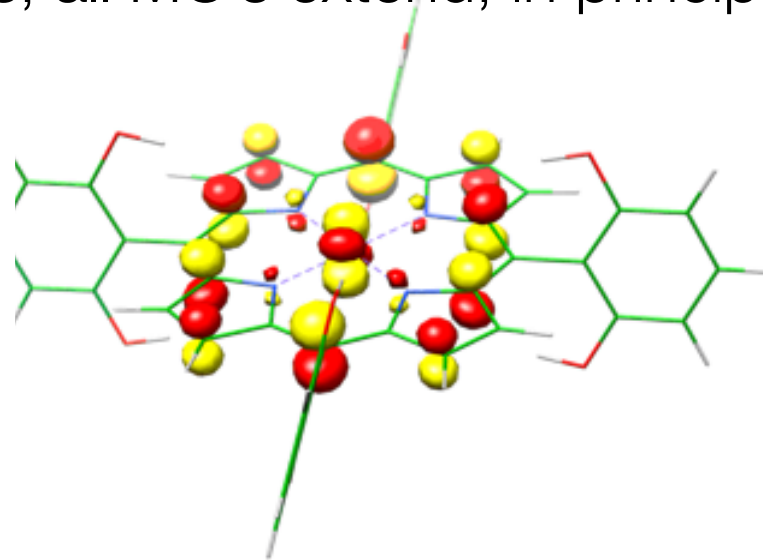
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The concept of the oxidation state rests on the idea that one can „count“ d-electrons

Oxidation States and MO Theory

In Molecular Orbital theories, all MO's extend, in principle over the whole molecule



The concept of the oxidation state rests on the idea that one can „count“ d-electrons

So, what does an oxidation state mean in a MO context?



(Pragmatic) Assignment of Oxidation States

- ▶ Analyze the occupied orbitals of the compound and determine the fractional d-character of the orbitals
(In **ORCA** this is possible via the `NormalPrint`, `Localize` or `UNO` keywords)

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- ▶ Count the number of electrons in such metal based orbitals. This gives you the **d^N configuration**

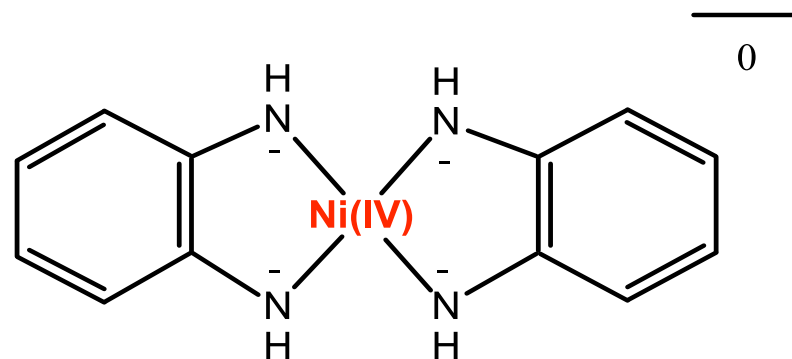
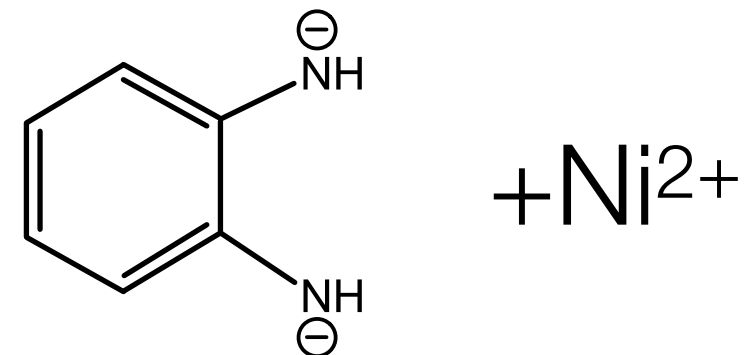
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- ▶ Count the number of electrons in such metal based orbitals. This gives you the **d^N configuration**
- ▶ The **local spin state** on the metal follows from the singly occupied metal-based orbitals

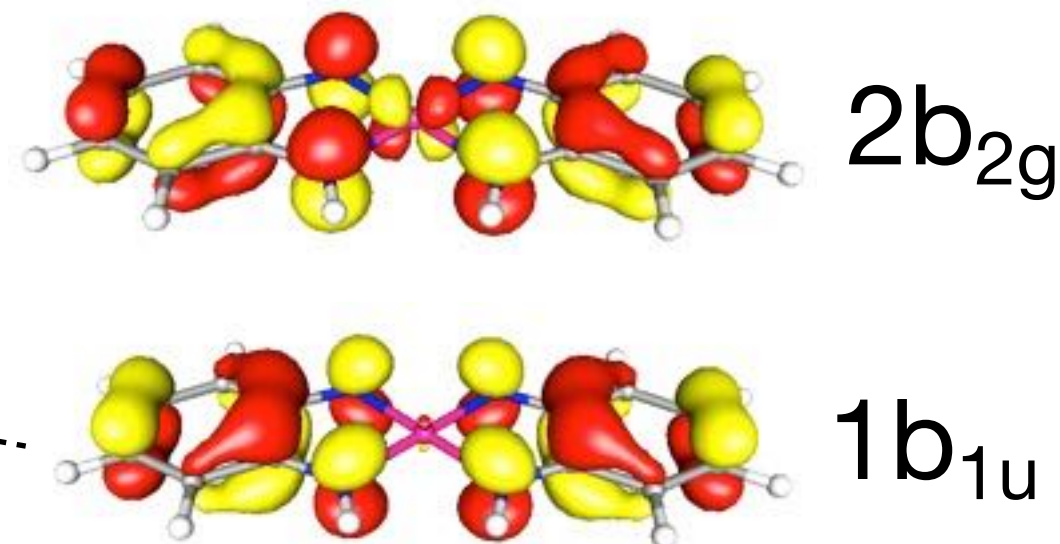
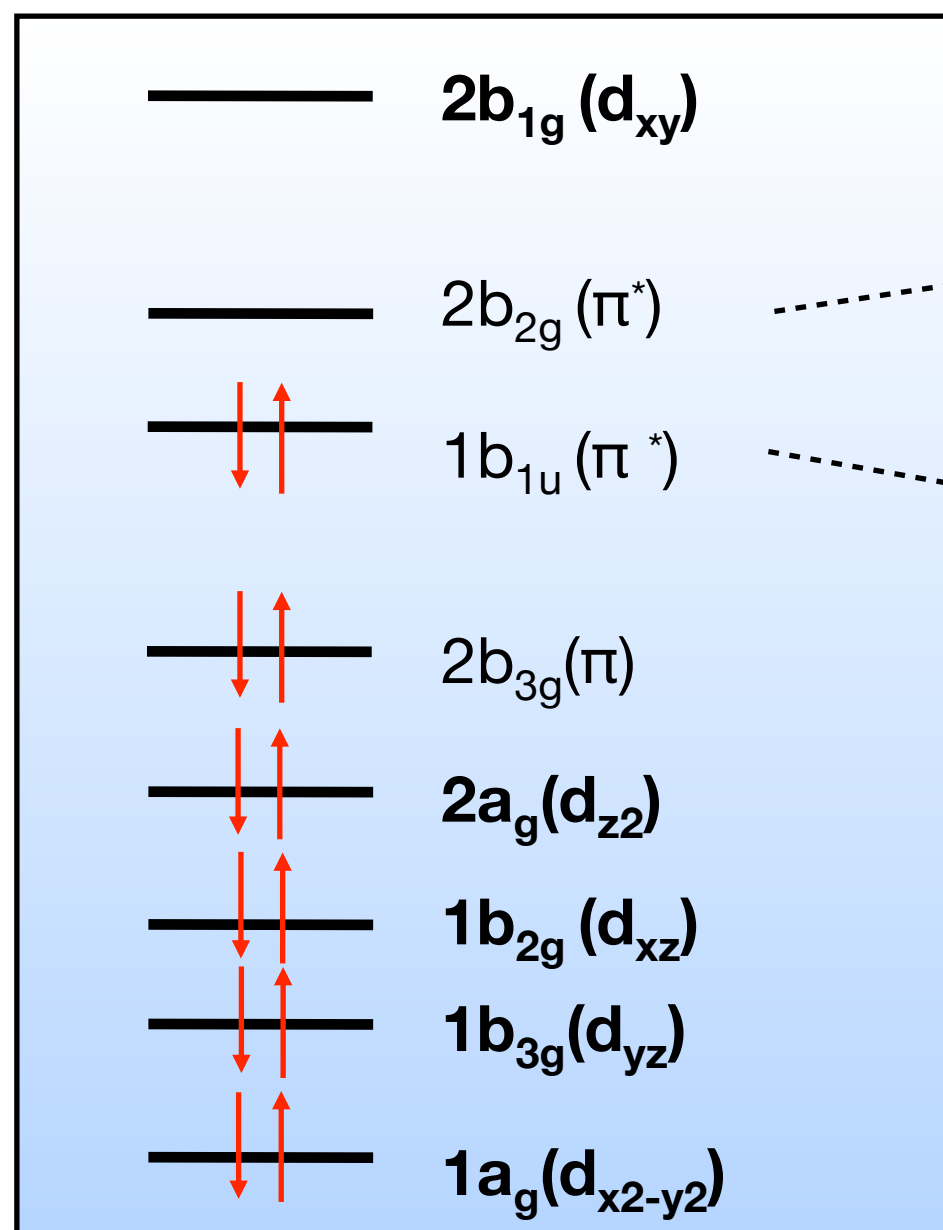
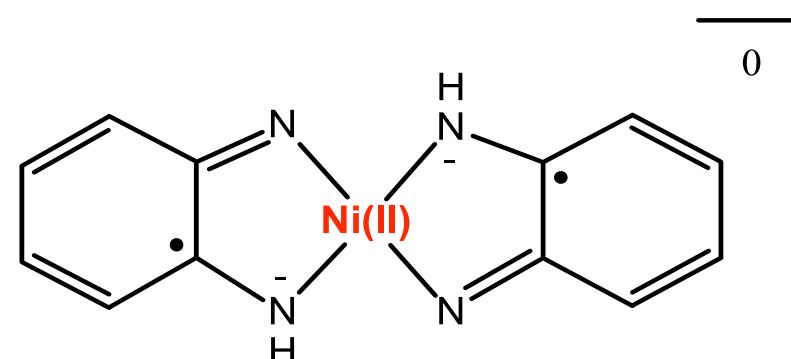
(Pragmatic) Assignment of Oxidation States

- ▶ Analyze the occupied orbitals of the compound and determine the fractional d-character of the orbitals
(In **ORCA** this is possible via the `NormalPrint`, `Localize` or `UNO` keywords)
- ▶ Orbitals that are centered more than, **say**, 70% on the metal are counted as pure metal d-orbitals
- ▶ Count the number of electrons in such metal based orbitals. This gives you the **d^N configuration**
- ▶ The **local spin state** on the metal follows from the singly occupied metal-based orbitals
- ▶ This fails, if there are some orbitals that are heavily shared with the ligands (metal character < 70%). In this case the oxidation state is ambiguous.
(Typically, experiments are ambiguous as well in these cases)

An Example



OR

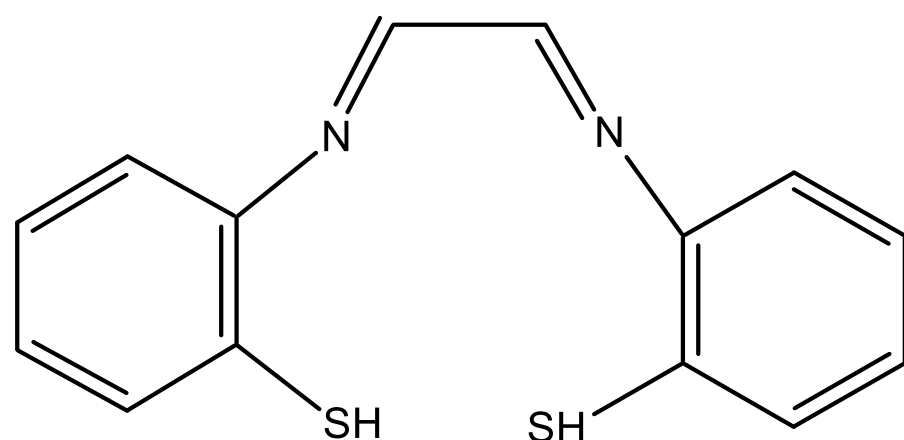


4 doubly occupied and one empty metal based orbitals

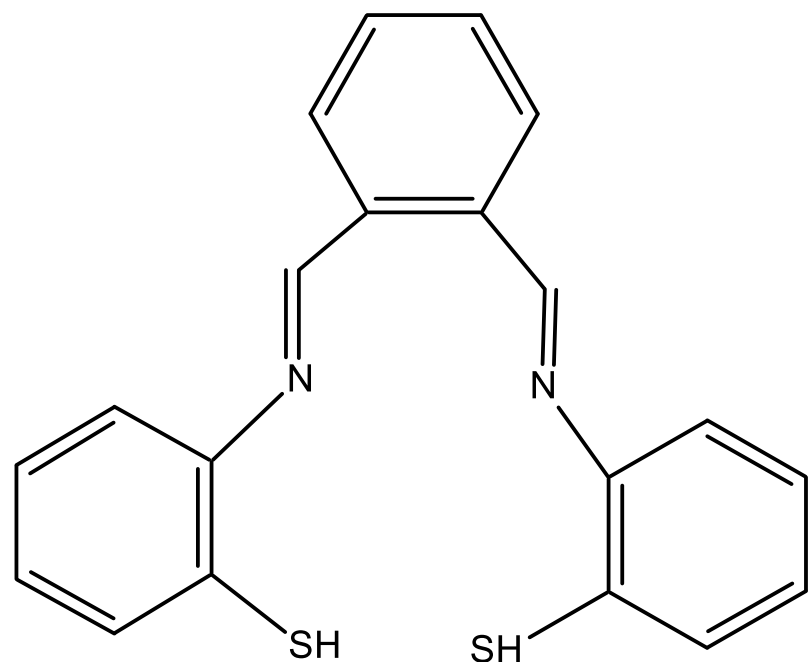
- ▶ Ni(II) oxidation state
- ▶ Ligand must be oxidized and has diradical character

This is evidenced by close lying HOMO and LUMO that are symmetric and antisymmetric combinations

An Example for an ,Exciting‘ Oxidation State



gma



PhBMA

[Zn(gma)]
S=0 (Holm,Gray)

[Zn(gma)]⁻
S=1/2 (Holm,Gray)

[Ni(gma)]
S=0 (Holm,Gray)

[Ni(gma)]⁻
S=1/2 (Holm,Gray)

[Fe(gma)(py)]

[Fe(gma)(PR₃)]

[Fe(gma)(CN)]⁻

S=1
(Strähle, Sellmann,
Wieghardt)

[Fe(PhBMA)]

[Fe(gma)(PR₃)₂]

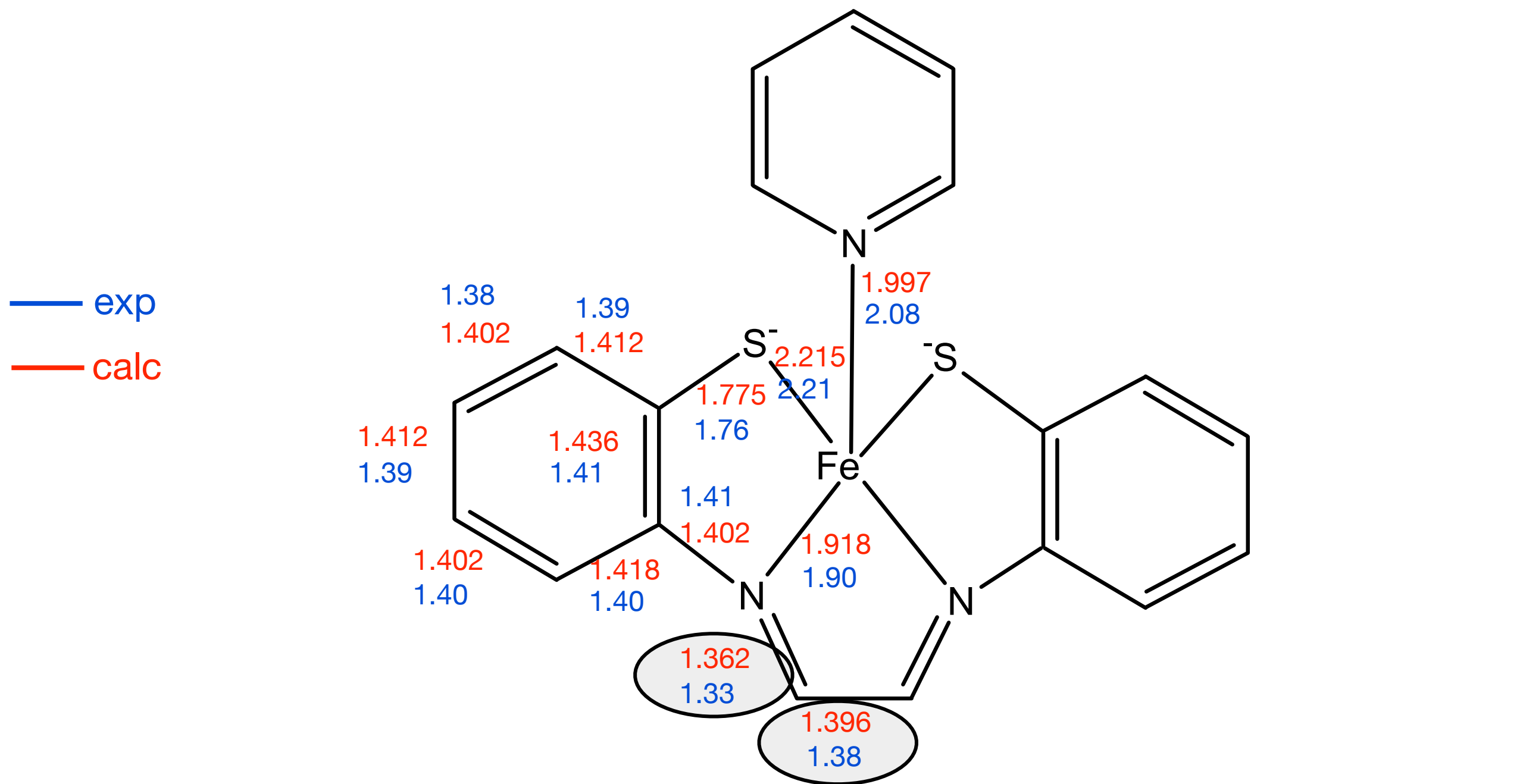
[Fe(gma)(py)]⁺

S=1
S=0 (Wieghardt)

S=1/2

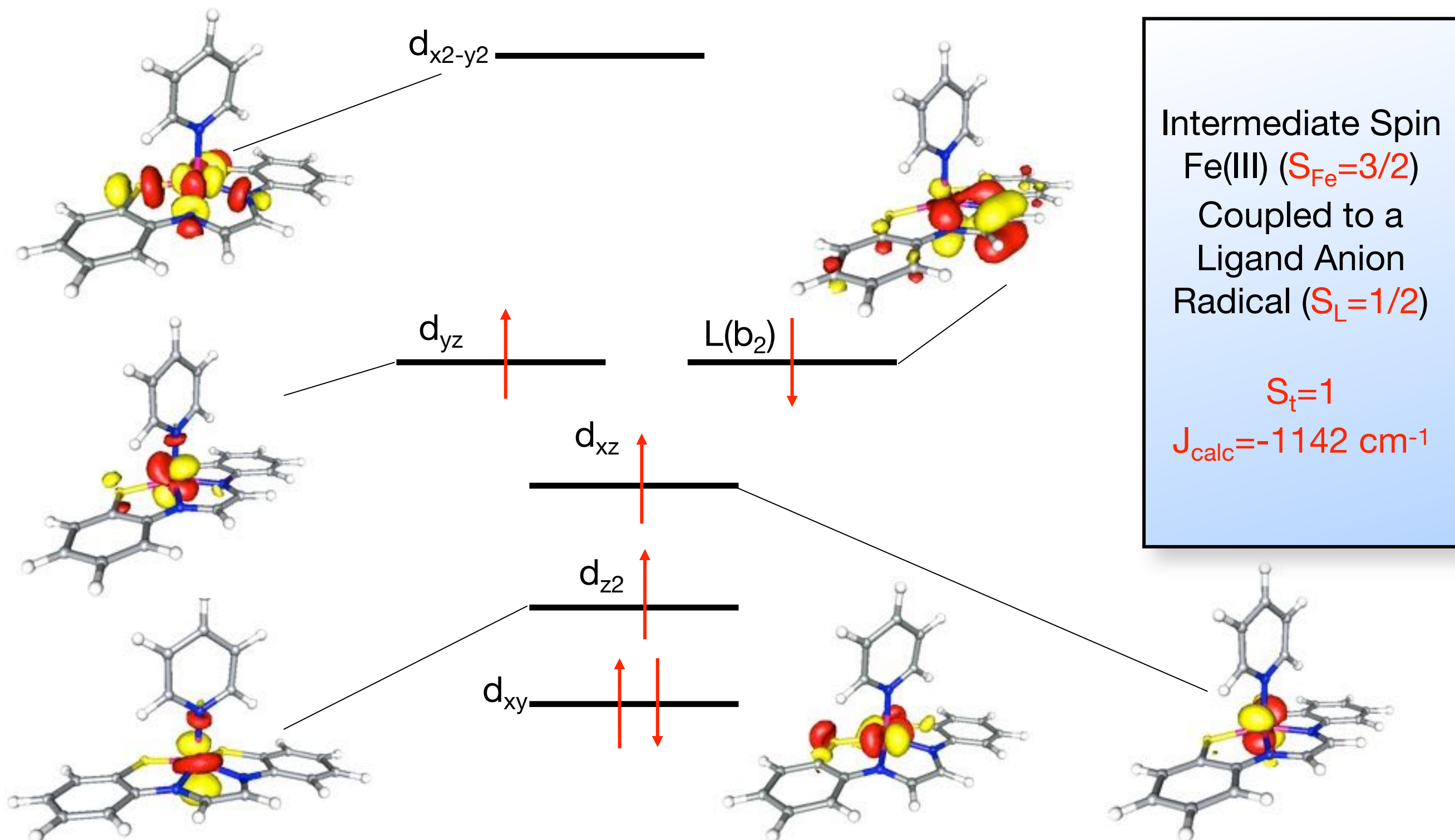
All described as ordinary
metal(II) chelates

Optimized Structure of [Fe(L)(py)]



		δ_{MB}	ΔE_Q	η
[Fe(L)(py)]	Calc.	0.205	+2.371	0.084
	Exp.	0.270	+2.33	small

Electronic Structure of [Fe(gma)(py)]

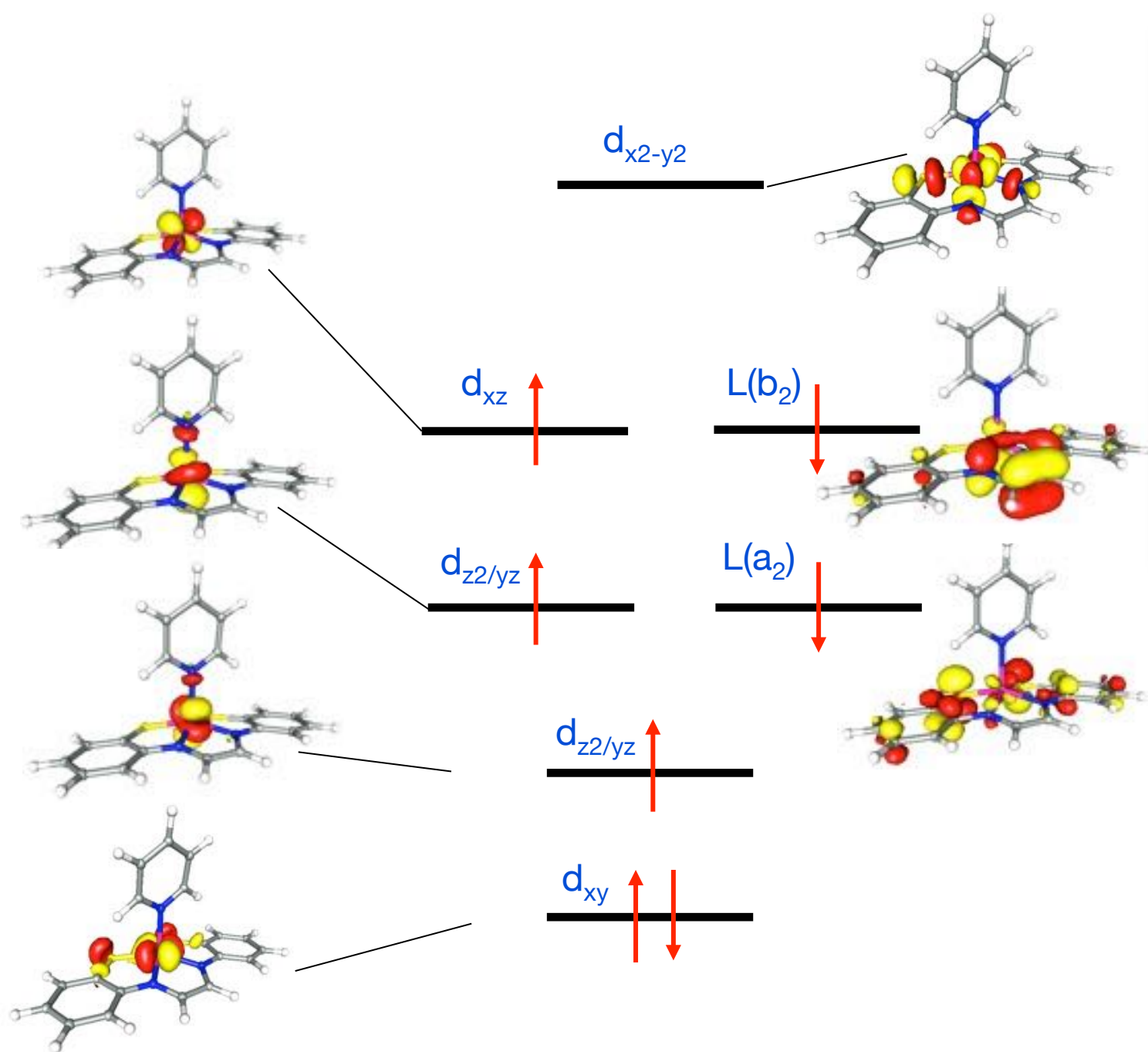


Electronic Structure of $[\text{Fe}(\text{gma})(\text{py})]^+$

Observation: Oxidation leaves the Mössbauer parameters invariant

- Dilemma:
- ★ $[\text{Fe}(\text{gma})(\text{py})]$ has a spin of 1
 - ▶ Intermediate Spin Fe(III)/Ligand Radical
 - ★ Taking the Electron out of the ligand LUMO
 - ▶ Intermediate Spin Fe(III)/Closed Shell Ligand
 - ▶ $S=3/2$ Expected but $S=1/2$ Observed
 - ★ Taking the Electron out of the Ligand but Changing the Spin at Fe
 - ▶ Would have been Detected in MB Experiments
 - ★ Taking the Electron from the Iron gives Fe(IV)
 - ▶ Would have been Detected in MB Experiments

Electronic Structure of $[\text{Fe}(\text{gma})(\text{py})]^+$

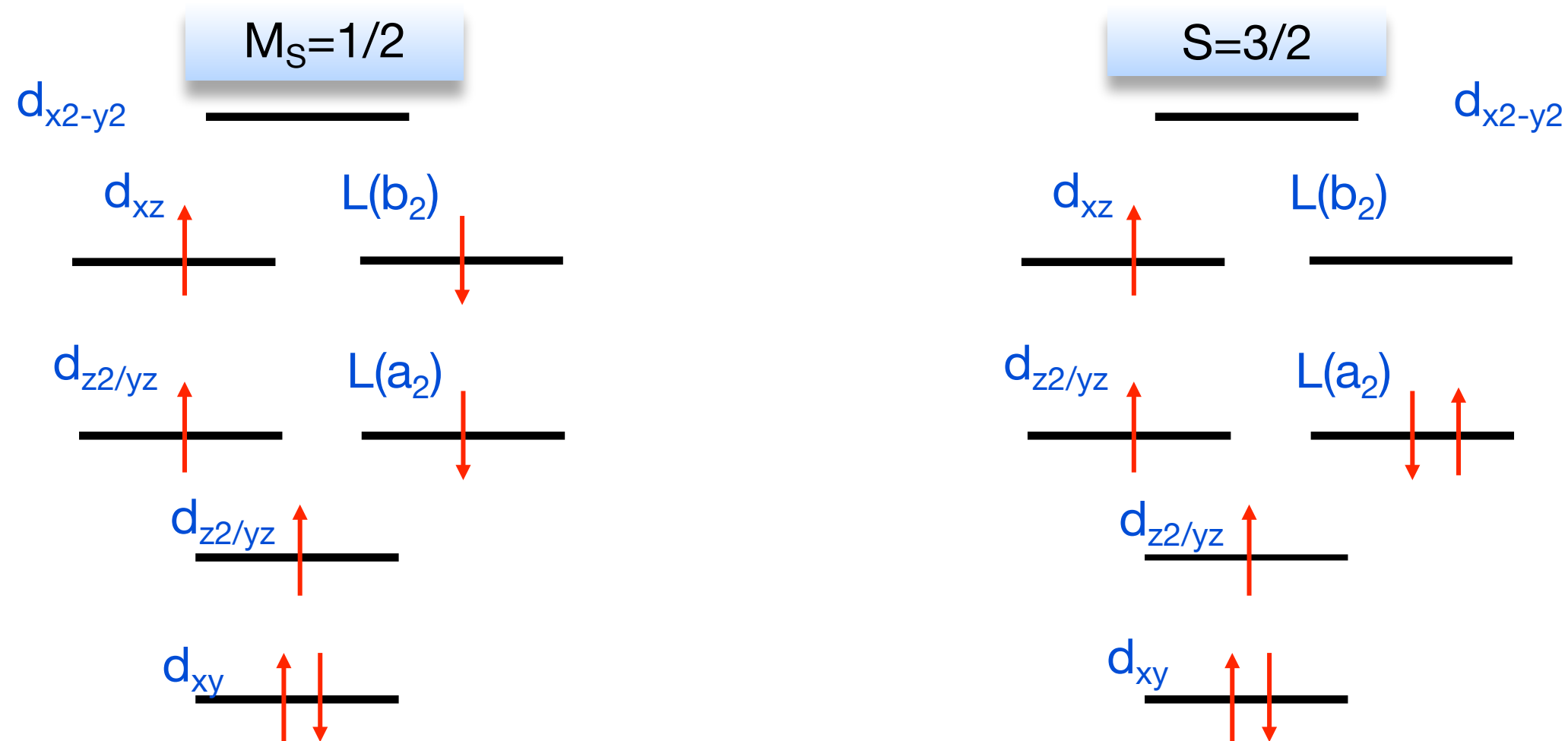


	Mössbauer Parameters		
	δ (mm/s)	ΔE_Q (mm/s)	η
Neutral	0.205	2.371	0.084
Cation	0.246	2.430	0.023
Exp.	(0.27)	(2.33)	?

Intermediate Spin Fe(III) ($S_{\text{Fe}}=3/2$)
Coupled to a Ligand Triplet
State ($S_L=1$)

$$S_t=1/2 \quad J_{\text{calc}}=-845 \text{ cm}^{-1}$$

Electronic Structure of $[\text{Fe}(\text{gma})(\text{py})]^+$



$$E(M_S = 1/2) = -0.6 \text{ kcal/mol}$$

$$E(S = 1/2) \sim -5.4 \text{ kcal/mol}$$

$$E(S = 3/2) = 0.0 \text{ kcal/mol}$$

Large Exchange Coupling Drives Ligand Triplet State Coordination
Basis for a „Metal Field Theory“ rather than „Ligand Field Theory“

(Guihery, N.; Robert, V.; FN *J. Phys. Chem.*, 2008, 112, 12975)

Making Ligand Field Theory quantitative

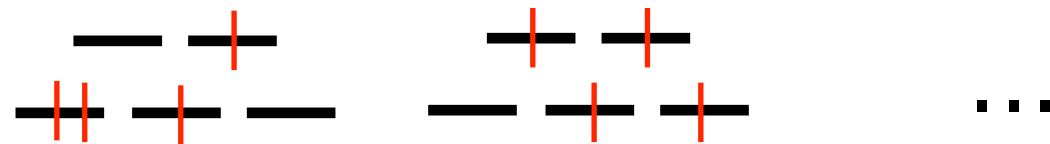
„Complete“ Ligand Field Theory

A complete LFT calculation in the strong field scheme proceeds as:

1. Build the one-electron matrix:

$$h_{\mu\nu} = f(e_{\sigma}^L, e_{\pi}^L) \quad \mu, \nu = d - \text{orbitals}$$

2. Build all configurations



3. Build all Slater determinants: $\Phi_I(\mathbf{x}_1, \dots, \mathbf{x}_N)$ $\Phi_1 = | d_{xy}^{\alpha} d_{xz}^{\alpha} d_{z^2}^{\beta} | \dots$ $\Phi_{65} = | d_{xz}^{\beta} d_{yz}^{\alpha} d_{x^2-y^2}^{\alpha} | \dots$

4. Build all Configuration State functions for total spin S and Irrep Γ $\Theta_I = \sqrt{\frac{1}{3}}\Phi_{23} - \sqrt{\frac{2}{3}}\Phi_{51}$

5. Calculate Hamiltonian matrix elements $H_{IJ}^{LFT}(\mathbf{e}, B, C) = \langle \Theta_I | \hat{H}_{LFT} + H_{SB} + H_{ZE} | \Theta_J \rangle$

6. Diagonalize the ligand field Hamiltonian $\mathbf{H}^{LFT} \mathbf{C}^{LFT} = E \mathbf{C}^{LFT}$

→ Yields all **ligand field multiplets** as a function of the LFT parameters.
Order them in the **Tanabe-Sugano diagrams**

Parameterization : The Angular Overlap Model

In the AOM the one-electron part of the ligand field is written as:

$$h_{ab} = \sum_L \sum_{\lambda} F_{\lambda a}(\theta_L, \phi_L, \psi_L) F_{\lambda b}(\theta_L, \phi_L, \psi_L)$$

L= sum over ligands

F= angular factor (symmetry!)

**$e_{\sigma,\pi}$ = Interaction parameter
(ligand specific, transferrable)**

Two-electron part of the ligand field

$$\langle d_{\sigma} d_{\sigma} | r_{12}^{-1} | d_{\pi} d_{\pi} \rangle = \langle d_{\sigma} d_{\sigma} | r_{12}^{-1} | d_{\pi'} d_{\pi'} \rangle = F_{dd}^0 + \frac{2}{49} F_{dd}^2 - \frac{24}{441} F_{dd}^4$$

$$\langle d_i d_i | r_{12}^{-1} | d_i d_i \rangle = F_{dd}^0 + \frac{4}{49} F_{dd}^2 + \frac{36}{441} F_{dd}^4$$

$$\langle d_{\delta} d_{\delta} | r_{12}^{-1} | d_{\delta'} d_{\delta'} \rangle = F_{dd}^0 + \frac{4}{49} F_{dd}^2 - \frac{34}{441} F_{dd}^4$$

etc.



$$A = F_{dd}^0 - \frac{49}{441} F_{dd}^4$$

$$B = \frac{1}{49} F_{dd}^2 - \frac{5}{441} F_{dd}^4$$

$$C = \frac{35}{441} F_{dd}^4 \approx 4B$$

**Racah parameters
(A=arbitrary, C=4B).**

Just two (one) parameter describes the electron-electron repulsion, 1-3 parameters for each ligand.

Examples of AOM parameters

Ligand	e_{σ} (cm ⁻¹)	e_{π} (cm ⁻¹)
CN ⁻	7530	-930
F ⁻	7390	1690
Cl ⁻	5540	1160
Br ⁻	4920	830
I ⁻	4100	670
NH ₃	7030	0
en	7260	0
py	5850	-670
H ₂ O	7900	1850
OH ⁻	8670	3000

Note: $10Dq = \Delta = 3e_{\sigma} - 4e_{\pi}$

Examples of AOM parameters

Ligand	e_{σ} (cm ⁻¹)	e_{π} (cm ⁻¹)
<div>BUT<ul style="list-style-type: none">- BIG overfitting problem- No unique, precise definition of ligand field parameters</div>		
OH ⁻	8670	3000

Note: $10Dq = \Delta = 3e_{\sigma} - 4e_{\pi}$

Good numbers:

Large-Scale Multireference Calculations

Properly calculating coordination compounds?

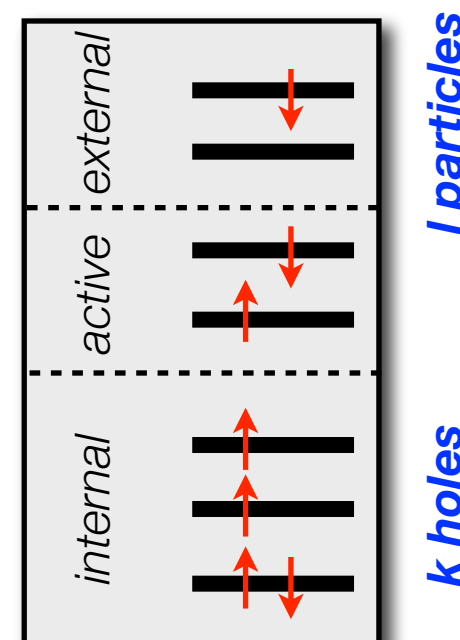
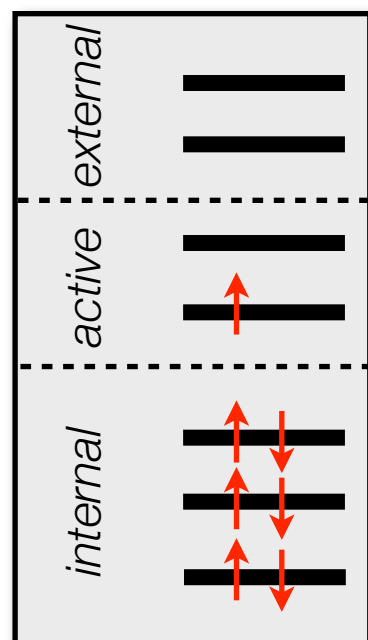
- ✓ Explicit representation of all spatial components of (nearly) degenerate states
 - ▶ Multireference wavefunction theory
- ✓ Realistic calculation of energy splittings between nearly degenerate spatial components
 - ▶ Accurate structural data!
 - ▶ Dynamic correlation
- ✓ Explicit representation of all magnetic sublevels $M = S, S-1, \dots, -S$ of each spatial component
 - ▶ Wigner-Eckart theorem
- ✓ Interaction of all of these magnetic sublevels via the SOC (+SSC) and Zeeman operators
 - ▶ Quasi-degenerate perturbation theory
- ✓ Perhaps a significant number of excited states of various multiplicities that also interact with the degenerate pair via SOC
 - ▶ Multiroot theory

Multireference Expansions

- ✓ Cover differential dynamic correlation by performing excitations relative to at least the important CSFs of the CAS

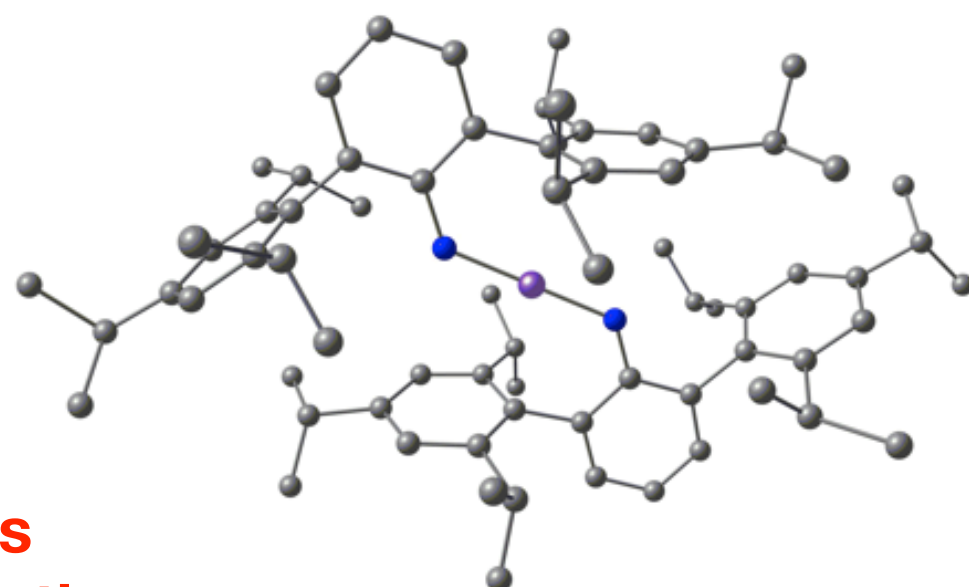
$$\left| \Psi_I^{SM;MRCI} \right\rangle = \sum_{K \in CAS} C_{KI} \left| \Phi_K^{SM} \right\rangle + \sum_{K \notin CAS} C_{KI} \left| \Phi_K^{SM} \right\rangle$$

Solve:
 $\mathbf{HC} = \mathbf{EC}$



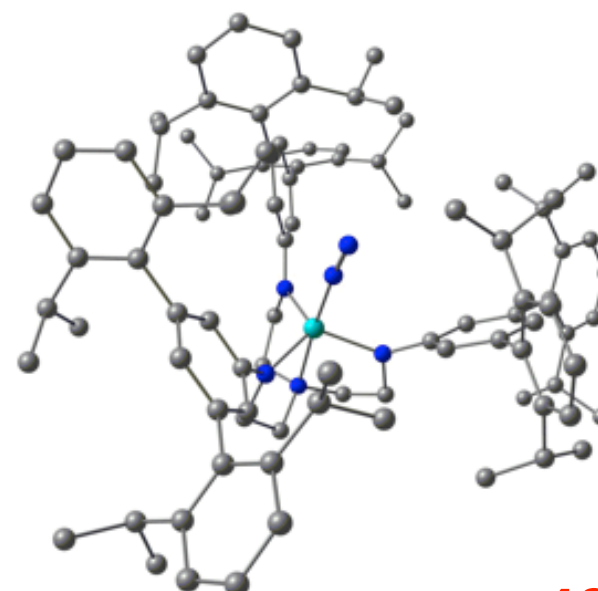
- ✓ Of the HUGE number of **(k-holes, l-particle)** CSFs, only a small fraction contributes to the differential dynamic correlation energy (idea of **DDCI** (Malrieu) and **SORCI** (FN))
- ▶ Computationally affordable computation protocols are **CASSCF/CASPT2** (MOLCAS) and **CASSCF/NEVPT2** (ORCA)

Large Scale Calculations with NEVPT2



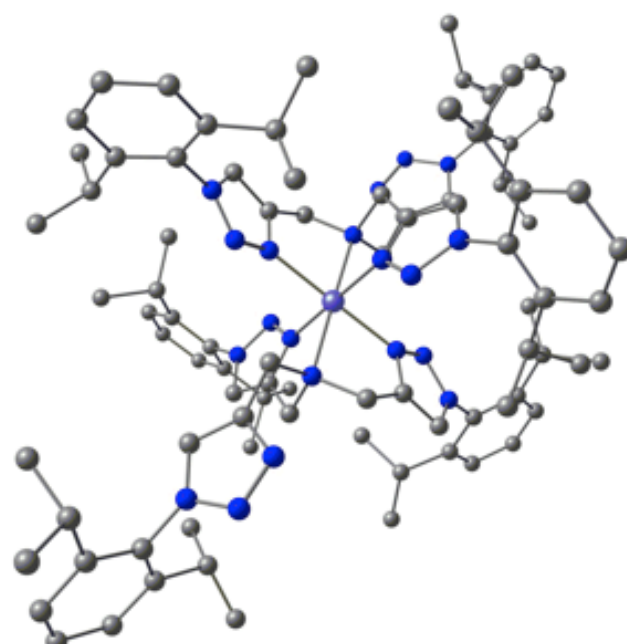
175 Atoms
2939 basis functions

$\text{FeC}_{72}\text{N}_2\text{H}_{100}$



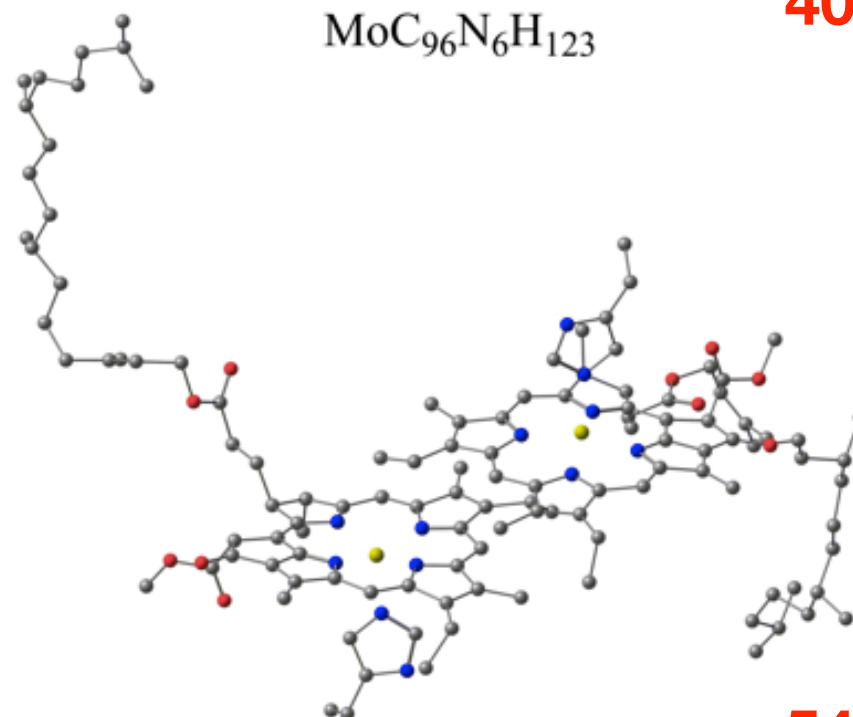
226 Atoms
4073 basis functions

$\text{MoC}_{96}\text{N}_6\text{H}_{123}$



231 Atoms
4175 basis functions

$\text{NiC}_{90}\text{N}_{20}\text{H}_{120}$



304 Atoms
5426 basis functions

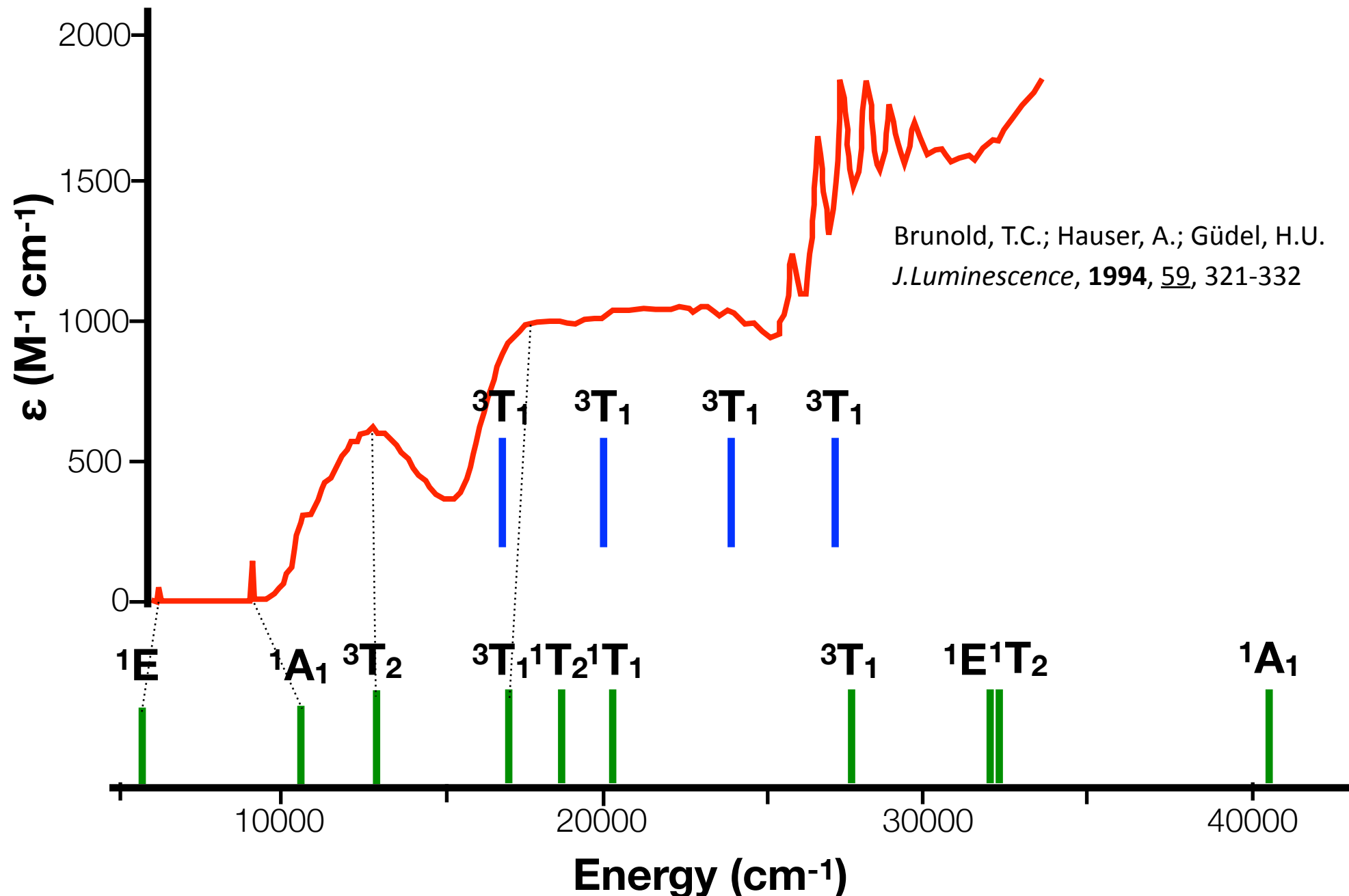
$\text{Mg}_2\text{C}_{120}\text{N}_{12}\text{H}_{160}\text{O}_{10}$

Ab Initio Ligand Field Theory

Ligand Field Theory as an Effective Hamiltonian

FeO_4^{2-} doped in K_2SeO_4

Atanasov, M.; Sivalingham, K.;
Ganyushin, D.; *FN Struc. Bond.*,
2012, pp 149



Brunold, T.C.; Hauser, A.; Güdel, H.U.
J.Luminescence, **1994**, 59, 321-332

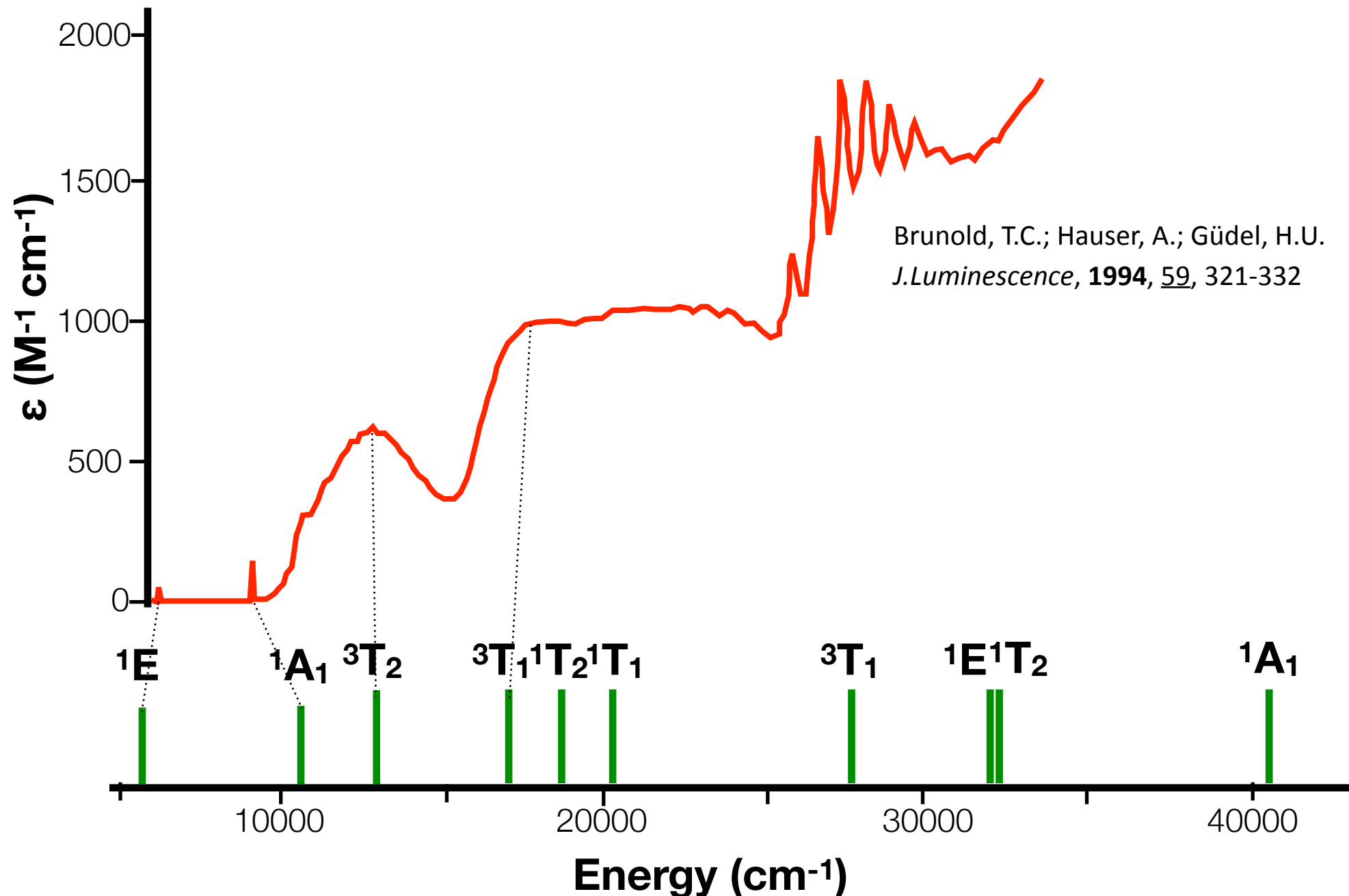
LMCT
NOT Covered by
Ligand Field Theory

d-d
Covered by Ligand
Field Theory

Ligand Field Theory as an Effective Hamiltonian

FeO_4^{2-} doped in K_2SeO_4

Atanasov, M.; Sivalingham, K.;
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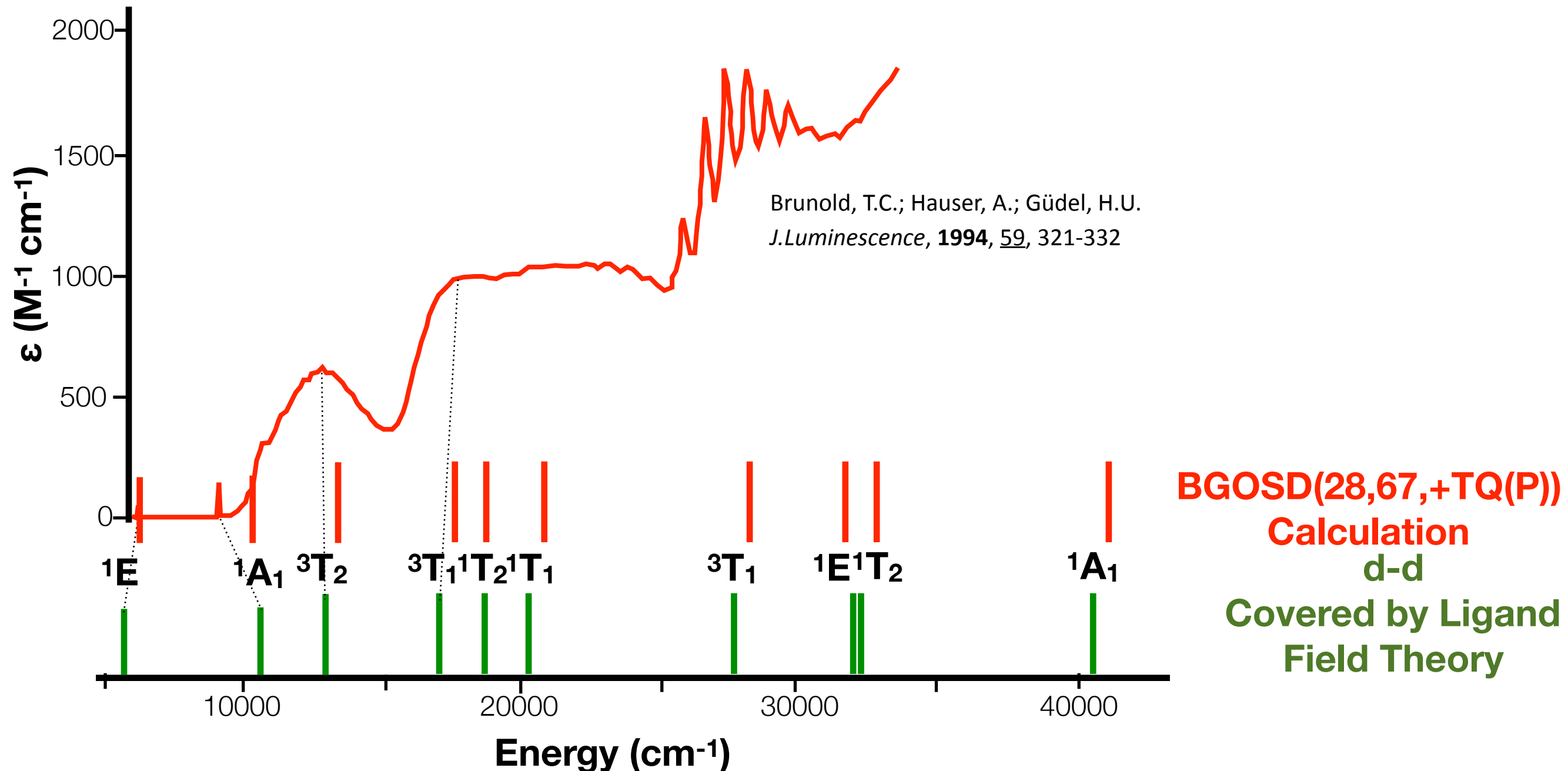


d-d
Covered by Ligand
Field Theory

Ligand Field Theory as an Effective Hamiltonian

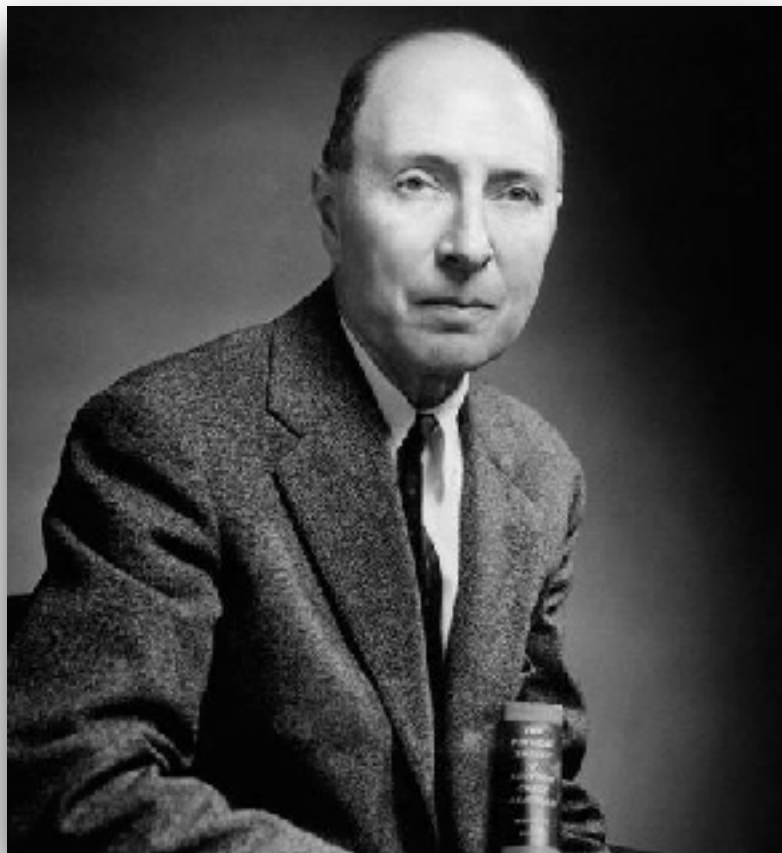
FeO_4^{2-} doped in K_2SeO_4

Atanasov, M.; Sivalingham, K.;
Ganyushin, D.; *FN Struc. Bond.*,
2012, pp 149



„It is nice to know that the computer understands the problem. But I would like to understand it too.“

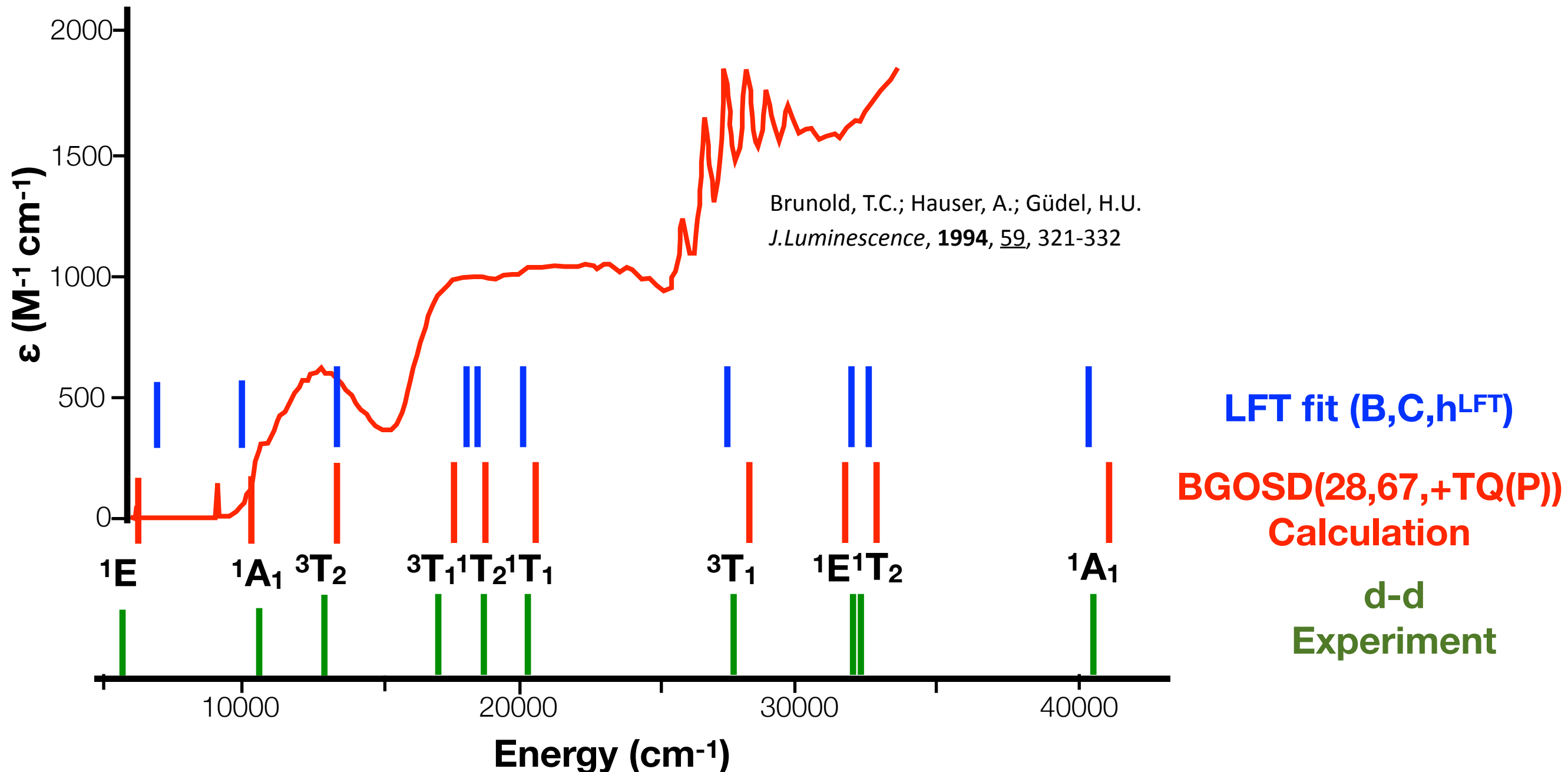
–Eugene Wigner



Ligand Field Theory as an Effective Hamiltonian

FeO_4^{2-} doped in K_2SeO_4

Atanasov, M.; Sivalingham, K.;
Ganyushin, D.; *FN Struc. Bond.*, **2012**, pp 149



Unambiguous Match between NEVPT2 and LFT

There is a 1:1 correspondance between the ligand field CSFs and the CAS-CI CSFs.

$$\Theta_I^{LFT} = | d_{xy}^\alpha d_{xz}^\beta \dots d_{z^2}^\alpha |$$

Ligand field pure d-orbital

$$\Theta_I^{CASSCF} = | \psi_{xy}^\alpha \psi_{xz}^\beta \dots \psi_{z^2}^\alpha |$$

Ab initio molecular orbital with metal d-parentage

Thus, all we have to ensure is that ligand field d-orbitals and CASSCF molecular orbitals of the same parentage are ordered in the same way and that CSFs are constructed in the same way.

The *ab initio* Intermediate Hamiltonian

1. We have a „**model space**“ that contains all the essential physics that we want to describe. This is the **CAS(n,5) space** of N-particle wavefunctions that cleanly maps onto the ligand field manifold
2. We have a „**outer space**“ that brings in all the remaining effects of dynamic correlation

$$H_{IJ}^{eff} = \langle \Phi_I | H | \Phi_J \rangle + \sum_{K \in roots} \langle \Phi_I | \Psi_K \rangle \Delta E_K^{(NEVPT2)} \langle \Psi_K | \Phi_J \rangle$$

- ✓ **Dynamic and non-dynamic correlation energy**
- ✓ **Same dimension as ligand field CI matrix**
- ✓ **Completely *ab initio*. No parameters!**
- ✓ **Yields a near exact eigenvalue spectrum**

The *ab initio* Intermediate Hamiltonian

1. We have a „**model space**“ that contains all the essential physics that we want to describe. This is the **CAS(n,5) space** of N-particle wavefunctions that cleanly maps onto the ligand field manifold
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- ✓ **Dynamic and non-dynamic correlation energy**
- ✓ **Same dimension as ligand field CI matrix**
- ✓ **Completely *ab initio*. No parameters!**
- ✓ **Yields a near exact eigenvalue spectrum**

The condition is then that the ligand field CI matrix should resemble the *ab initio* effective Hamiltonian as closely as possible

$$\left| H_{IJ}^{LFT} - H_{IJ}^{eff} \right| = \min$$

For each matrix element!

While this looks at first sight to be a nonlinear optimization problem, in reality things are easy **because the ligand field matrix is linear in each and every ligand field parameter!**

$$\mathbf{H}^{LFT}(e, B, C) = \mathbf{H}^{LFT}(0) + \frac{\partial \mathbf{H}^{LFT}}{\partial B} B + \frac{\partial \mathbf{H}^{LFT}}{\partial C} C + \sum_L \frac{\partial \mathbf{H}^{LFT}}{\partial e_L} e_L$$

This ensures that there is a **unique least squares solution** that provides the unambiguous best fit of the ligand field and effective Hamiltonian matrices:

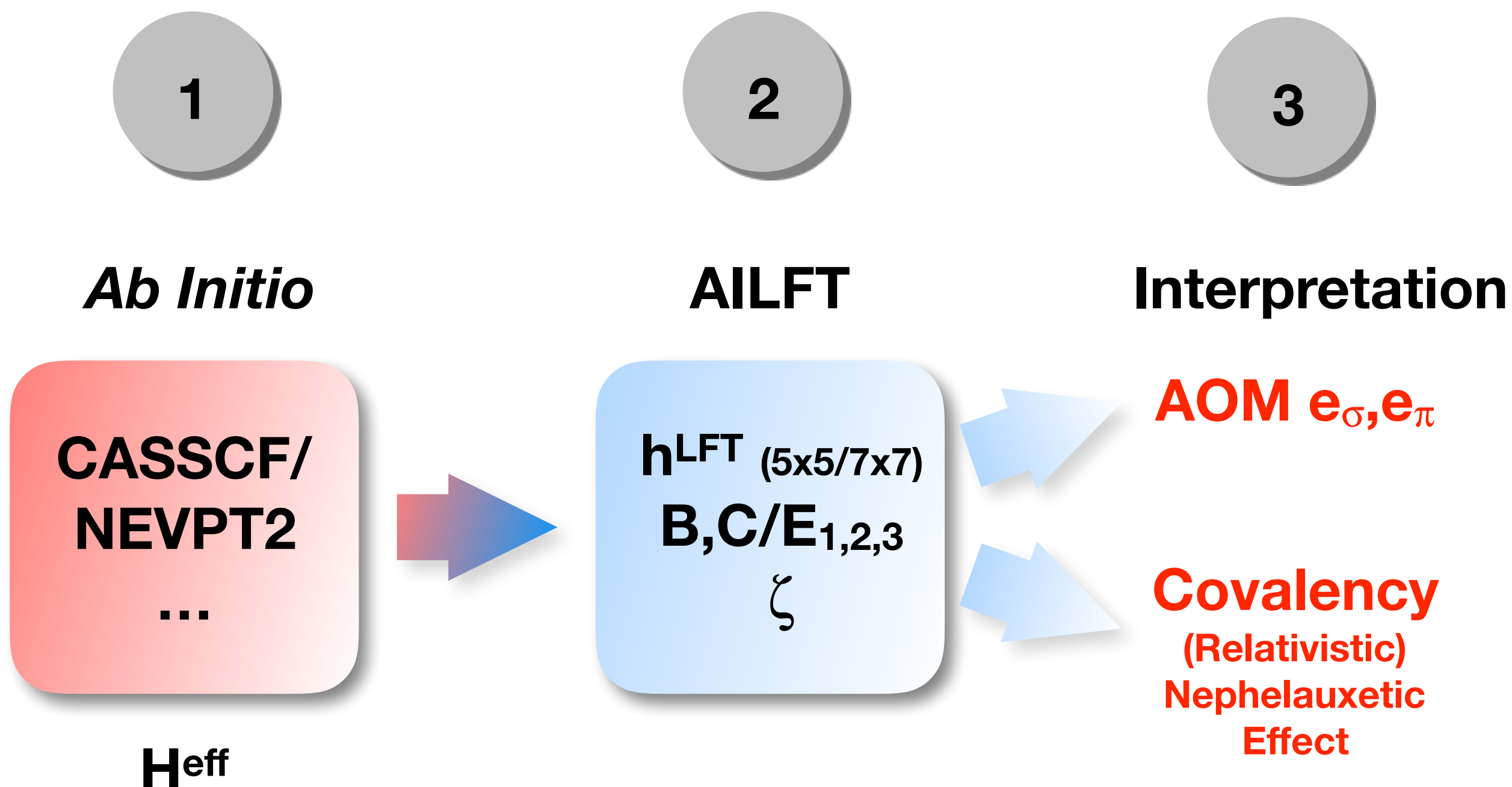
$$\mathbf{A}\mathbf{p} = -\mathbf{b} \Leftrightarrow \mathbf{p} = -\mathbf{A}^{-1}\mathbf{b}$$

p_K = The k'th ligand field parameter

$$A_{KL} = \sum_{IJ} \frac{\partial H_{IJ}^{LFT}}{\partial p_K} \frac{\partial H_{IJ}^{LFT}}{\partial p_L}$$

$$b_K = \sum_{IJ} \frac{\partial H_{IJ}^{LFT}}{\partial p_K} H_{IJ}^{eff}$$

Summary: Information Flow in AILFT



Application to CrX_6^{3-} (X=F,Cl,Br,I)

**Prof. Mihail
Atanasov**



Calculated spectra

	CASSCF	NEVPT2	exp
$^4\text{T}_2$	13380	15365	15200
$^4\text{T}_1(1)$	21424	23449	21800
$^4\text{T}_1(2)$	34778	35307	35000
$^2\text{E}(1)$	19409	17579	16300
$^2\text{T}_1(1)$	20430	18681	16300
$^2\text{T}_2(1)$	27372	25288	23000
$^2\text{A}_1$	29674	30022	
$^2\text{T}_2(2)$	32664	32979	
$^2\text{T}_1(2)$	33736	33800	
$^2\text{E}(2)$	36009	35316	
$^2\text{T}_1(3)$	39668	39324	
$^2\text{T}_2(3)$	46237	45791	
$^2\text{T}_1(4)$	46546	48689	
$^2\text{A}_2$	51121	47612	
$^2\text{T}_2(4)$	52510	49902	
$^2\text{E}(3)$	55140	55611	

Application to CrX_6^{3-} ($\text{X}=\text{F}, \text{Cl}, \text{Br}, \text{I}$)

Prof. Mihail Atanasov



Calculated spectra

	CASSCF	NEVPT2	exp
$^4\text{T}_2$	13380	15365	15200
$^4\text{T}_1(1)$	21424	23449	21800

Extracted parameters

	CrF_6^{3-}			CrCl_6^{3-}			CrBr_6^{3-}			CrI_6^{3-}	
	CASSCF	NEVPT2	Exp.	CASSCF	NEVPT2	Exp.	CASSCF	NEVPT2	Exp.	CASSCF	NEVPT2
10Dq	13359	15255	15297	10941	14109	12630	9876	13542	13089	9258	13533
B	1071	863	734	988	767	632	972	729	471	943	699
C	4018	3720	3492	3907	3605	3180	3886	3631	3249	3841	3627
C/B	3.75	4.31	4.76	3.95	4.70	5	4.00	4.98	6.90	4,07	5.18

$^2\text{T}_1(3)$	46207	48751	
$^2\text{T}_1(4)$	46546	48689	
$^2\text{A}_2$	51121	47612	
$^2\text{T}_2(4)$	52510	49902	
$^2\text{E}(3)$	55140	55611	

Application to CrX_6^{3-} ($\text{X}=\text{F}, \text{Cl}, \text{Br}, \text{I}$)

Prof. Mihail
Atanasov



Calculated spectroscopic parameters			CrF ₆ ³⁻	CASSCF	NEVPT2	Deviations from Ligand field fit and ab initio calculations Extracted parameters					
			⁴ T ₂	-20	-111						
		⁴ T ₂	⁴ T ₁ (1)	-42	-533	0					
		⁴ T ₁ (1)	⁴ T ₁ (2)	-16	439	0					
		C	² E(1)	-64	-283			CrBr ₆ ³⁻		CrI ₆ ³⁻	
			² T ₁ (1)	-38	-575	Exp.	CASSCF	NEVPT2	Exp.	CASSCF	NEVPT2
			² T ₂ (1)	-85	-119						
			² A ₁	25	-156	12630	9876	13542	13089	9258	13533
10Dq	13359	1	² T ₂ (2)	88	-527	632	972	729	471	943	699
B	1071	8	² T ₁ (2)	20	-678	3180	3886	3631	3249	3841	3627
C	4018	3	² E(2)	10	-151	5	4.00	4.98	6.90	4,07	5.18
C/B	3.75		² T ₁ (3)	90	-694						
			² T ₁ (4)								
			² A ₂								
			² T ₂ (4)								
			² E(3)								
			² T ₂ (3)	-24	-797						
			² T ₁ (4)	89	-767						
			² A ₂	-4	-480						
			² T ₂ (4)	-11	762						
			² E(3)	-166	-1017						

Application to CrX_6^{3-} (X=F,Cl,Br,I)

Prof. Mihail Atanasov



Calculated s		CrF_6^{3-}	CASSCF	NEVPT2
		4T_2	-20	-111
4T_2		$^4T_1(1)$	-42	-533
$^4T_1(1)$		$^4T_1(2)$	-16	439
		$^2E(1)$	-64	-283
		$^2T_1(1)$	-38	-575
		$^2T_2(1)$	-85	-119
10Dq	13359	2A_1	25	-156
B	1071	$^2T_2(2)$	88	-527
C	4018	$^2T_1(2)$	20	-678
C/B	3.75	$^2E(2)$	10	-151
		$^2T_1(3)$	90	-694
		$^2T_2(3)$	-24	-797
		$^2T_1(4)$	89	-767
		2A_2	-4	-480
		$^2T_2(4)$	-11	762
		$^2E(3)$	-166	-1017

Deviations from Ligand field fit
and ab initio calculations
Extracted parameters

		CrBr_6^{3-}		CrI_6^{3-}	
Exp.	CASSCF	NEVPT2	Exp.	CASSCF	NEVPT2
12630	9876	13542	13089	9258	13533
632	972	729	471	943	699
3180	3886	3631	3249	3841	3627
5	4.00	4.98	6.90	4,07	5.18

- ✓ Very good agreement between *ab initio* values and empirical values for 10Dq
- ✓ Still slight overestimation of electron repulsion (basis set incompleteness + 2nd order perturbation theory)
- ✓ Ligand field fit to CASSCF near perfect, to NEVPT2 within 1000 cm^{-1}

Periodic Trends in Lanthanide Compounds through the Eyes of Multireference ab Initio Theory

Daniel Aravena,[†] Mihail Atanasov,^{*,‡,§} and Frank Neese^{*,‡}

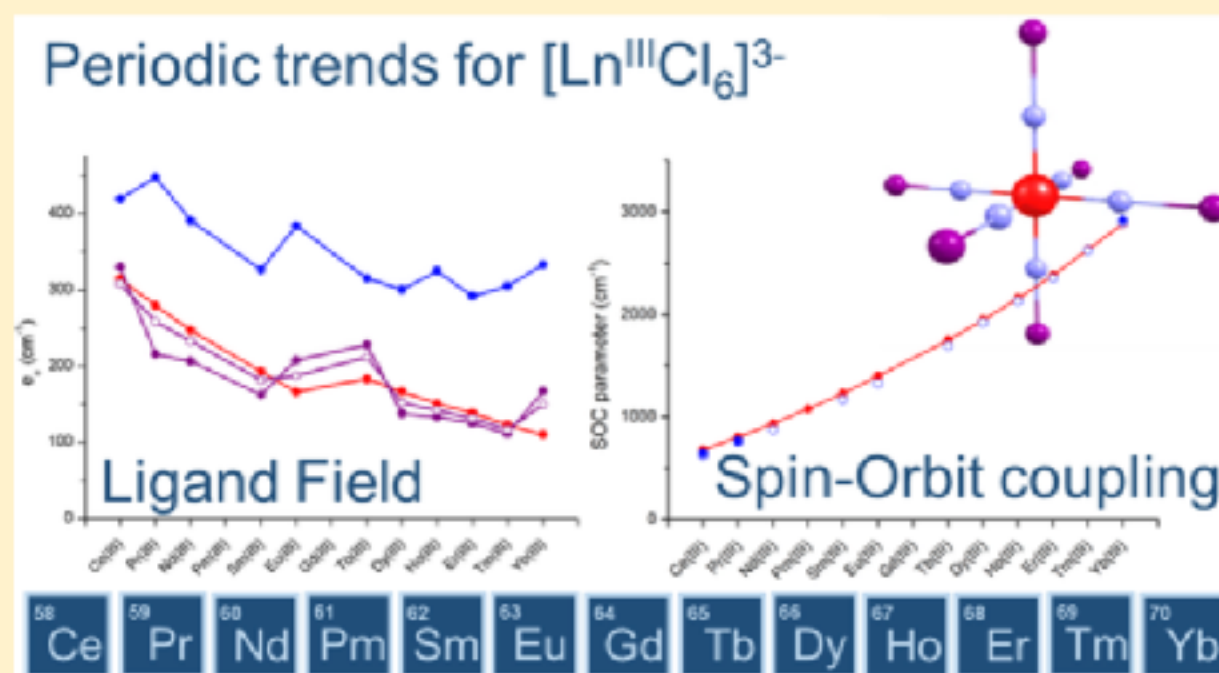
[†]Facultad de Química y Biología, Universidad de Santiago de Chile, Casilla 40, Correo 33, Santiago, Chile

[‡]Max Planck Institut für Chemische Energiekonversion, Stifstr 34-36, D-45470 Mülheim an der Ruhr, Germany

[§]Bulgarian Academy of Sciences, Institute of General and Inorganic Chemistry, Akad. Georgi Bontchev Street 11, 1113 Sofia, Bulgaria

Supporting Information

ABSTRACT: Regularities among electronic configurations for common oxidation states in lanthanide complexes and the low involvement of f orbitals in bonding result in the appearance of several periodic trends along the lanthanide series. These trends can be observed on relatively different properties, such as bonding distances or ionization potentials. Well-known concepts like the lanthanide contraction, the double–double (tetrad) effect, and the similar chemistry along the lanthanide series stem from these regularities. Periodic trends on structural and spectroscopic properties are examined through complete active space self-consistent field (CASSCF) followed by second-order *N*-electron valence perturbation theory (NEVPT2) including both scalar relativistic and spin–orbit



Inorganic Chemistry, **2016,2016**, 55 (9), pp 4457–4469

Case Study: Tetrahedral Co(II) Complexes

J | A | C | S
JOURNAL OF THE AMERICAN CHEMICAL SOCIETY

Communication

pubs.acs.org/JACS

Slow Magnetic Relaxation at Zero Field in the Tetrahedral Complex $[\text{Co}(\text{SPh})_4]^{2-}$

Joseph M. Zadrozny and Jeffrey R. Long*

[dx.doi.org/10.1021/ja2100142](https://doi.org/10.1021/ja2100142) | *J. Am. Chem. Soc.* 2011, 133, 20732–20734

Inorganic Chemistry

ARTICLE

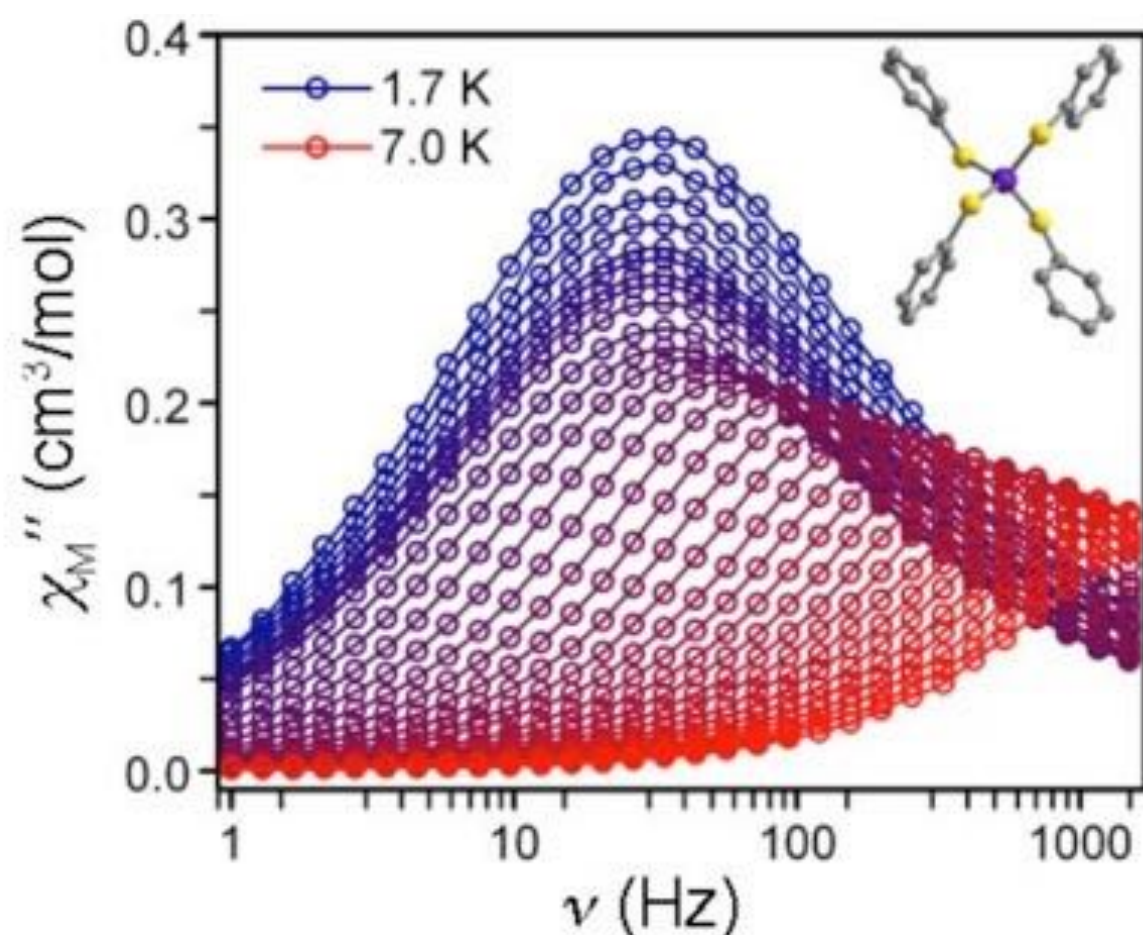
pubs.acs.org/IC

Theoretical Analysis of the Spin Hamiltonian Parameters in $\text{Co}^{(\text{II})}\text{S}_4$ Complexes, Using Density Functional Theory and Correlated ab initio Methods

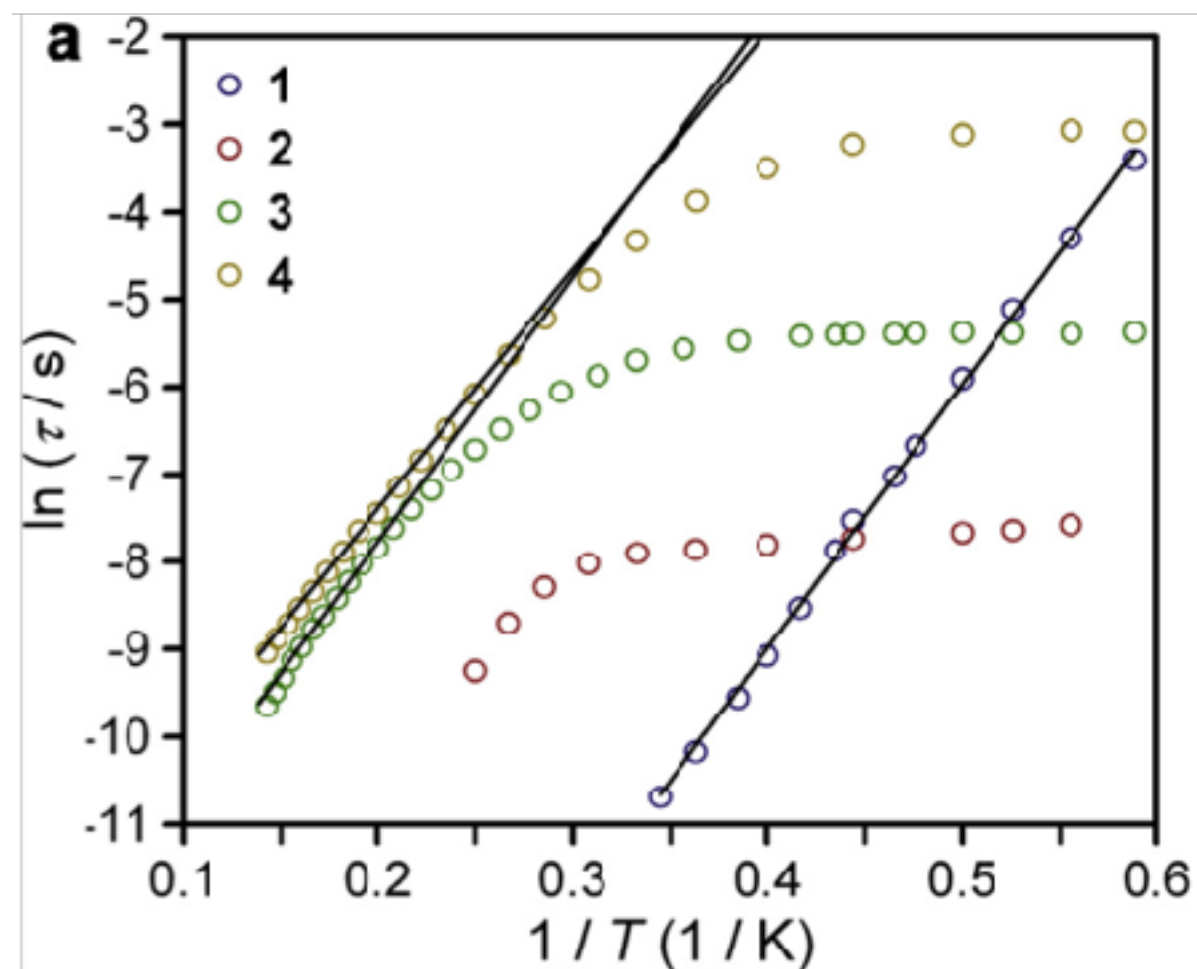
Dimitrios Maganas,^{†,§} Silvia Sottini,[‡] Panayotis Kyritsis,^{*,†} Edgar J. J. Groenen,[‡] and Frank Neese^{*,§}

Inorg. Chem. 2011, 50, 8741–8754

$[\text{Co}^{\text{II}}(-\text{S}-\text{Ph})_4]^{2-}$ shows slow magnetic relaxation



First mononuclear TM compound with Single Molecular Magnet behavior in the absence of any external magnetic field

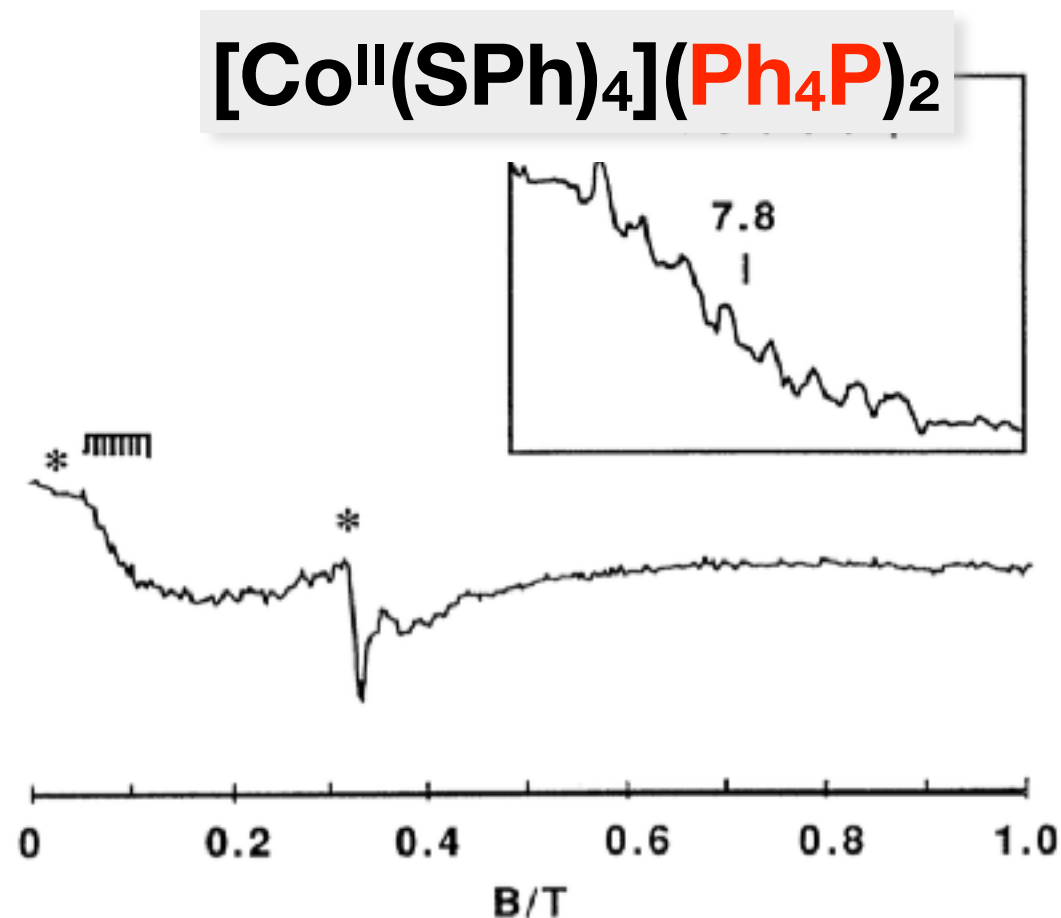


zfs is approximated from linear fit of the Arrhenius plot of:

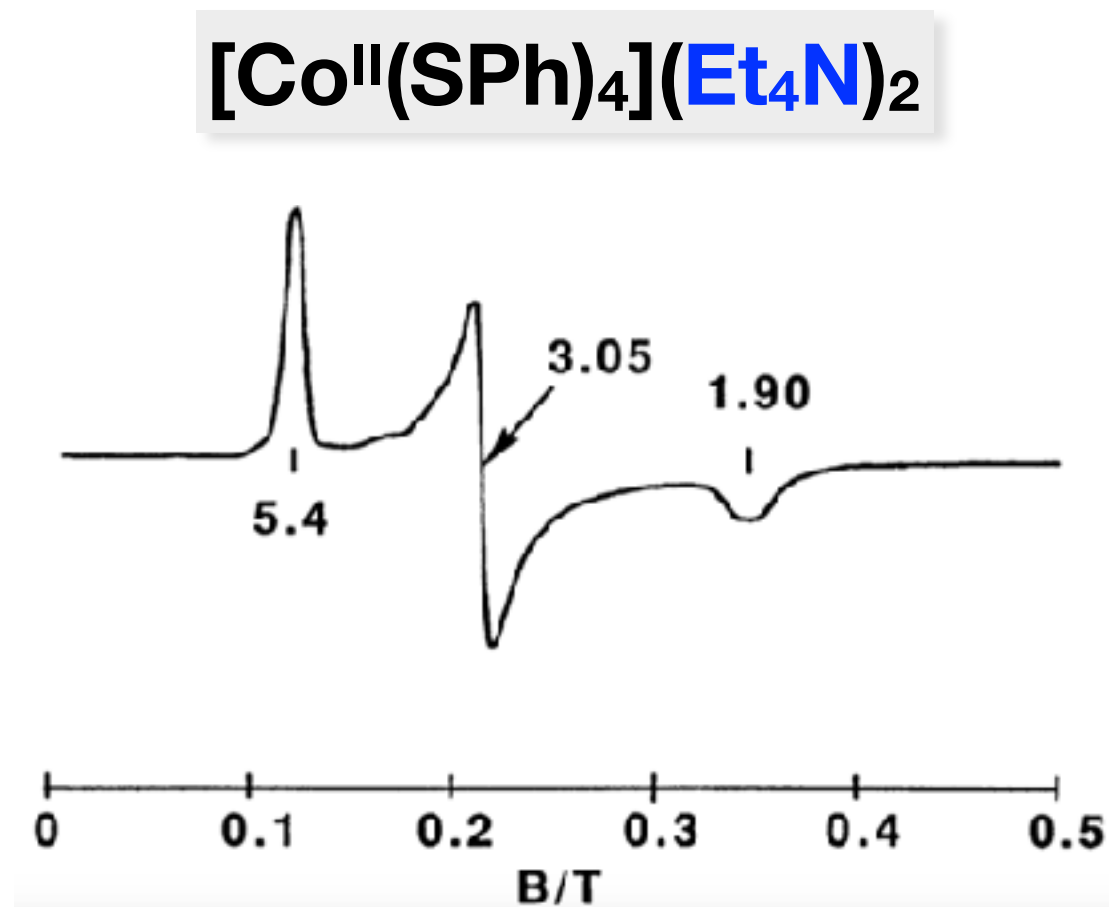
$$\tau \text{ vs } 1/T \text{ as } U_{\text{eff}} \sim |D|S^2 - 1/4 = 2D$$

- Due to relaxation processes $U_{\text{eff}} = 21 \text{ cm}^{-1}$ (But $D \sim 70 \text{ cm}^{-1}$ e.g. $\ll 140 \text{ cm}^{-1}$)
- This method is not conclusive for the magnitude of D

Serendipity: Giant Counterion Effects in $[\text{Co}^{\text{II}}(\text{RS})_4]^{2-}$



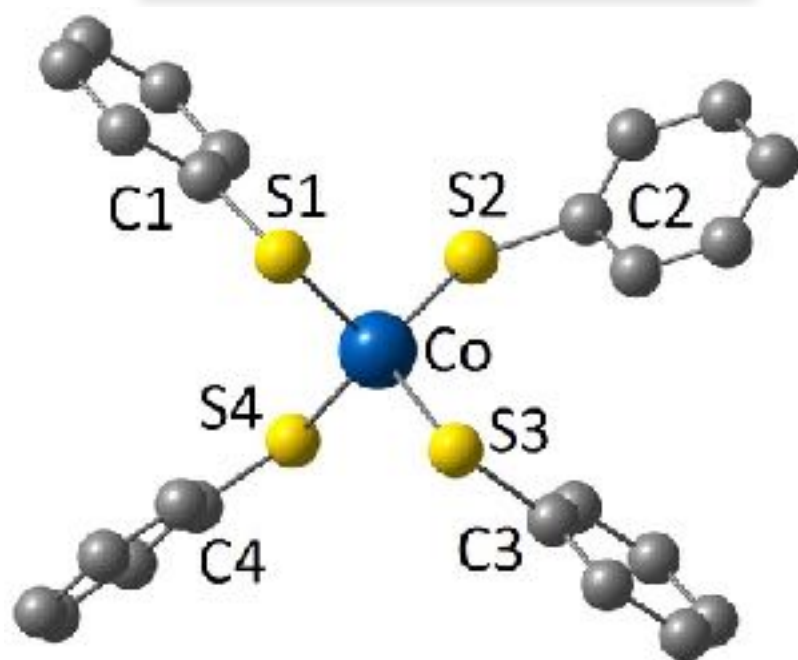
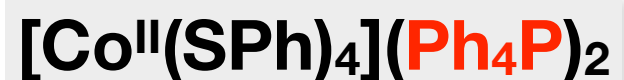
Strange Co(II) with large ZFS



Regular Co(II) with moderate ZFS

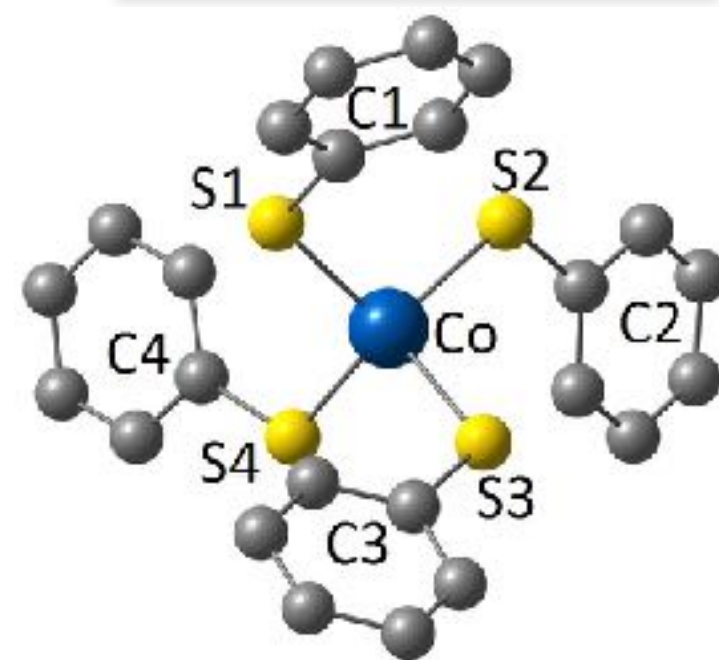
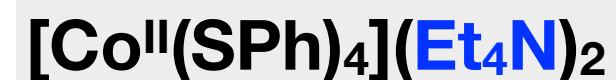
We need to understand these giant changes if we want to design a SMM

Geometric Origin of Counterion Effects



Tasks:

- ▶ Careful experimental study
- ▶ Detailed understanding of the electronic structure origin of the observed magnetic parameters



Elongated tetrahedron

Compressed tetrahedron

	[Co ^{II} (SPh) ₄](Ph ₄ P) ₂		Opt1	[Co ^{II} (SPh) ₄](Et ₄ N) ₂		Opt2
S1-Co1-S2	114	116	117	106	105	105
S2-Co1-S3	120			105		
S3-Co1-S4	112			101		
S4-Co1-S1	116			107		
S1-Co1-S3	98	98	96	119	120	119
S2-Co1-S4	97			121		

Energy difference between optimized structures ~1.3 kcal/mol

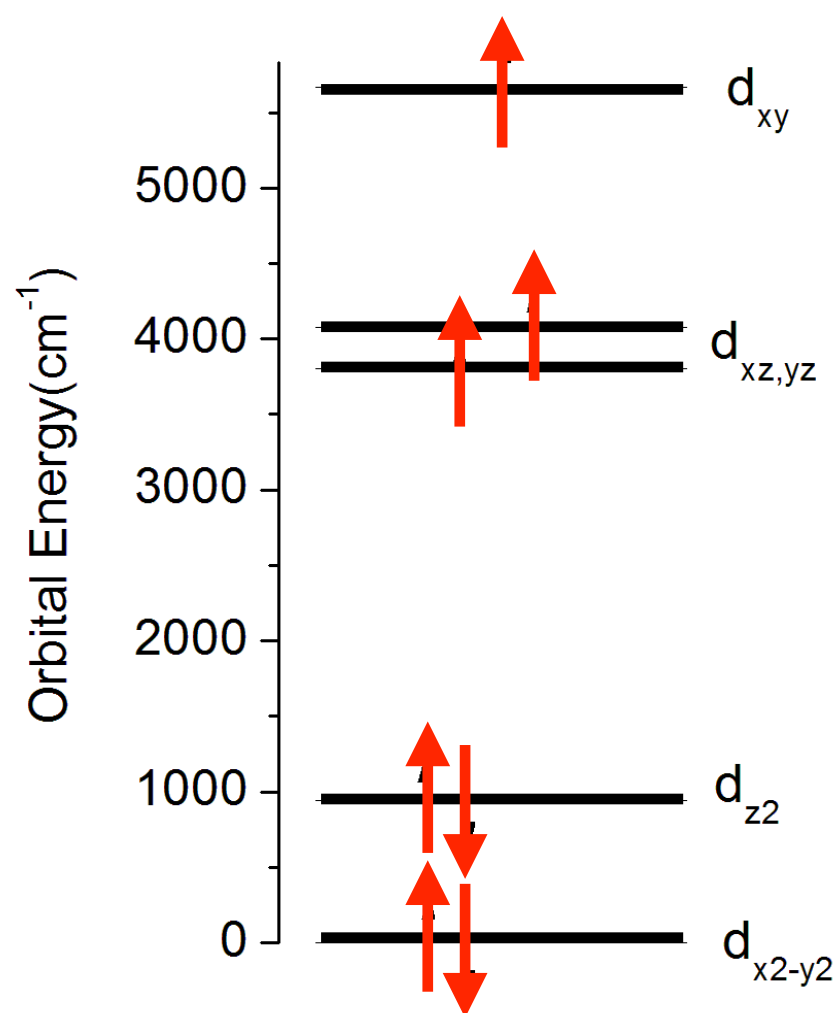
Quantitative Comparison

		D (cm⁻¹)	E/D	g_z	g_x	g_y
[Co^{II}(SPh)₄] (Ph₄P)₂	M and χ T fitting	-54(5)	<0.01	2.587(3)	2.1(1)	
	single crystal χ T and	-70(10)	<0.09	2.6(1)	2.2(1)	
	MCD	<-20	~0	2,5	2,8	
	NEVPT2/X-Ray	-46,7	0,03	2,71	2,13	2,18
	NEVPT2/Opt	-48,0	0,00	2,69	2,12	2,12
[Co^{II}(SPh)₄] (Et₄N)₂	M and χ T fitting	12(1)	0.31(2)	2.10(1)		
	MCD	+7	0.3	2.35		
	NEVPT2/X-Ray	11,5	0,16	2,16	2,29	2,33
	NEVPT2/Opt	3,0	0,00	2,19	2,23	2,23

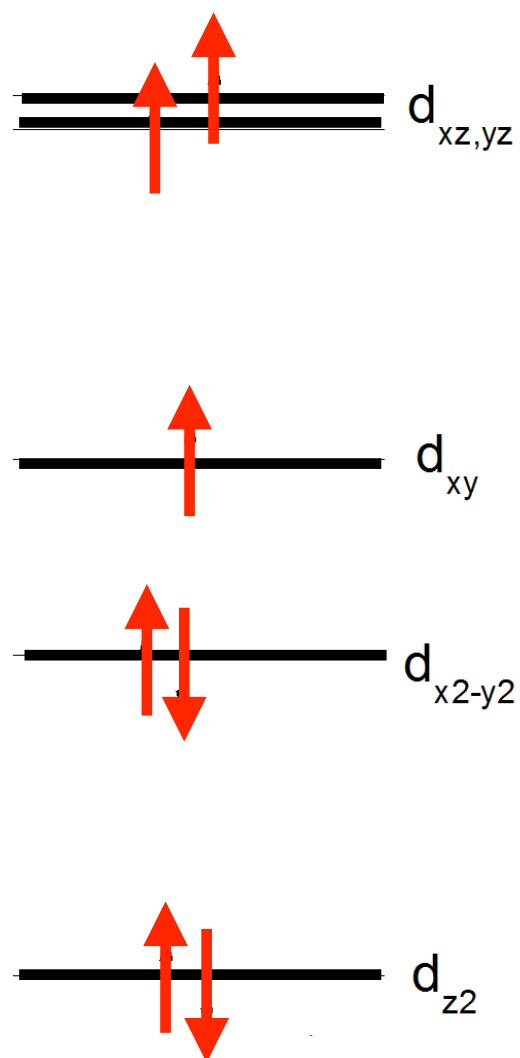
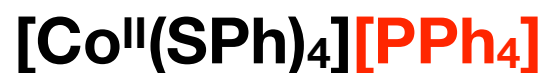
Good to excellent numerical agreement with experiment

Can now explore structural variations

Ab Initio Ligand Field Reconstruction



D > 0

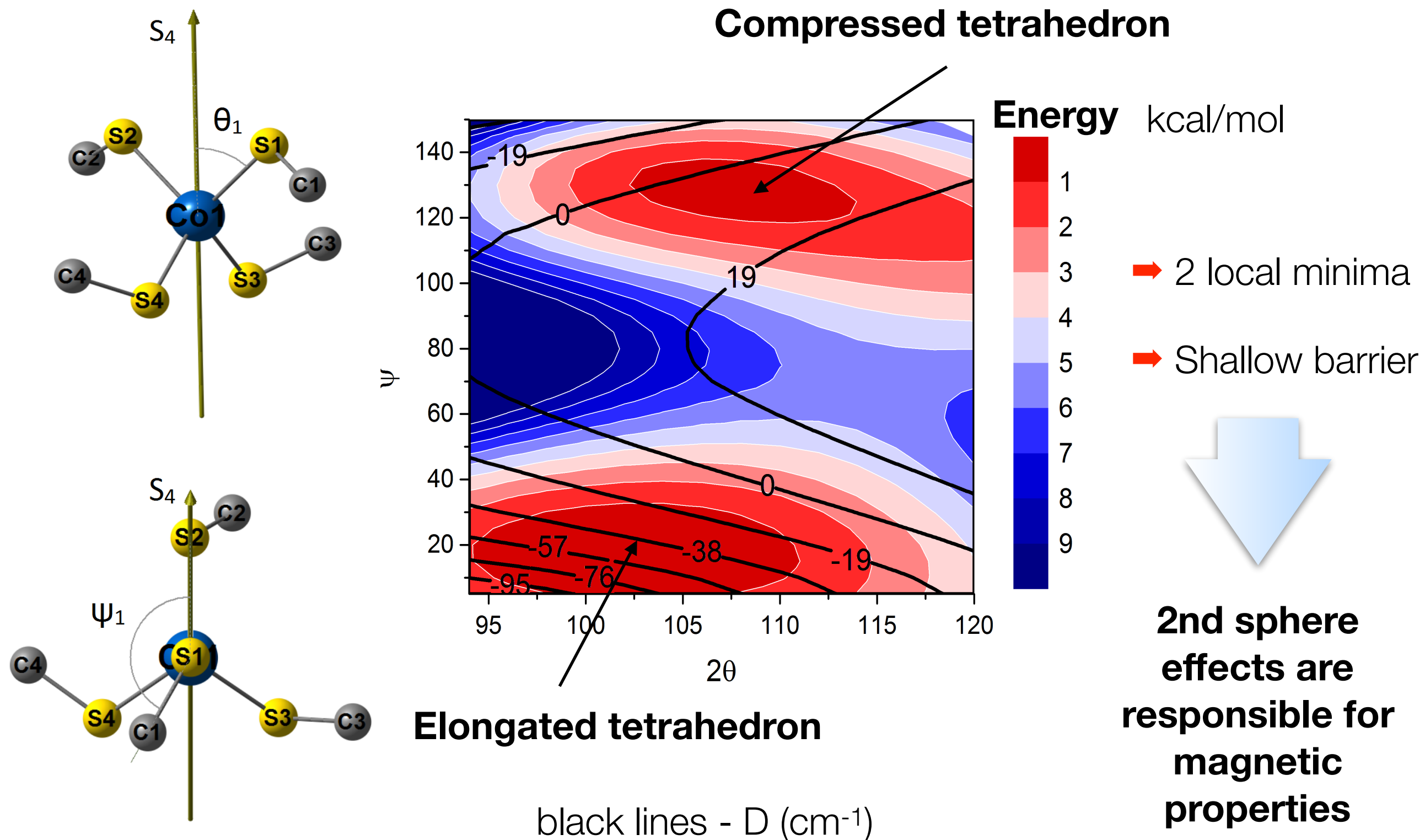


D < 0

$$D \simeq \frac{4}{9} \zeta_{eff}^2 \frac{E(^4T_{2z}) - E(^4T_{2x,y})}{(10Dq)^2}$$

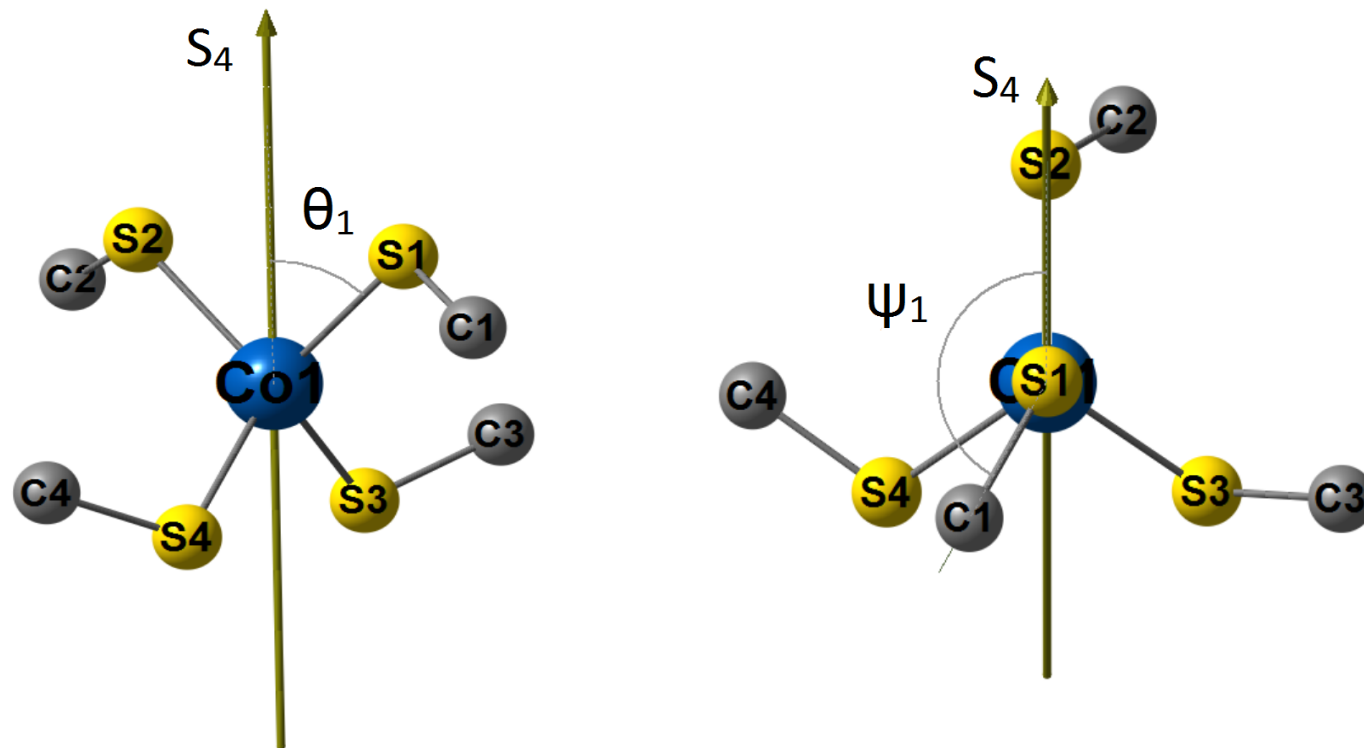
- ✓ Splittings of opposite sign induced by 'low-symmetry' ligand field.
- ✓ Much larger splittings for [PPh₄].
- ✓ Predicted sign of D consistent with observation

Magneto-Structural Correlations

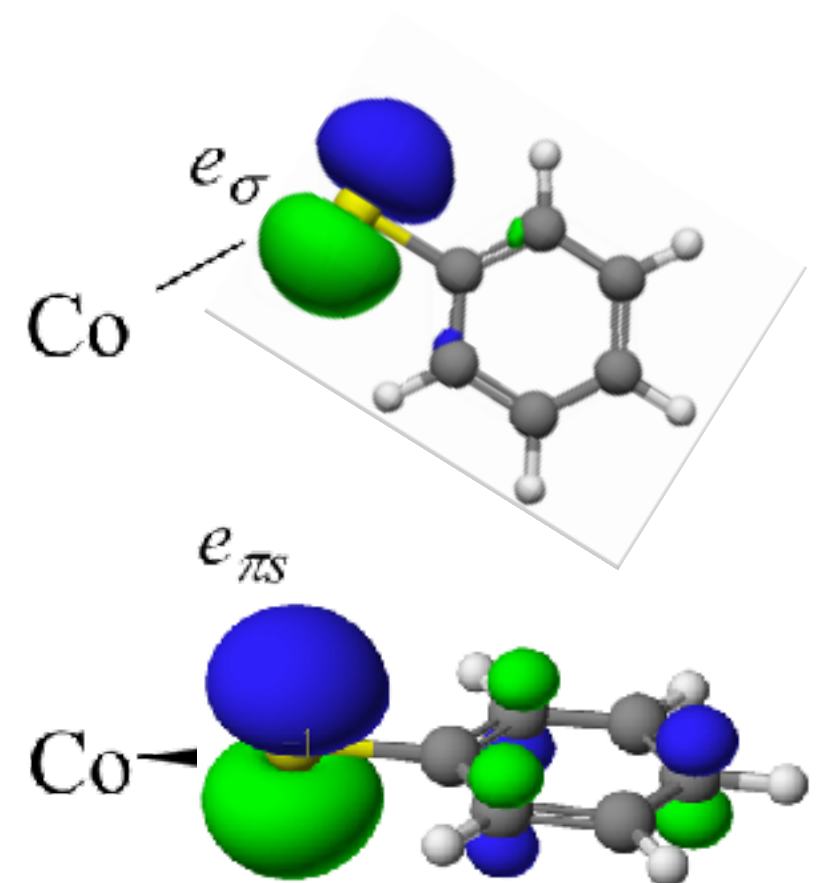


Qualitative Interpretation: Misdirected Valence

π -anisotropy of metal-ligand bonding:



- ➔ Counter Ion/Ligand interactions determine the ligand orientation.
- ➔ Ligand orientation determines lone pair orientations
- ➔ Metal/lone-pair interaction determines orbital splittings
- ➔ Orbital splittings determine magnetic properties through spin-orbit coupling

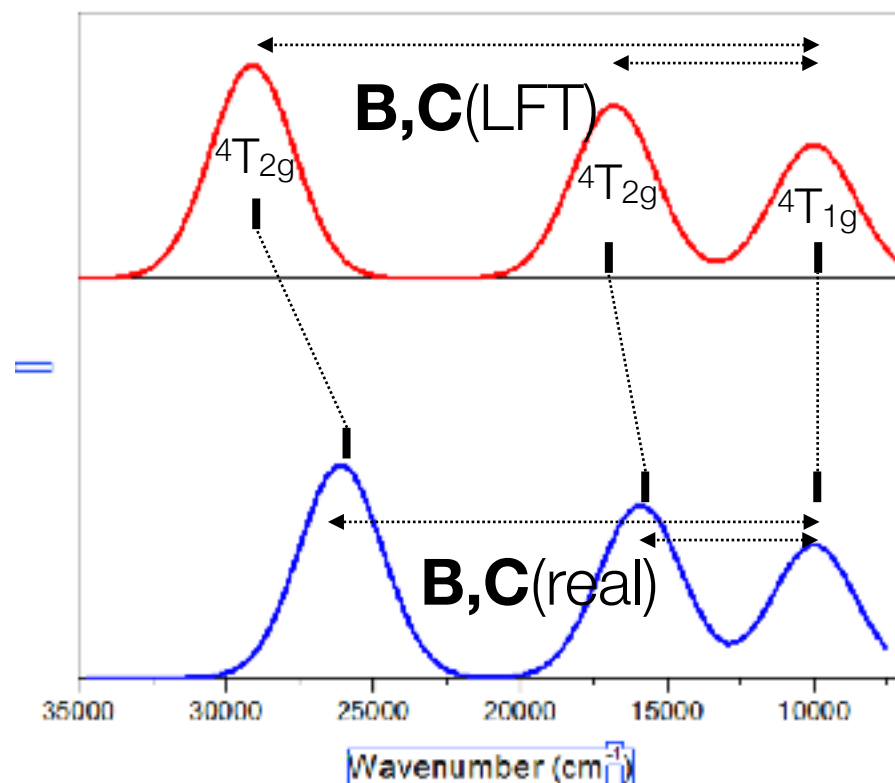


} **HUGE second sphere effect on magnetic properties**

Covalency, Oxidation States and MO Theory

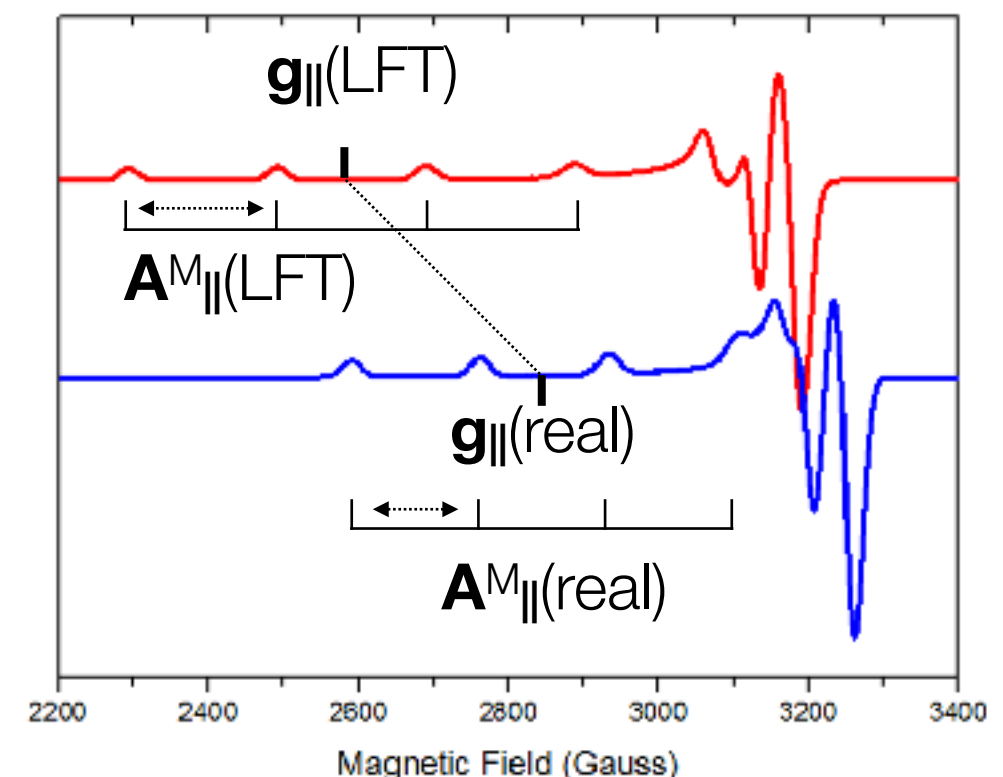
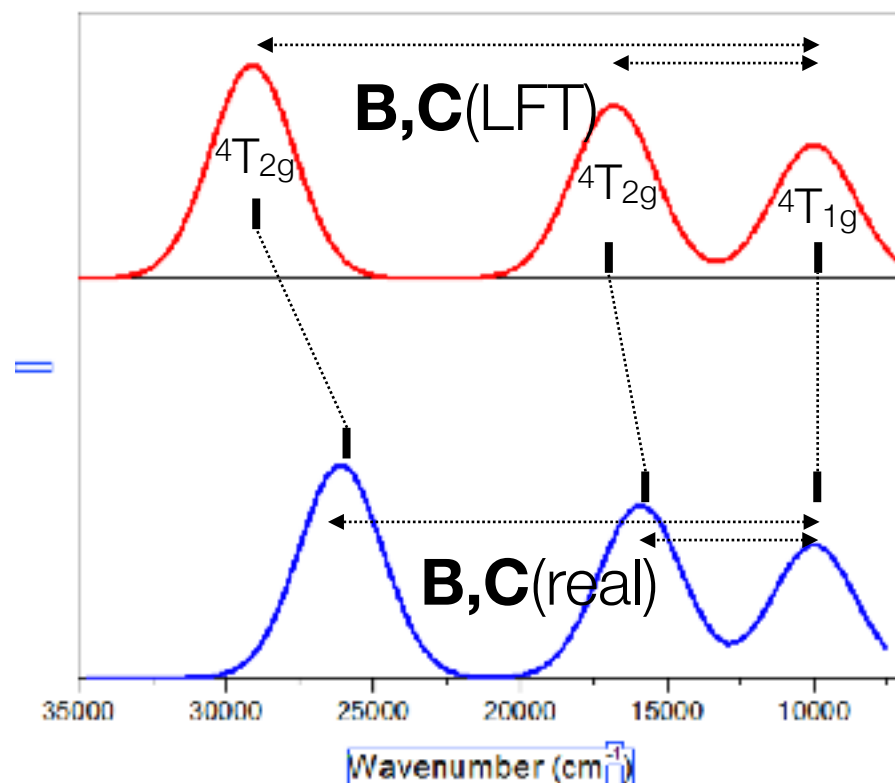
The Nephelauxetic Effects

- ➔ Observation 1: **Racah parameters** fitted for ions in complexes are smaller than those obtained for free-ions from atomic spectroscopy (**nephelauxetic effect**)



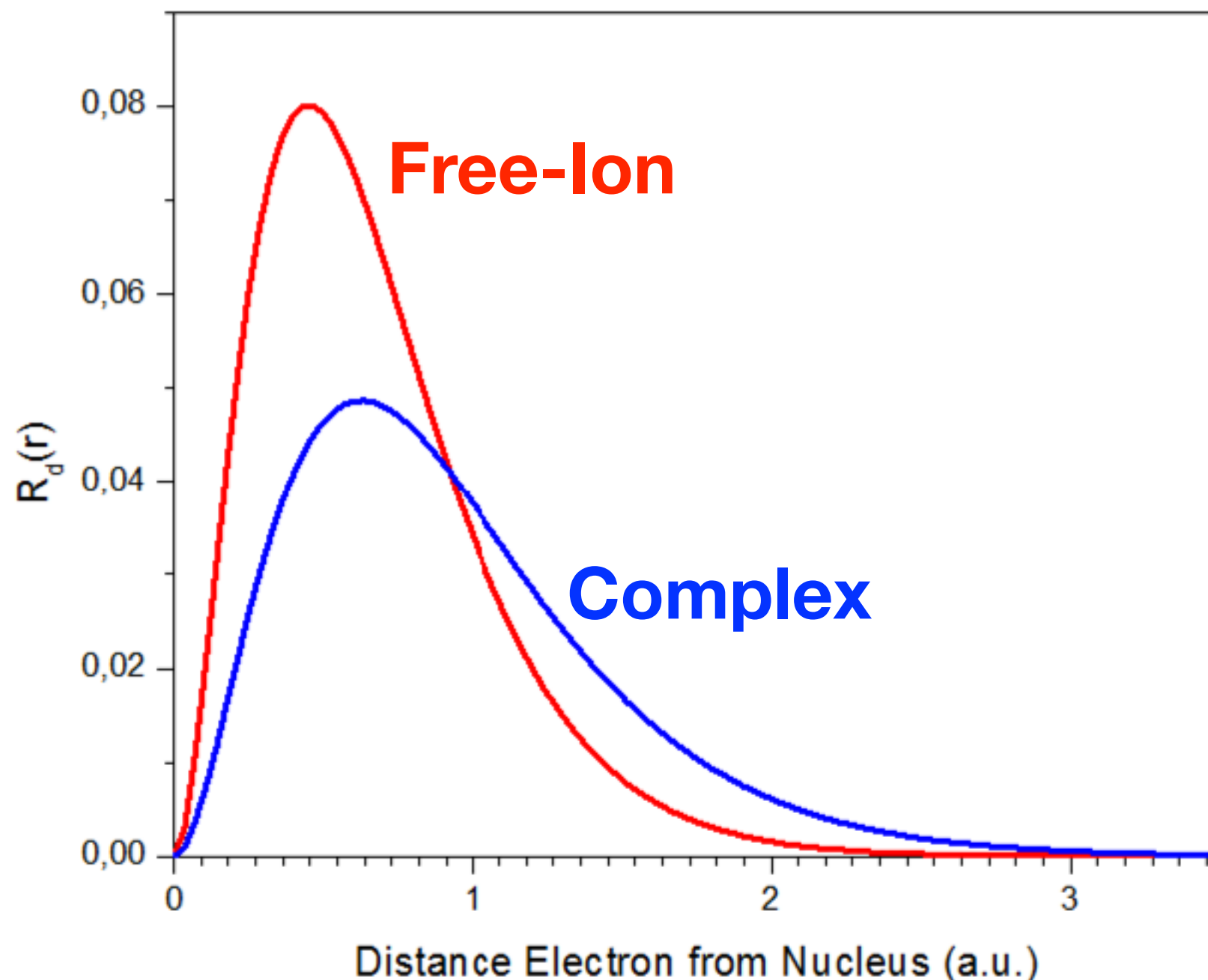
The Nephelauxetic Effects

- ➔ Observation 1: **Racah parameters** fitted for ions in complexes are smaller than those obtained for free-ions from atomic spectroscopy (**nephelauxetic effect**)
- ➔ Observation 2: **Spin-Orbit coupling parameters** fitted for ions in complexes are smaller than those obtained for free-ions from atomic spectroscopy (**relativistic nephelauxetic effect**)



Interpretation of the Nephelauxetic Effects

Traditional Interpretation: The radial wavefunction of the metal d-orbitals expand in the complex and hence B and ζ are reduced (*nephelauxetic=,cloud expanding'*)



Radial wavefunction

$$R_d(r) = \sqrt{\frac{(2\kappa)^7}{720}} r^2 \exp(-\kappa r)$$



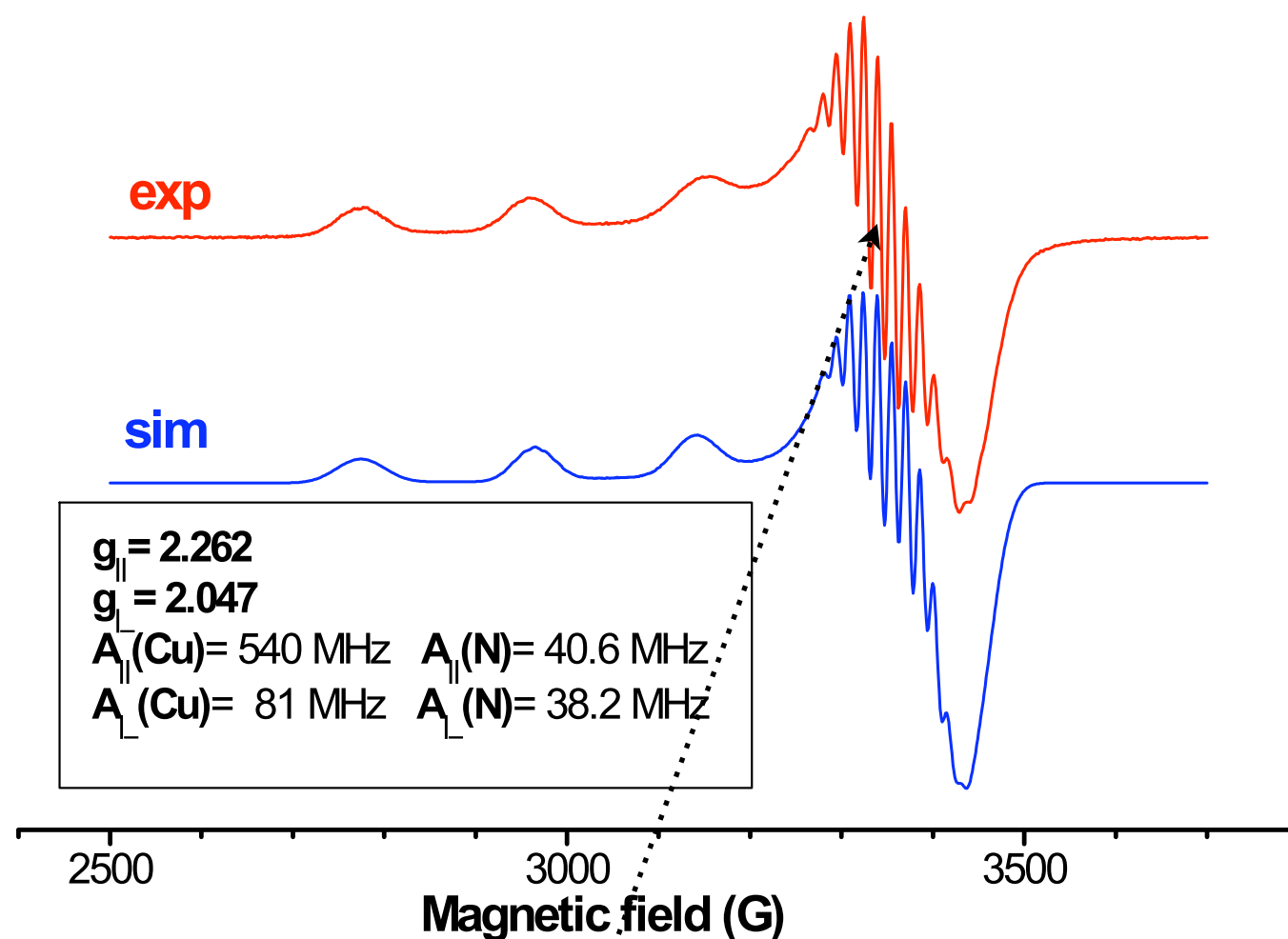
$$\zeta_M = \frac{1}{30} Z_M^{\text{eff}} c^{-2} \kappa^3$$

$$B = \frac{143}{80640} \kappa$$

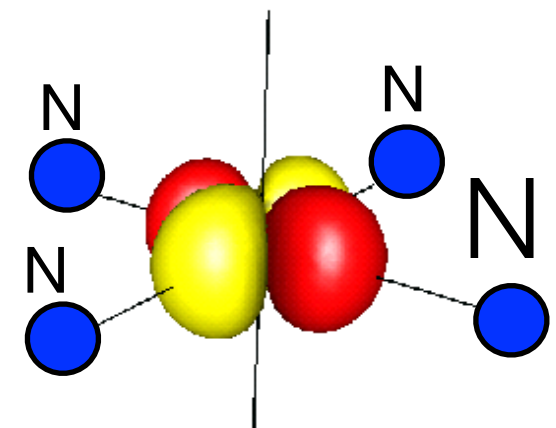
Experimental Proof of the Inadequacy of LFT



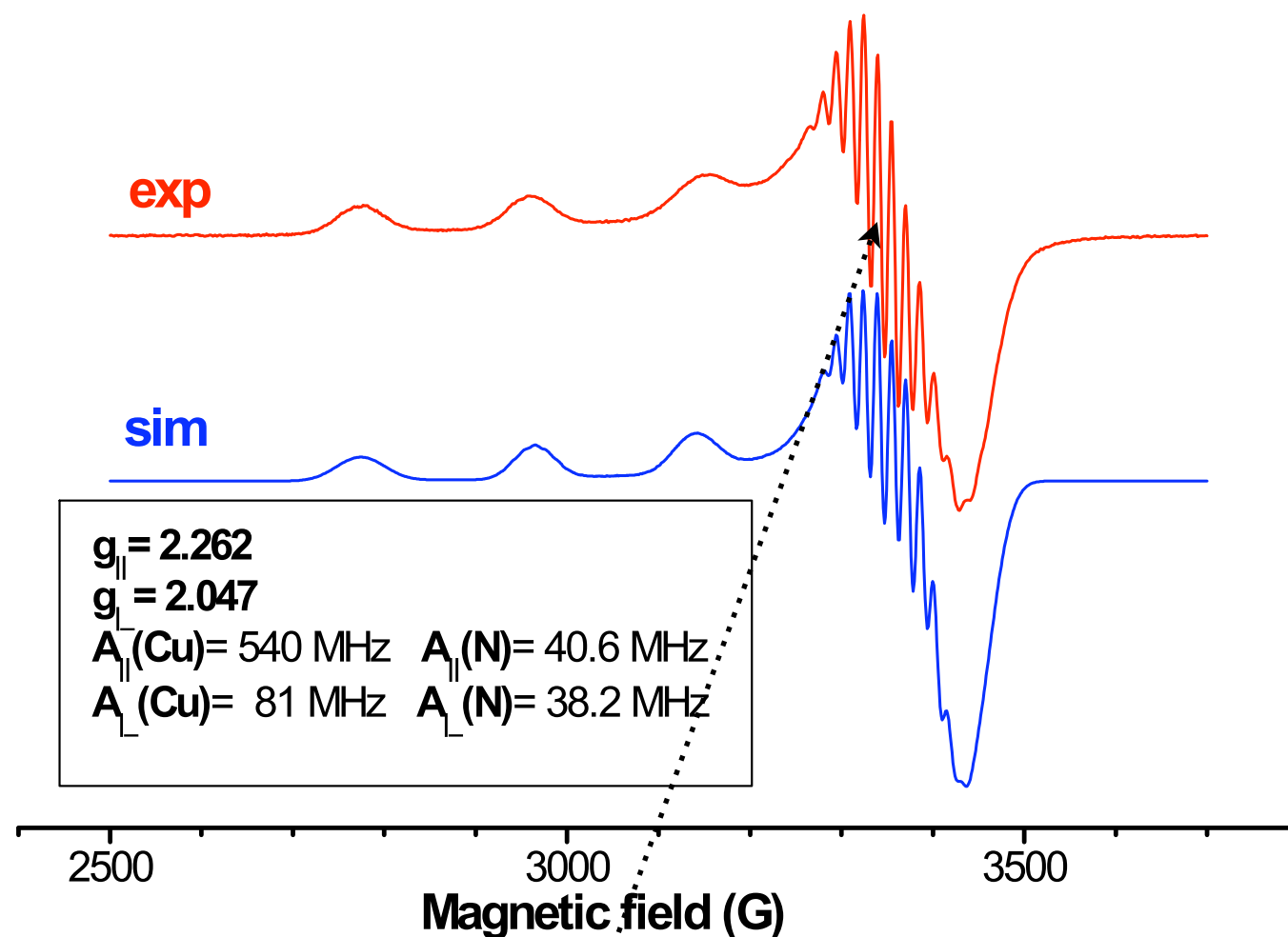
Ligand Field Picture



Clearly Observed Coupling Between
The Unpaired Electron and the Nuclear Spin
of Four ^{14}N Nitrogens ($I=1$)

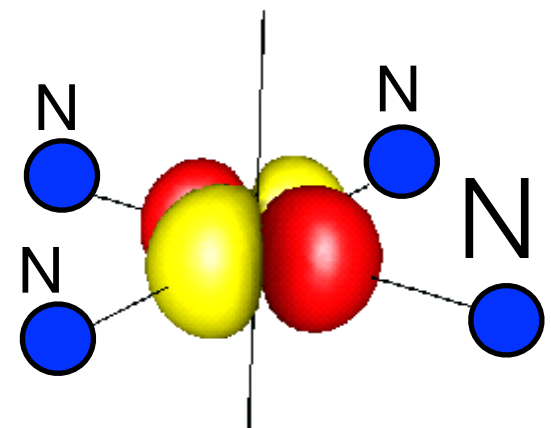


Experimental Proof of the Inadequacy of LFT

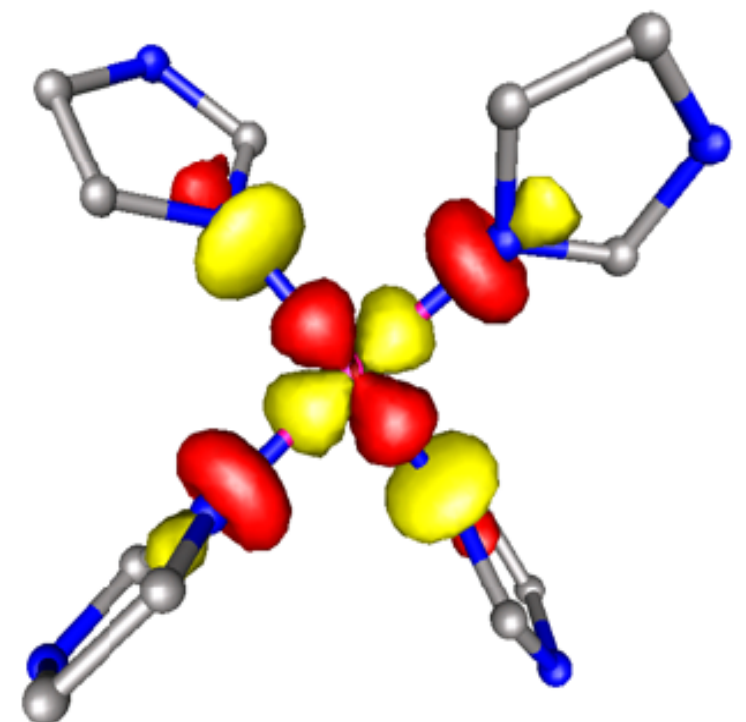


Clearly Observed Coupling Between
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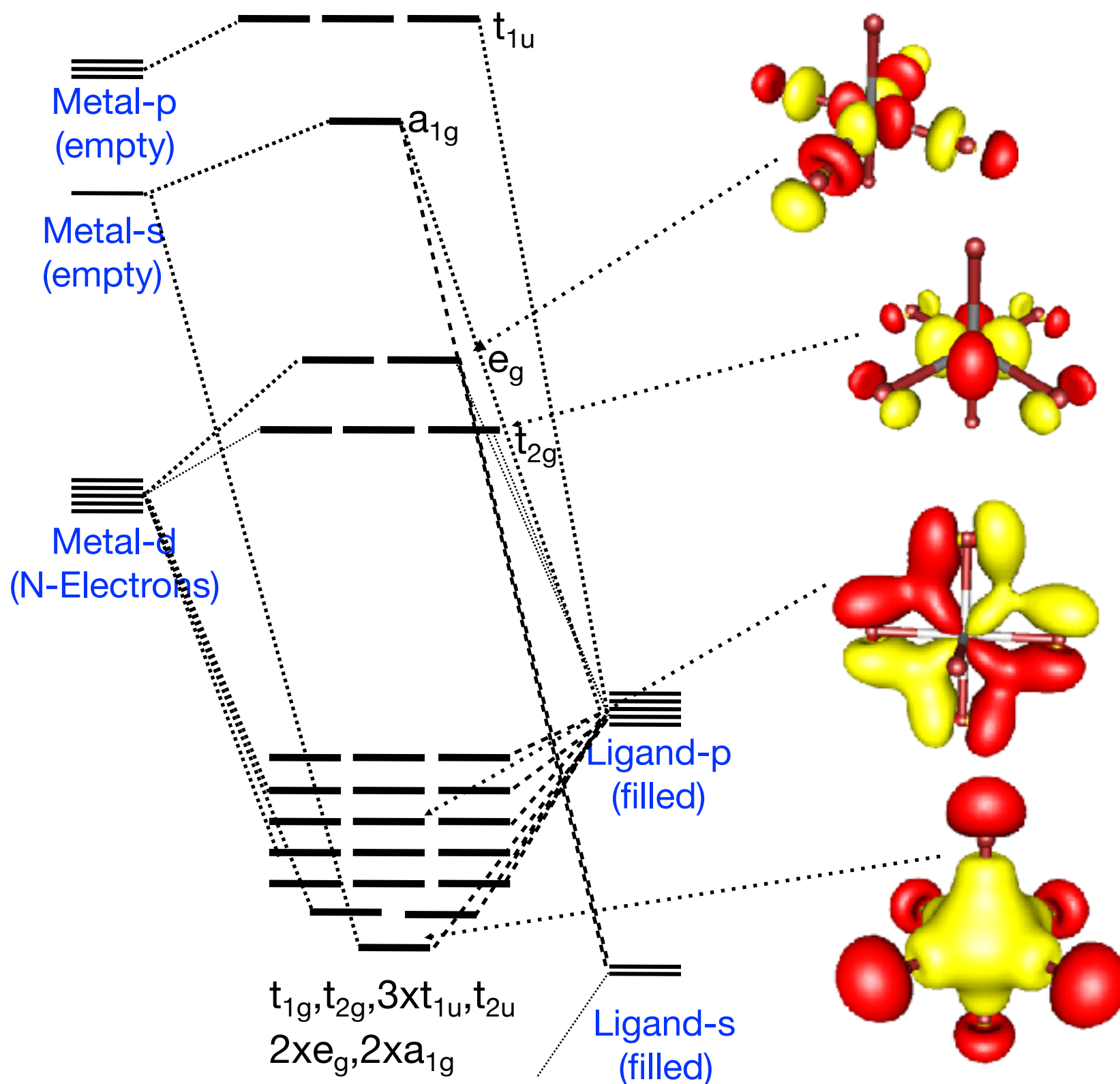
Ligand Field Picture



Molecular Orbital Picture



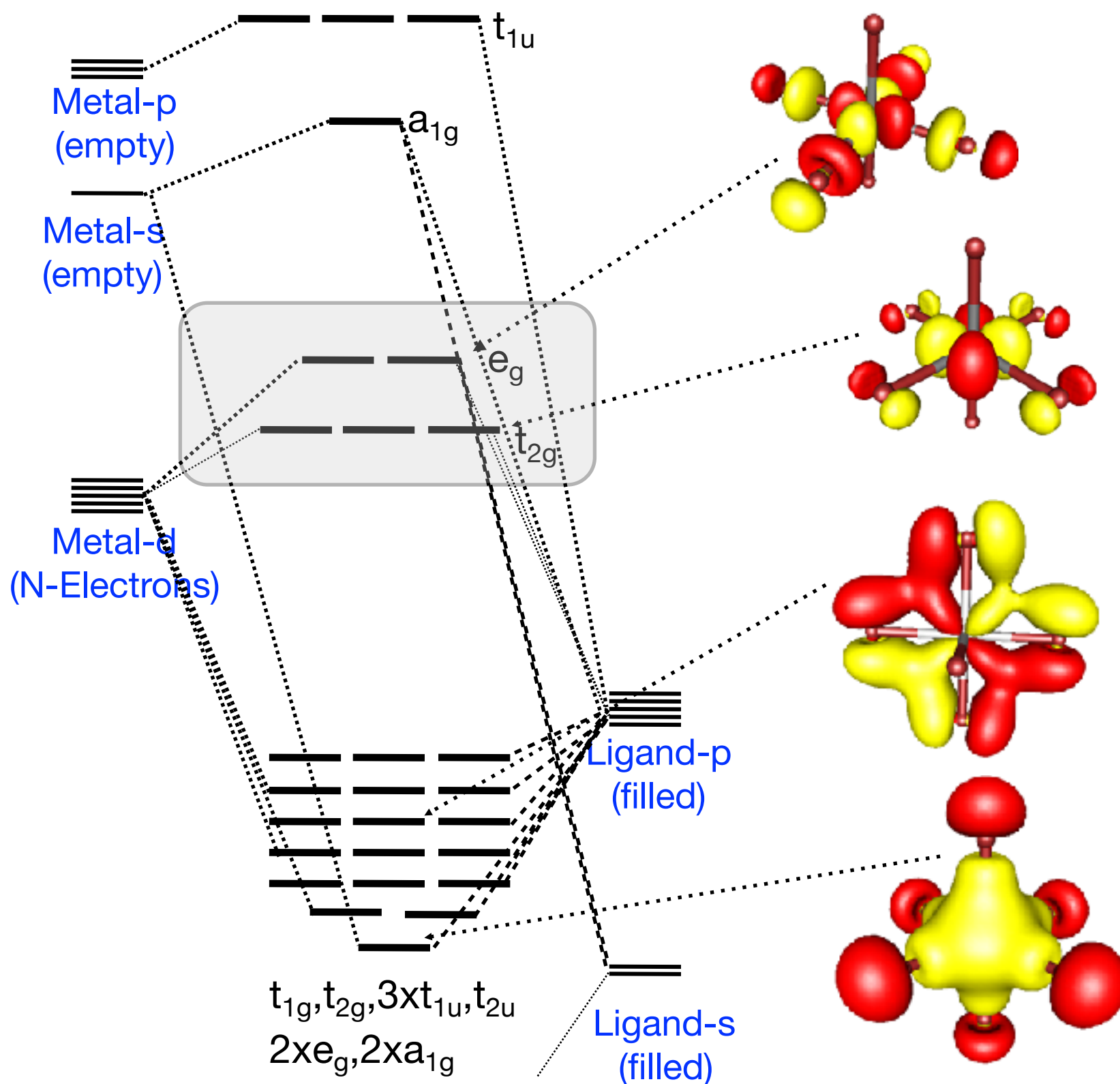
MO Theory of ML_6 Complexes



Key Points:

- ▶ The orbitals that are treated in **LFT** correspond to the **anti-bonding metal-based orbitals** in **MO Theory**
- ▶ Through bonding some **electron density** is transferred from the ligand to the metal
- ▶ The extent to which this takes place defines the **covalency** of the M-L bond

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Alright then: Now we understand

„Covalency is just the dilution of the d-orbitals with ligand orbitals and the founding fathers where wrong about the cloud expansion“

$$\psi_d \cong \alpha |M_d\rangle - \sqrt{1 - \alpha^2} |L\rangle$$

Aha effect: The (in)famous ‚Stevens reduction factor‘ is just the d-character in the respective MO squared (α^2)

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„Measurements“ of Covalency?

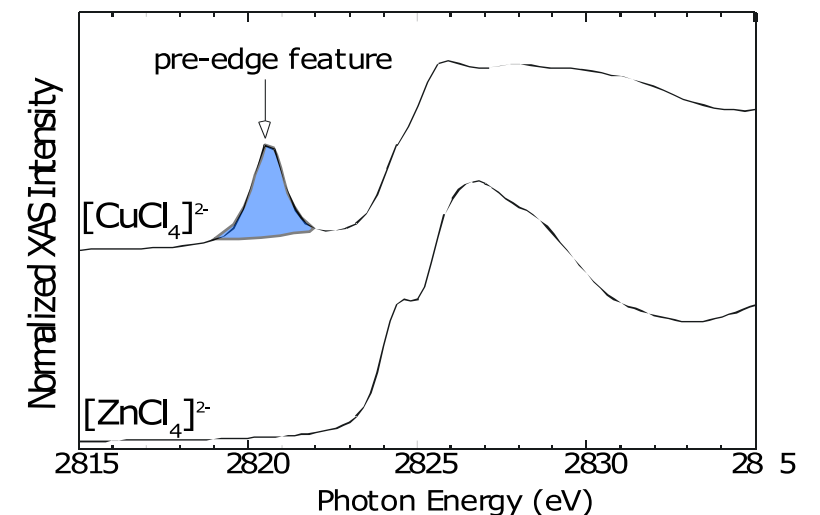
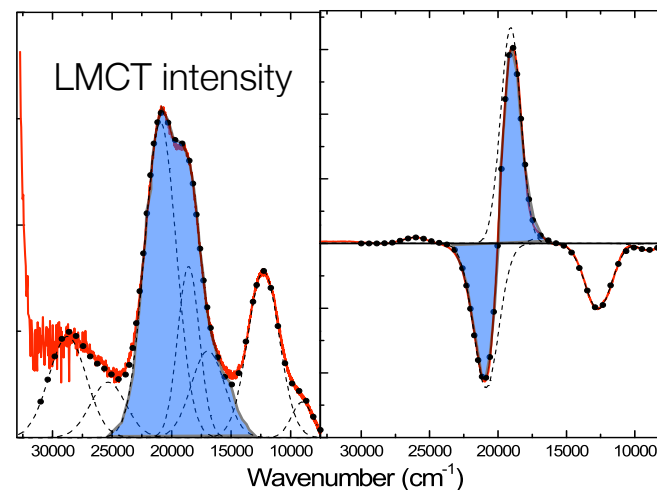
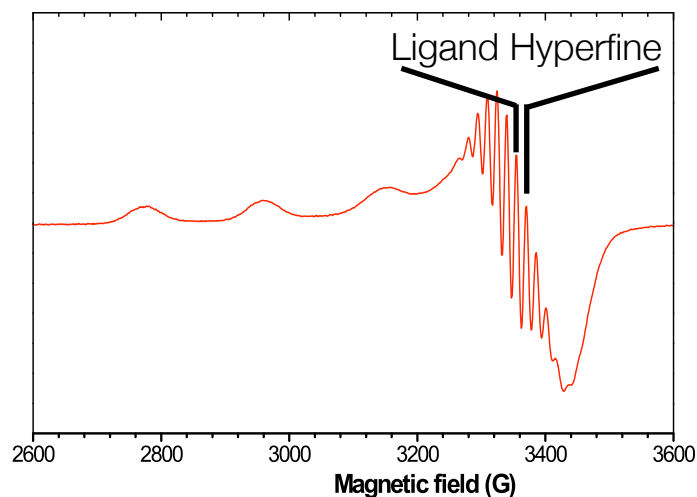
★ Can covalency be measured?

➔ Rigorously speaking: **NO!** Orbitals are not observables!

➔ On a practical level: (more or less) **YES**. Covalency can be correlated with a number of spectroscopic properties

- ▶ **EPR metal- and ligand hyperfine couplings**
- ▶ **Ligand K-edge intensities**
- ▶ **Ligand-to-metal charge transfer intensities**

➔ As all of this is „semi-qualitative“ you can not expect numbers that come out of such an analysis to agree perfectly well. If they do this means that you have probably been good at fudging!

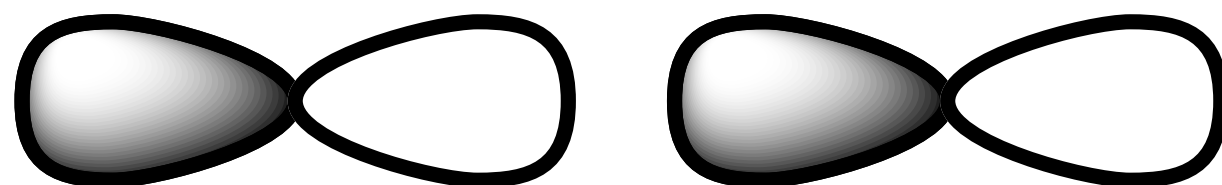


ALL proportional to $1-\alpha^2$

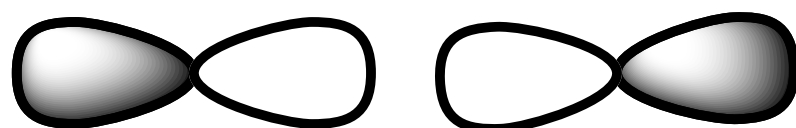
... but that is still too simple ...

Rüdenberg's Analysis of the Chemical Bond

In order to maintain the balance between potential and kinetic energy required by the virial theorem ($-2T=V$) orbitals have to expand or contract upon bond formation



Antibonding orbital: **expansion**



Bonding orbital: **contraction**

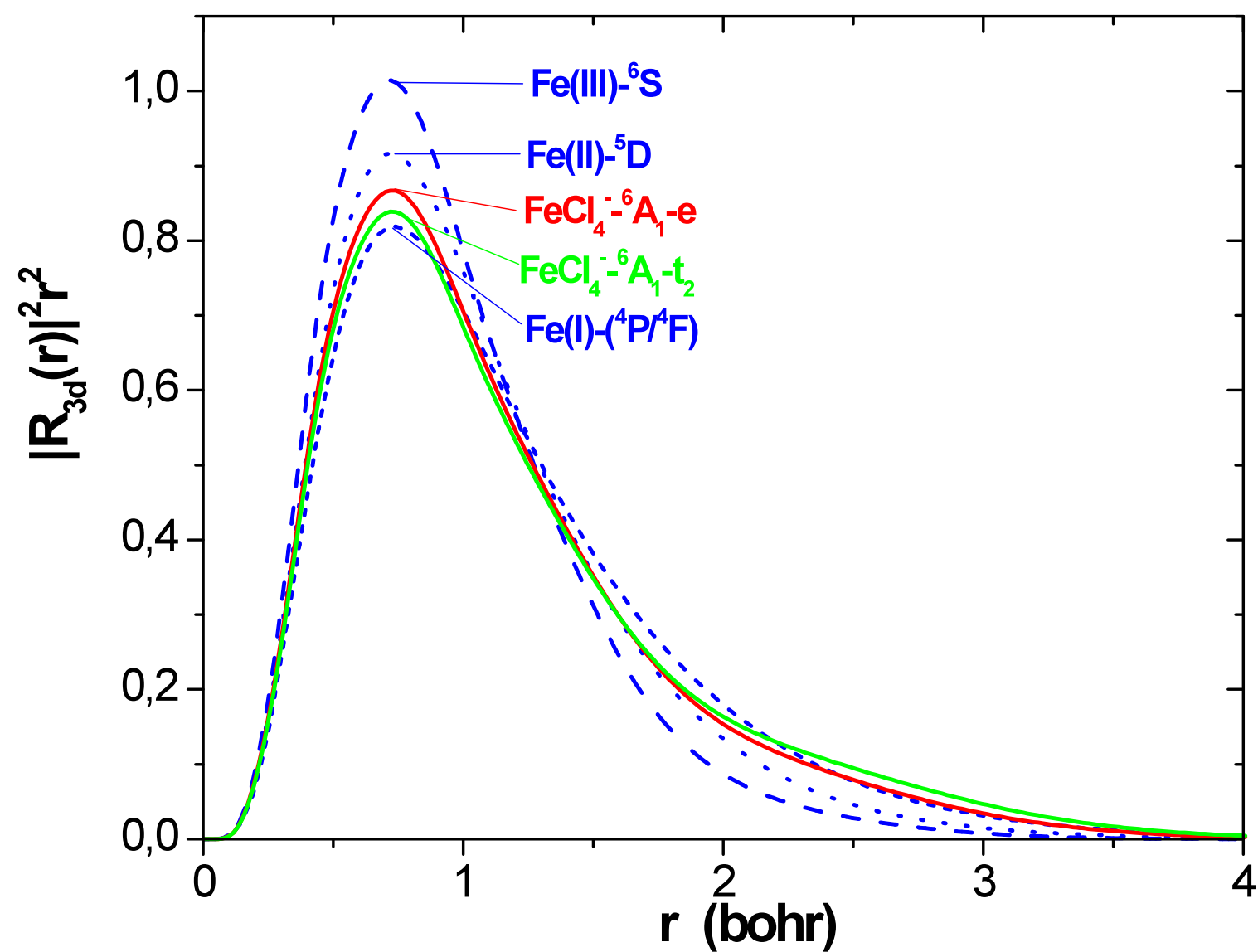


Klaus Rüdenberg
(1920-)

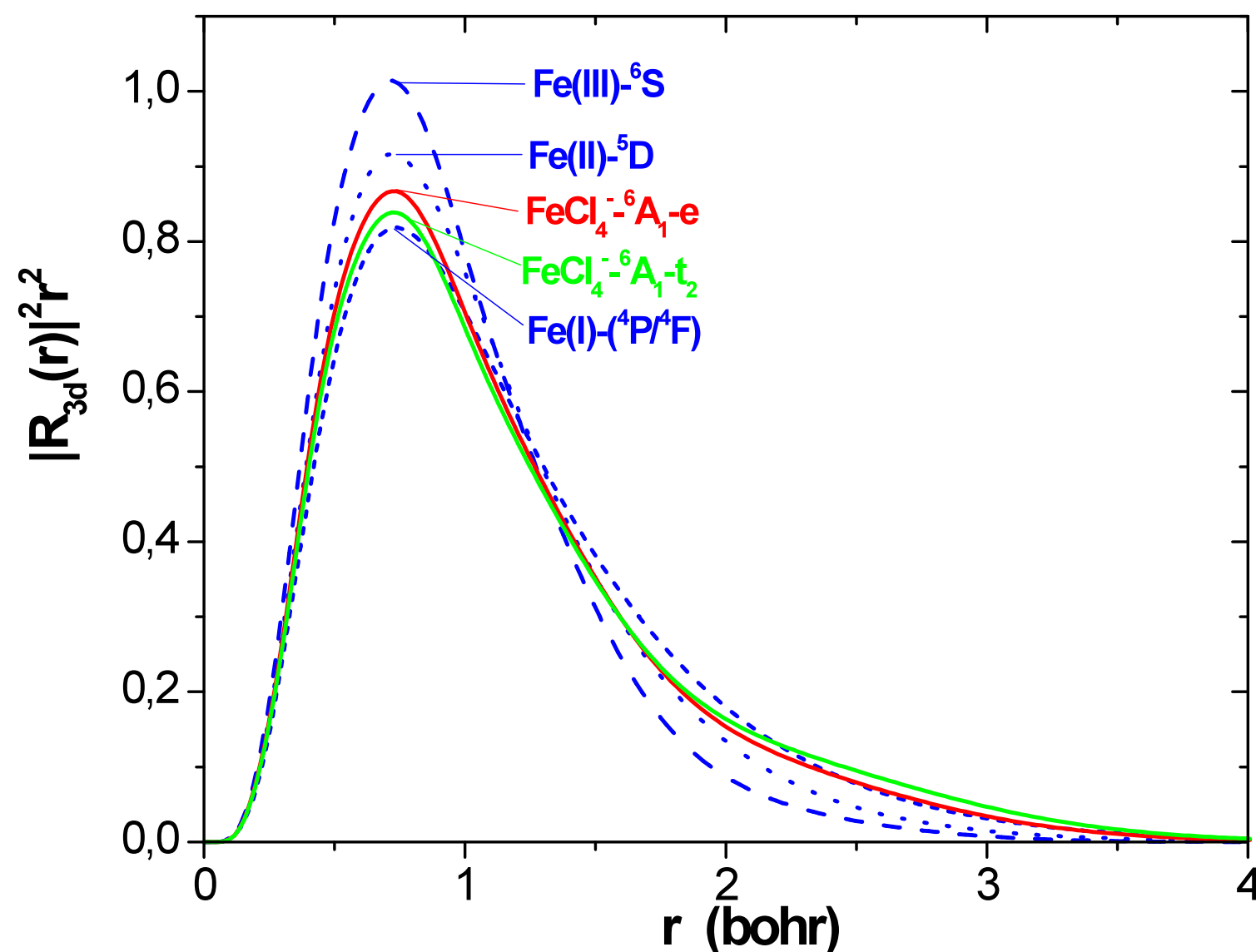
➔ Metal d-based molecular orbitals are anti bonding - hence there is radial expansion!

a) Rüdenberg, K. *Rev. Mod. Phys.* **1962**, 34, 326 b) Feinberg, M. J.; Rüdenberg, K.; Mehler, E. L. *Adv. Quantum Chem.* **1970**, 5, 27 c) Feinberg, M. J.; Rüdenberg, K. *J. Chem. Phys.* **1971**, 59, 1495.

Central Field Covalency

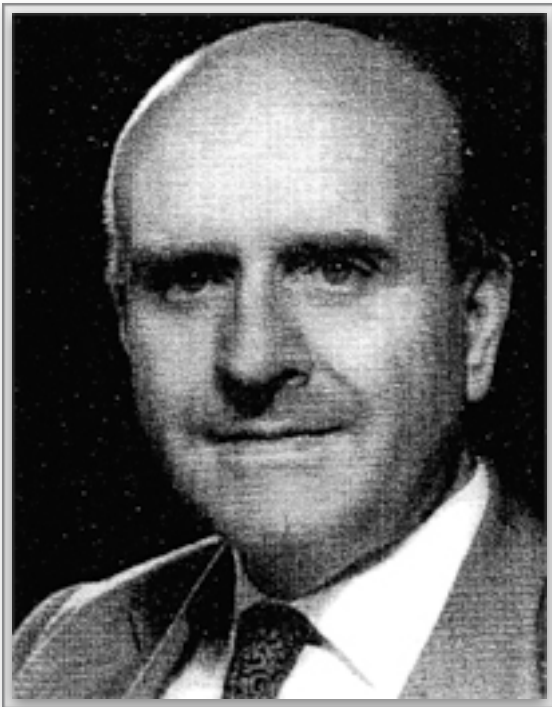


Central Field Covalency



- ▶ The founding fathers of ligand field were partially right by assuming cloud expansion
- ▶ However: the magnitude of the effect was grossly overestimated
- ▶ The explanation given (electrostatic effects, effective nuclear charge) was wrong

Bottom line up to this point:



Christian Klixbüll Jörgensen
(1931-2001)

There are 2 types of covalency:

1) Symmetry restricted covalency

Metal d-orbital dilution

(entirely correct)

2) ,Central field' covalency

Metal radial expansion

(effect correct but mechanism incorrect)

... but is that the whole picture then?

... but is that the whole picture then?

... We need to be able to calculate ligand field parameters ($10Dq$, B , C) from accurate first principles electronic structure theory!

Unambiguous Match between NEVPT2 and LFT

Overwhelming importance:

There is a 1:1 correspondance between the ligand field CSFs and the CAS-CI CSFs.

$$\Theta_I^{LFT} = | d_{xy}^\alpha d_{xz}^\beta \dots d_{z^2}^\alpha |$$

Ligand field pure d-orbital

$$\Theta_I^{CASSCF} = | \psi_{xy}^\alpha \psi_{xz}^\beta \dots \psi_{z^2}^\alpha |$$

Ab initio molecular orbital with metal d-parentage

Thus, all we have to ensure is that ligand field d-orbitals and CASSCF molecular orbitals of the same parentage are ordered in the same way and that CSFs are constructed in the same way.

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The condition is then that the ligand field CI matrix should resemble the *ab initio* effective Hamiltonian as closely as possible

$$\left| H_{IJ}^{LFT} - H_{IJ}^{eff} \right| = \min$$

For each matrix element!

While this looks at first sight to be a nonlinear optimization problem, in reality things are easy **because the ligand field matrix is linear in each and every ligand field parameter!**

$$\mathbf{H}^{LFT}(e, B, C) = \mathbf{H}^{LFT}(0) + \frac{\partial \mathbf{H}^{LFT}}{\partial B} B + \frac{\partial \mathbf{H}^{LFT}}{\partial C} C + \sum_L \frac{\partial \mathbf{H}^{LFT}}{\partial e_L} e_L$$

This ensures that there is a **unique least squares solution** that provides the unambiguous best fit of the ligand field and effective Hamiltonian matrices:

$$\mathbf{A}\mathbf{p} = -\mathbf{b} \Leftrightarrow \mathbf{p} = -\mathbf{A}^{-1}\mathbf{b}$$

p_K = The k'th ligand field parameter

$$A_{KL} = \sum_{IJ} \frac{\partial H_{IJ}^{LFT}}{\partial p_K} \frac{\partial H_{IJ}^{LFT}}{\partial p_L}$$

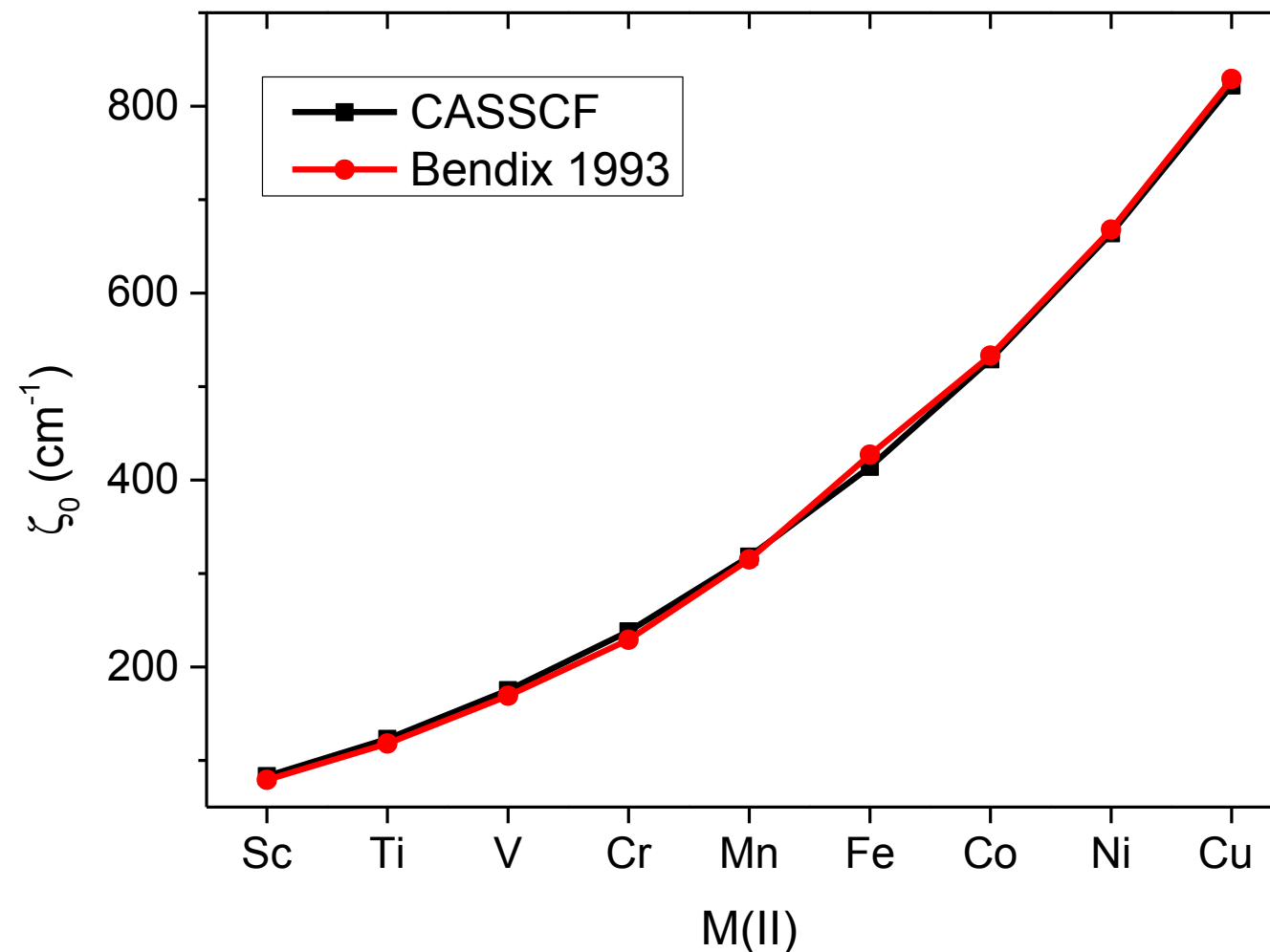
$$b_K = \sum_{IJ} \frac{\partial H_{IJ}^{LFT}}{\partial p_K} H_{IJ}^{eff}$$

This implies the strategy:

1. Choose your AOM scheme
2. Perform a QD-NEVPT2 calculation to obtain H^{eff}
3. Solve linear equation system to obtain the ligand field parameters

Revisiting Covalency from an *Ab Initio* Ligand Field Perspective

Results for The Spin-Orbit Coupling Constant



➔ Spin-Orbit Coupling parameters are *extremely* well reproduced, already at CASSCF

A model for the Relativistic Mephelauxetic Effect

For Molecules: Decomposition of the SOC constant into contributions from symmetry restricted and central field covalency

$$\zeta_{complex} = \zeta_{ion} \left(1 - \underbrace{\gamma_M^{SR}}_{\substack{\text{covalent} \\ \text{dilution} \\ \text{proportional to } \alpha^2}} - \underbrace{\gamma_{CF}}_{\substack{\text{Central} \\ \text{Field} \\ \text{proportional to } \kappa^3}} - \underbrace{\gamma_L}_{\substack{\text{Ligand} \\ \text{SOC part}}} \right)$$

$$h_{SOC} \propto \frac{1}{c^2} \sum_A \frac{Z_A^{eff}}{|\mathbf{r} - \mathbf{R}_A|^3} \mathbf{l}^{(A)} \mathbf{s} \quad \text{Hence} \quad \langle \psi_d | h_{SOC} | \psi_{d'} \rangle \approx \sum_A \alpha_A^2 \zeta^{(A)} \langle \psi_d | \mathbf{l}^{(A)} \mathbf{s} | \psi_{d'} \rangle$$

A very local operator! To a good approximation sum of quasi-atomic contributions weighted by covalent dilution factors

Very similar to an anisotropic version of the ,classical‘ Steven’s reduction factors

Results for Molecules

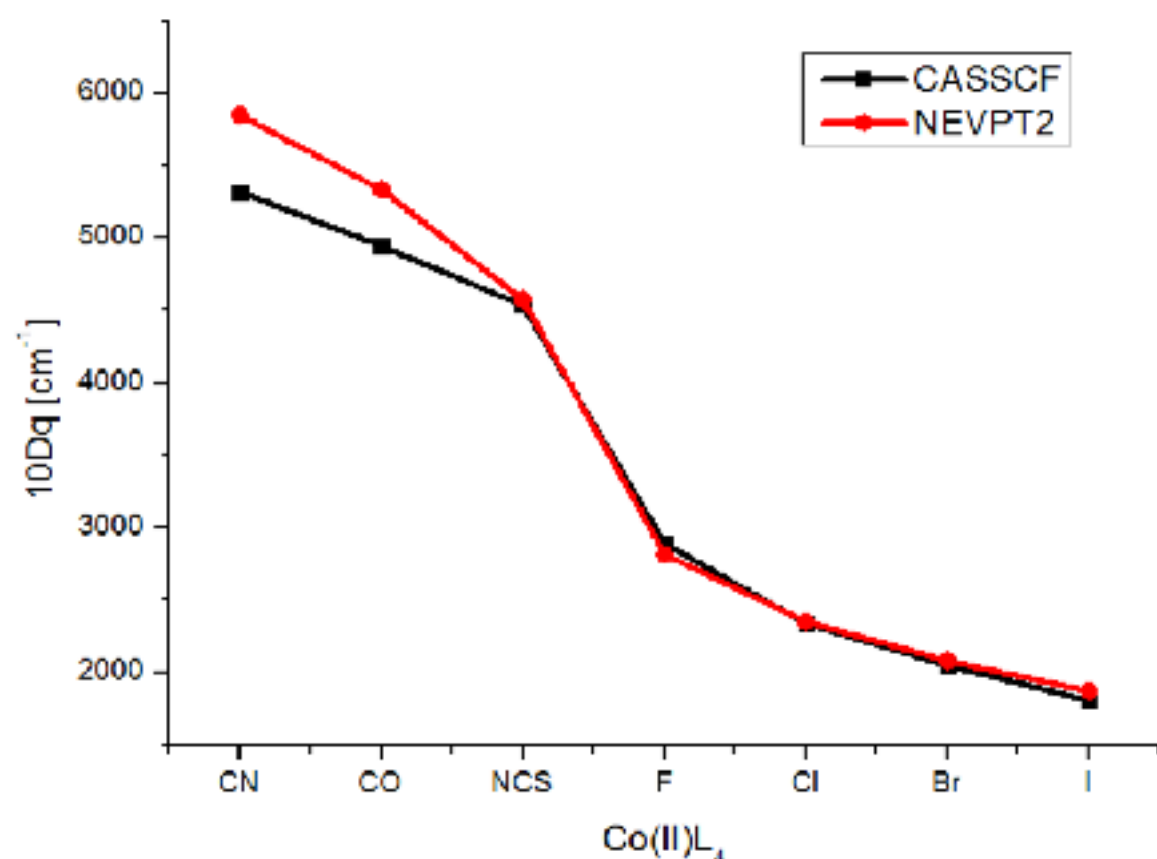
In silico Studies:



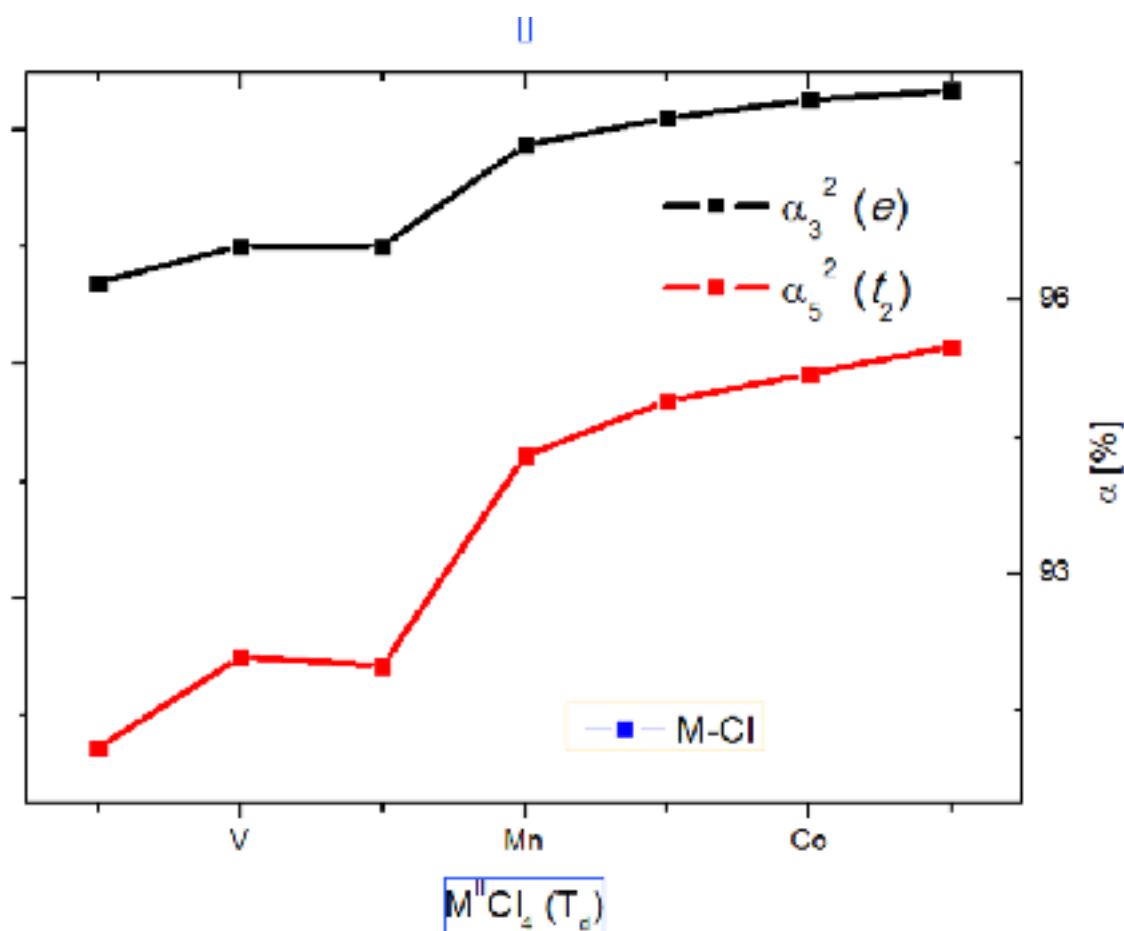
Metal-variation



Ligand-variation



➔ Expected trends in Ligand Fields

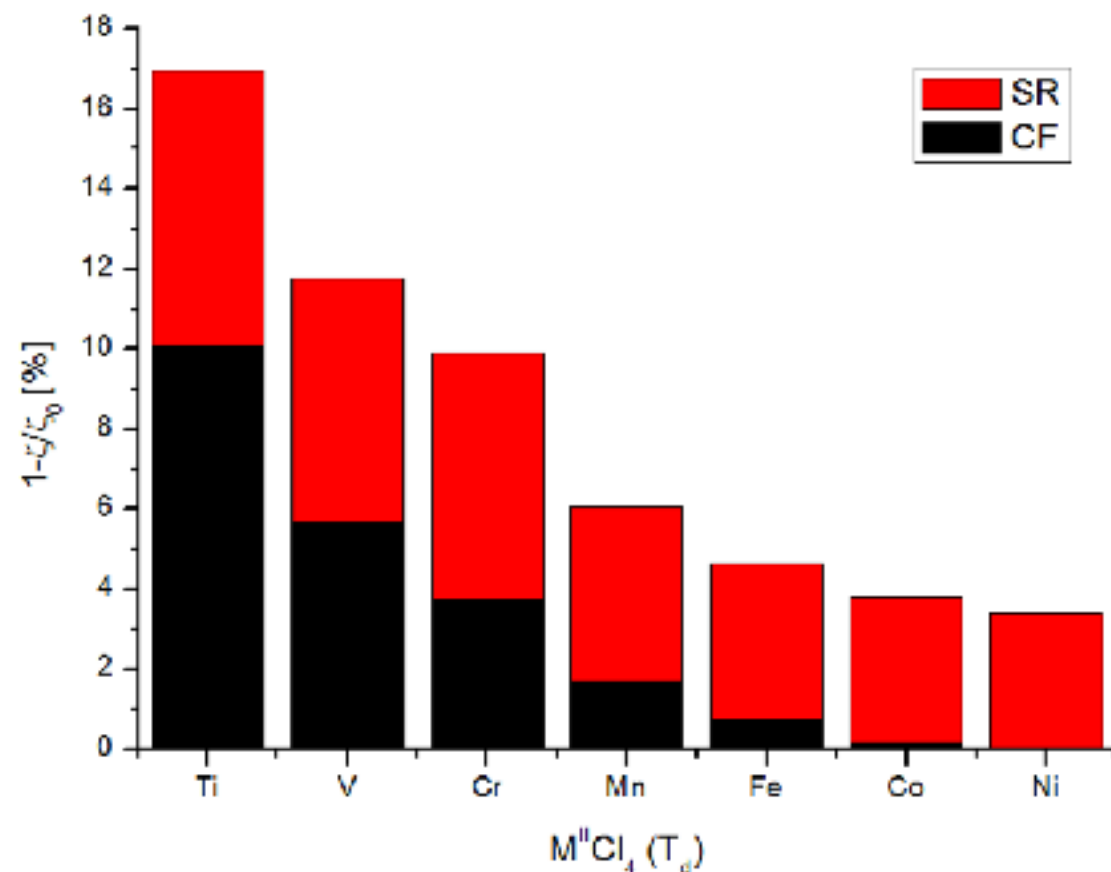


➔ Expected trends in symmetry restricted covalency

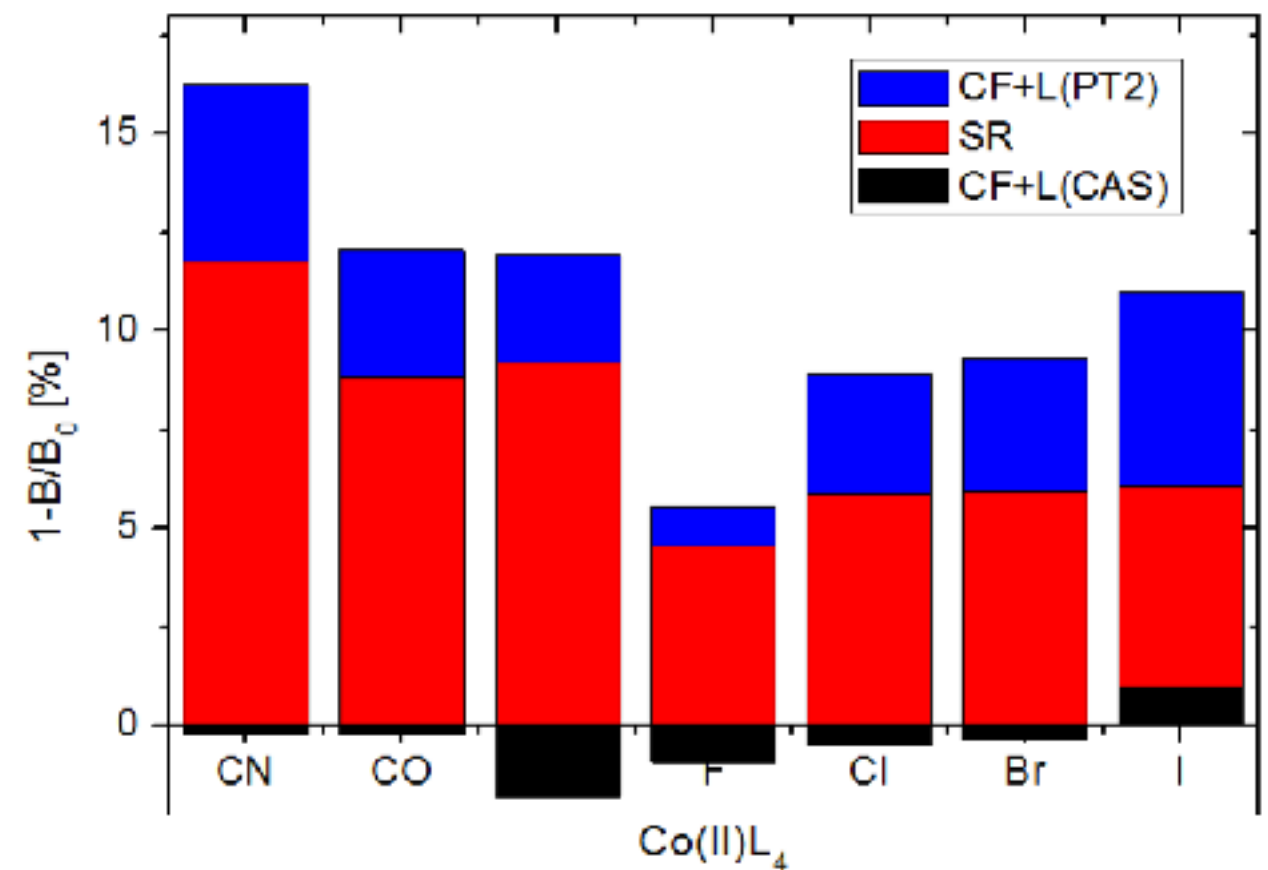
Everything chemically reasonable: Now analyze

The Relativistic Nephelauxetic Effect: Results

$$\gamma_{SR} = 1 - \alpha^2 \quad \gamma_{CF} = 1 - \frac{\zeta_M}{\alpha^2 \zeta_{ion}} \quad \gamma_L = 1 - \frac{\zeta_{complex} - \zeta_M}{\zeta_{ion}}$$



→ For late transition metals symmetry restricted covalency dominates, for early, central field covalency



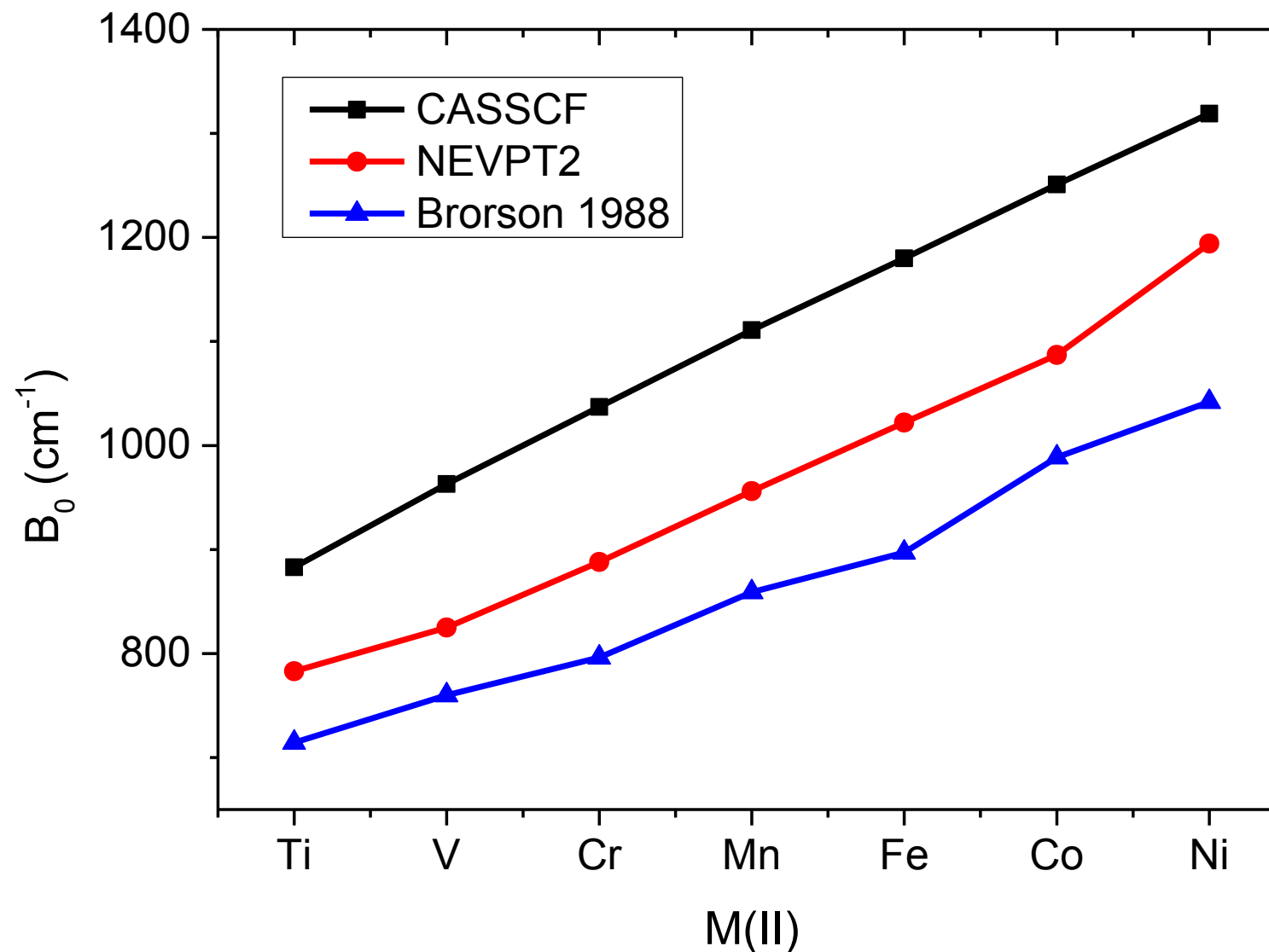
→ Ligand contribution only sizable for heavy ligand; typically negative

Bottom Line:

The classical picture of interplay between central field and symmetry restricted covalency works well for the interpretation of the relativistic nephelauxetic effect



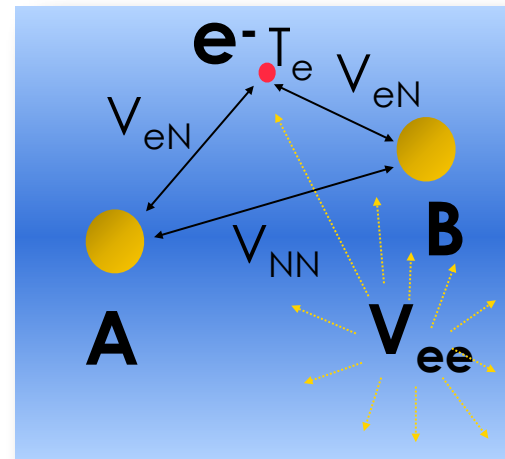
AILFT Results for B-Parameters of Free Ions



- ➔ Racah parameters are overestimated by CASSCF and only slightly overestimated by NEVPT2
- ➔ **DYNAMIC ELECTRON CORRELATION**

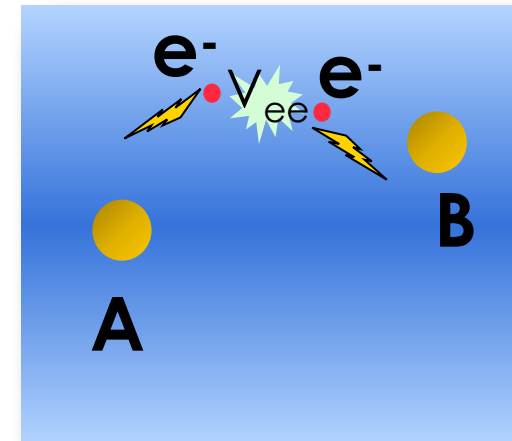
What is Dynamic Electron Correlation?

Exact Energy =



“Mean Field”
Hartree-Fock

+



Instantaneous electron-
electron interaction

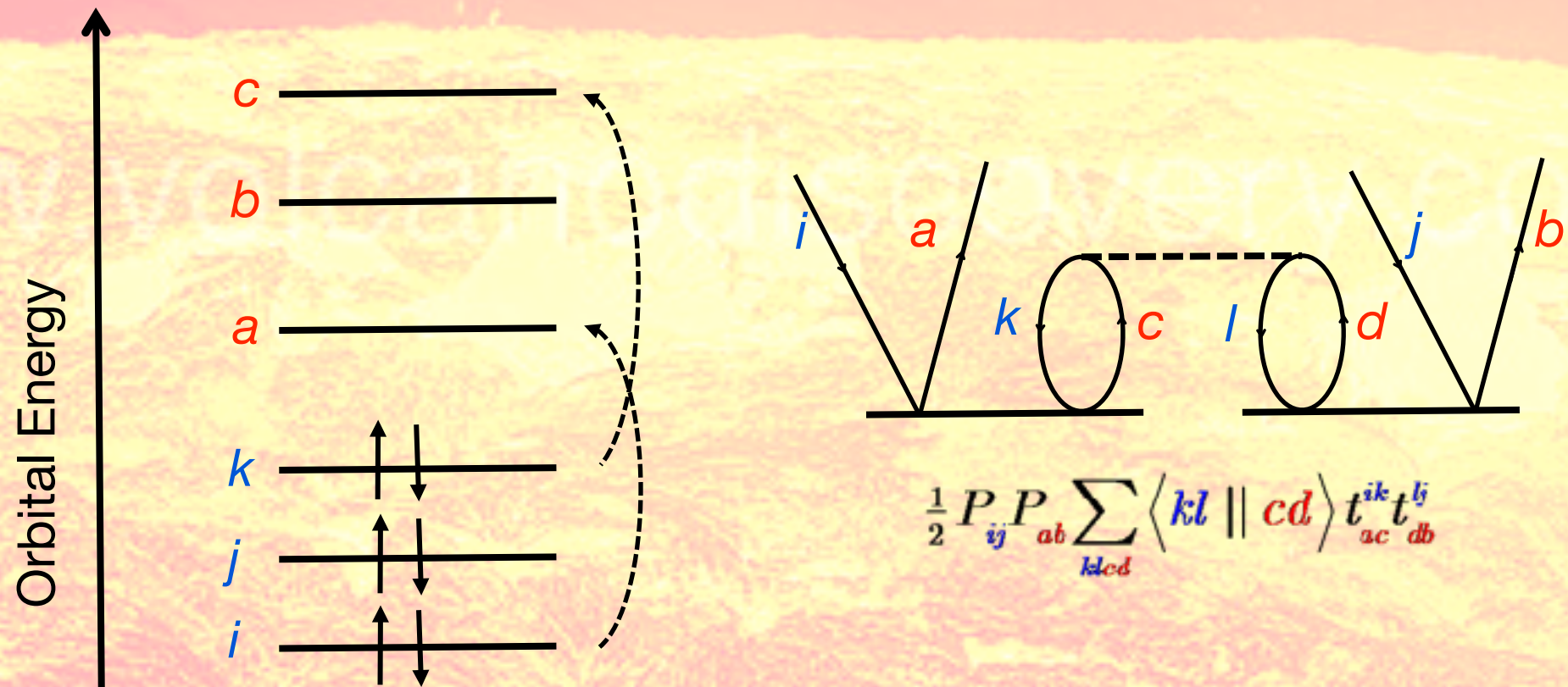
$$\text{Correlation energy} = \frac{1}{2} \sum_{i,j \text{ Electron pairs}} \mathcal{E}_{ij}(\uparrow\uparrow) + \mathcal{E}_{ij}(\uparrow\downarrow)$$

Fermi-Correlation

Coulomb-correlation

*Relatively easy due to
“Fermi hole” in the
mean-field*

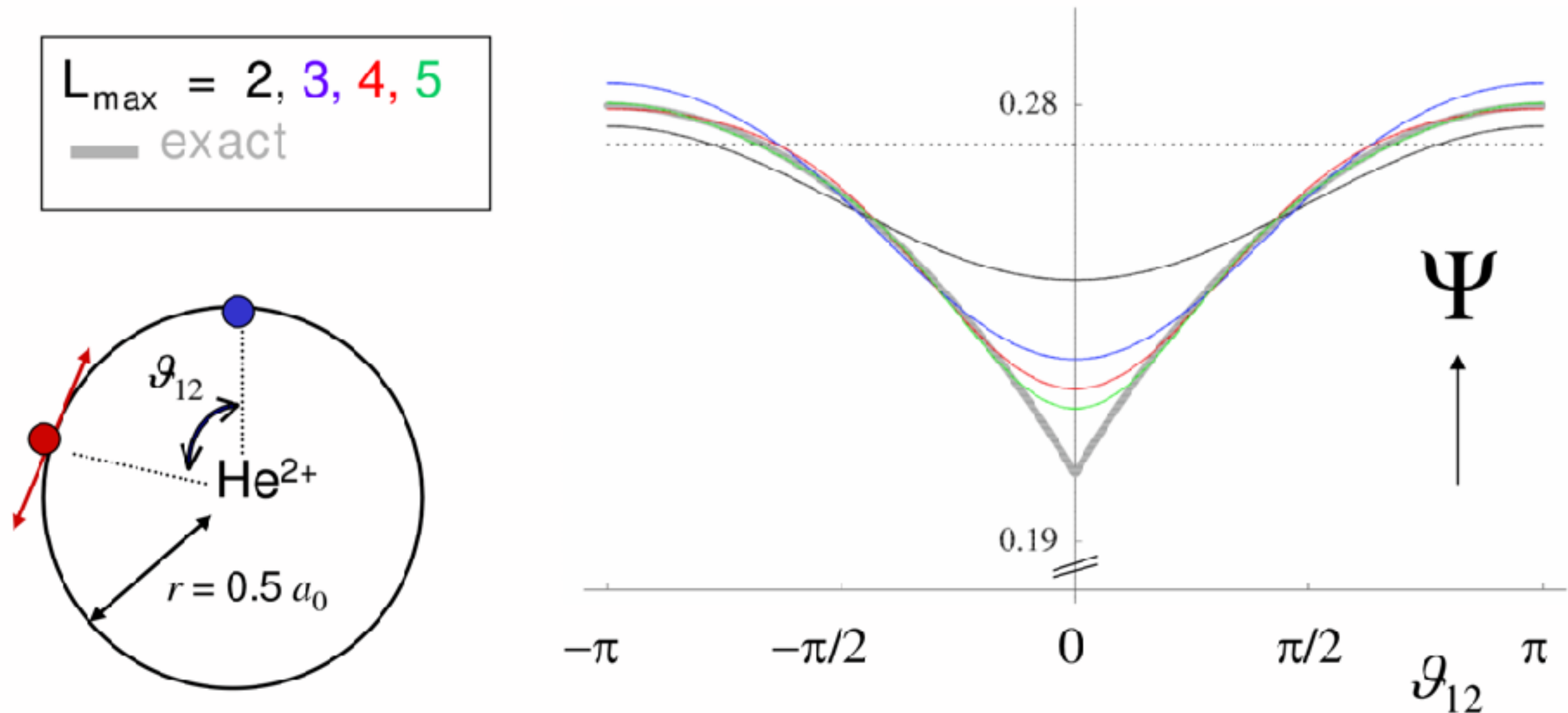
*Extremely hard to
calculate due to
interelectronic cusp at
the coalescence point
 $\mathbf{r}_1 = \mathbf{r}_2$*



$$\frac{1}{2} P_{ij} P_{ab} \sum_{klcd} \langle kl || cd \rangle t_{ac}^{ik} t_{db}^{lj}$$

$$|\Psi\rangle = |\Psi_0\rangle + \sum_{ia} C_a^i |\Psi_i^a\rangle + \frac{1}{4} \sum_{ijab} C_{ab}^{ij} |\Psi_{ij}^{ab}\rangle + \frac{1}{36} \sum_{ijkabc} C_{abc}^{ijk} |\Psi_{ijk}^{abc}\rangle + \dots$$

Another View: The Electron-Electron Cusp



$$\Psi(r_{12}) = \Psi(r_{12}=0)(1 + 0.5r_{12} + \dots)$$

Racah B Parameters for Complexes

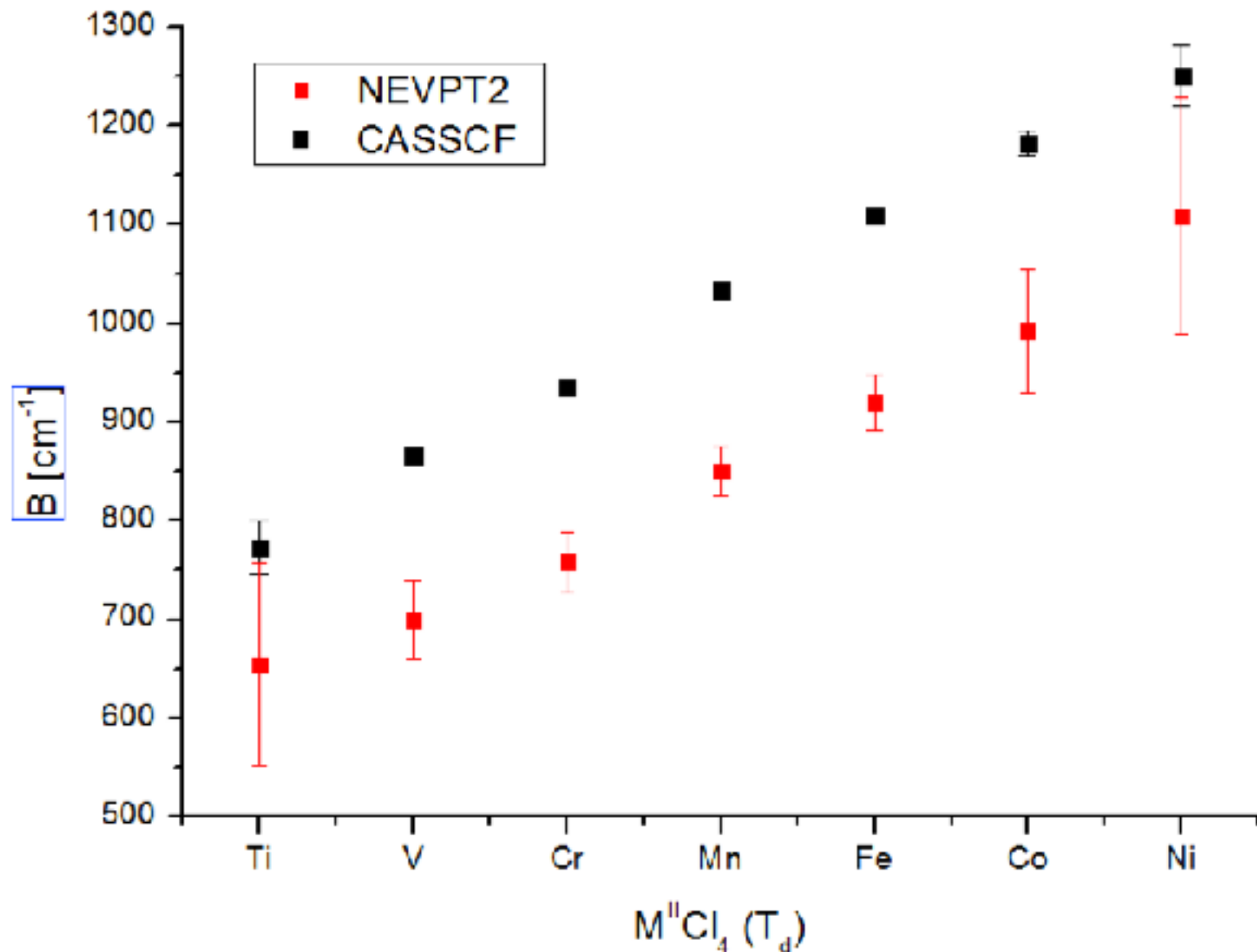
In Silico Studies:



Metal-variation



Ligand-variation



- ➔ Expected trends in Racah B -parameters
- ➔ Similar **LARGE** effects of dynamic correlation (10-30% of B!)

Attempted Decomposition of B

Try the equivalent decomposition

$$B_{complex} = B_{ion} (1 - \delta_M^{SR} - \delta_{CF} - \delta_L)$$

**covalent
dilution
prop. to α^4**

**Central
Field
prop. to κ**

*Ligand
repulsion part*

**Compared
to α^2 for ζ**

**Compared
to κ^3 for ζ**

- ➔ **There is no intrinsic locality in B**
- ➔ **B can NOT be decomposed as ζ**

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**Compared
to α^2 for ζ**

**Central
Field
prop. to κ**

**Compared
to κ^3 for ζ**

*Ligand
repulsion part*

**NOT small
NOT local**

- ➔ **There is no intrinsic locality in B**
- ➔ **B can NOT be decomposed as ζ**

Effect of Dynamic Correlation on B

Gedankenexperiment:

We can nicely decompose ζ .

- ✓ Hence, we can use the value of ζ_M to deduce an effective orbital exponent κ
- ✓ We can then use this exponent to calculate B
- ✓ This would give us the central field effect on B

However:

- ➔ Because of dynamic correlation this value of B is MUCH too large!
- ➔ Which exponent would we need to reproduce the molecular B?

Example: $[\text{Co}^{\text{II}}(\text{SR})_4]^{2-}$

From ζ : $\kappa = 4.46$

From B: $\kappa = 2.78$

Compared to $\kappa=3.95$ for a free ion
(Clementi)

**MUCH too diffuse
effective treatment of dynamic correlation!**

Bottom Line:

The analysis of the covalent dilution and radial distortion is missing two key contributions to B:

- (1) DYNAMIC ELECTRON CORRELATION**
- (2) The fact that B is NOT LOCAL**

Summing up: What really is Covalency?

➔ **Covalency** is strictly a **ONE-ELECTRON PHENOMENON**:

Symmetry restricted covalency: Covalent ,dilution' of metal orbitals

Central Field covalency: Radial distortion due to BONDING

Both effects work differently for B and ζ .

	Symmetry	Central Field	Local
B	α^4	K_{eff}	NO
ζ	α^2	K_{eff}^3	YES

➔ **A single reduction factor does NOT cover both effects**

➔ B can **not** be understood without recourse to ,dynamic electron correlation', which is strictly a **TWO-ELECTRON PHENOMENON**. It is not part of covalency
(disclaimer: covalency as I understand it)

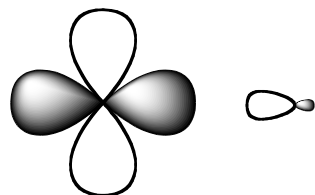
Measuring Covalency?

What is Covalency?

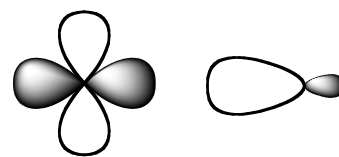
- ★ **Covalency refers to the ability of metal and ligand to share electrons** („soft“ concept with no rigorous definition)
- ★ Operationally, covalency can be defined in MO theory from the mixing coefficients of metal- and ligand orbitals

$$\psi_i \cong \alpha_i |M_i\rangle - \sqrt{1 - \alpha_i^2} |L_i\rangle \quad (\text{overlap neglected})$$

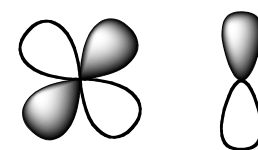
- ★ The value $1 - \alpha^2$ can be referred to as „the covalency“ of the specific metal ligand bond. It is the probability of finding the electron that occupies ψ_i at the ligand
 - ▶ *The maximal covalency is 0.5, e.g. complete electron sharing*
 - ▶ *The covalency might be different in σ - and π -bonds (e.g. it is anisotropic)*
 - ▶ *In σ -donor and π -donor bonds these are antibonding. The bonding counterparts are occupied and lower in energy*
 - ▶ *In π -acceptor bonds these orbitals are bonding. The antibonding counterparts are higher in energy and unoccupied*



Typical Ionic bond; hard ligands
Werner type complexes
 $\alpha^2 = 0.8 - 0.9$



Typical covalent bond
Organometallics; soft ligands
 $\alpha^2 = 0.5 - 0.8$

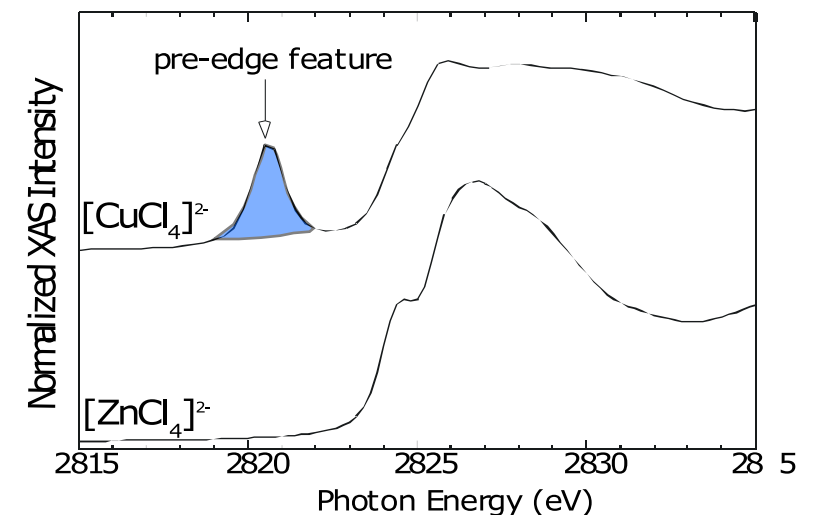
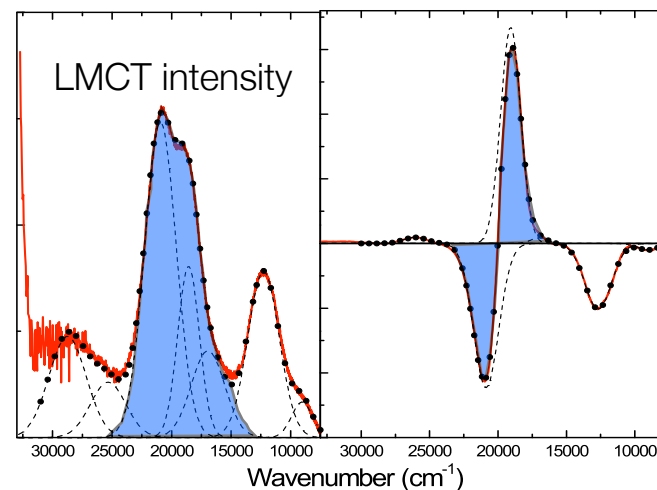
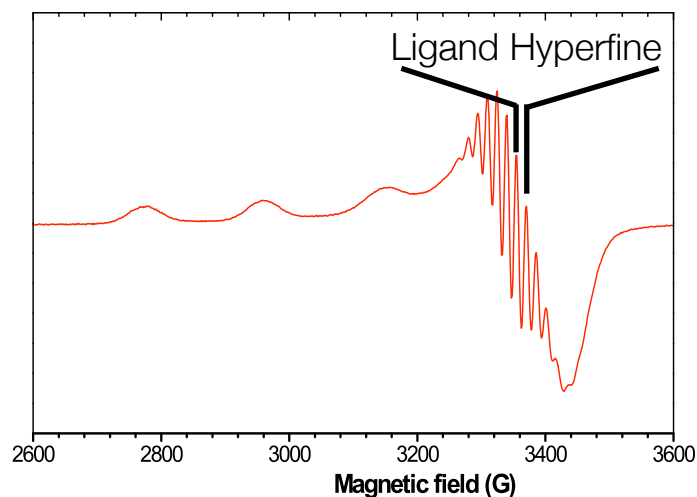


Typical π -backbond
Heterocyclic aromatic ligands ; CO, NO+...
 $\alpha^2 = 0.7 - 0.95$

„Measurements“ of Covalency?

★ Can covalency be measured?

- ➔ Rigorously speaking: **NO!** Orbitals are not observables!
- ➔ On a practical level: (more or less) **YES**. Covalency can be correlated with a number of spectroscopic properties
 - ▶ **EPR metal- and ligand hyperfine couplings**
 - ▶ **Ligand K-edge intensities**
 - ▶ **Ligand-to-metal charge transfer intensities**
- ➔ As all of this is „semi-qualitative“ you can not expect numbers that come out of such an analysis to agree perfectly well. If they do this means that you have probably been good at fudging!



ALL proportional to $1-\alpha^2$

Covalency and Molecular Properties

Metal-Ligand Covalency Affects Many Chemical Properties!

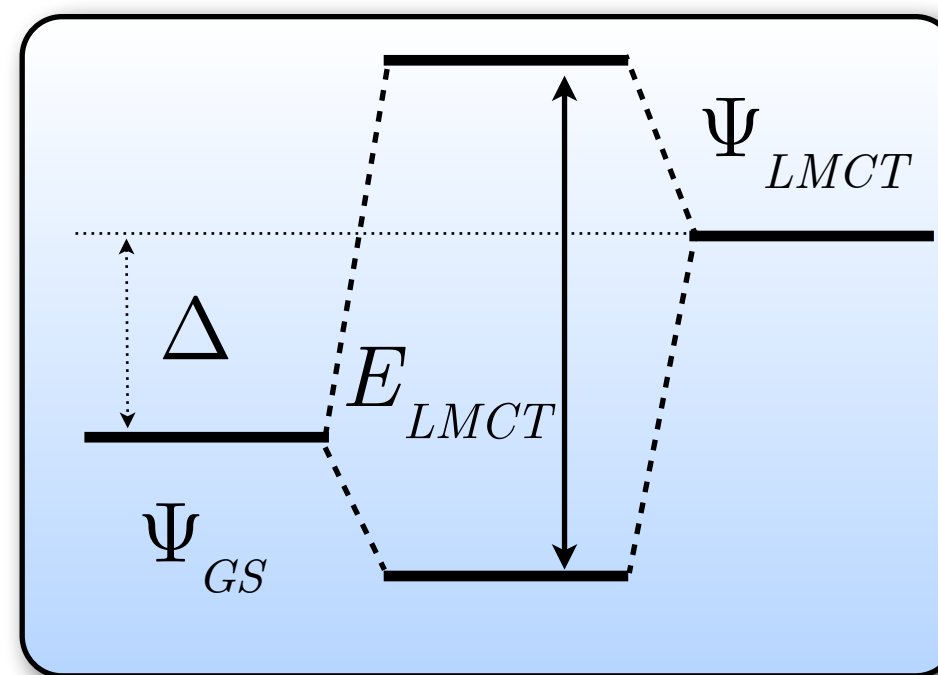
1. The **stability** of a complex increases with metal-ligand covalency
2. Covalency reflects **charge-donation**. The larger the charge donation the more negative the **redox potential**
3. Covalency may affect ,**electron transfer pathways**‘
4. Covalency taken to the extreme might mean that ligands are activated for **radical chemistry**
5. ...

Covalency and Ligand-to-Metal Charge Transfer Spectra

$$\Psi_{GS} = |\psi_L \bar{\psi}_L \psi_M| \quad \Psi_{LMCT} = |\psi_L \bar{\psi}_M \psi_M|$$

Energy Difference $\Delta = I_L - A_M$

Interaction $\beta = F_{LM} \propto S_{LM}$



Transition Energies:

- ★ Low if ligand is easy to ionize
- ★ Low if metal is strongly oxidizing (high oxidation state)
- ★ Increases for large ML overlap
- ★ Overlap increases for highly polarizable (soft) ligands

Transition Intensities:

- ★ High for large covalent binding (β =large, Δ =small)
- ★ Maximal for equal mixing (Δ =0)
- ★ Transitions are always most intense for bonding to antibonding excitations (polarized along the M-L bond)

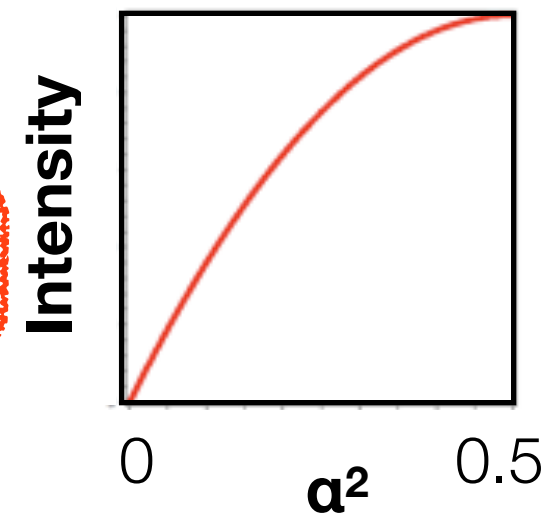
Estimating Covalency

From the little valence bond model, we can obtain the two eigenstates as:

➡
$$\begin{aligned} |\Psi'_{GS}\rangle &= \alpha |\Psi_{GS}\rangle + \sqrt{1 - \alpha^2} |\Psi_{LMCT}\rangle & |\Psi'_{LMCT}\rangle &= \sqrt{1 - \alpha^2} |\Psi_{GS}\rangle - \alpha |\Psi_{LMCT}\rangle \end{aligned}$$

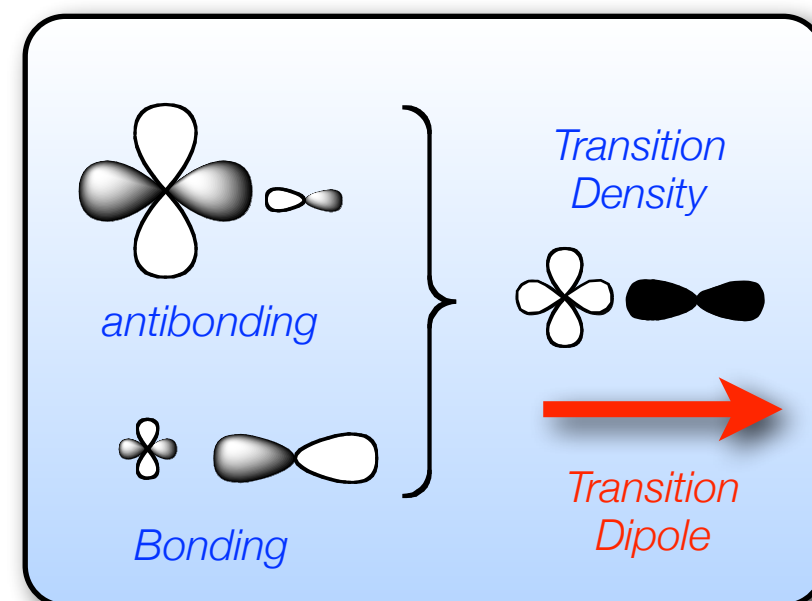
The **Transition Energy**: $E_{LMCT} = \Delta + 2\frac{\beta^2}{\Delta} - 2\frac{\beta^4}{\Delta^3} + O(\beta^6)$

The **Transition Intensity**: $\left| \langle \Psi'_{GS} | \vec{\mu} | \Psi'_{LMCT} \rangle \right|^2 \equiv D_{ML}^2 \approx \alpha^2(1 - \alpha^2)R_{ML}^2 \approx \frac{\beta^2}{\Delta^2} R_{ML}^2$



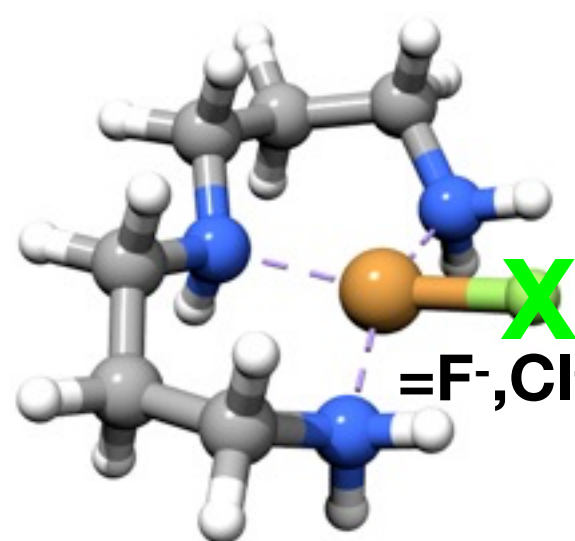
This can be turned around to obtain the model parameters from the measurable quantities: **R_{ML}** , **D_{ML}** and **E_{LMCT}**

$$\beta = \pm D_{ML} \frac{R_{ML}}{R_{ML}^2 + 2D_{ML}^2} E_{LMCT} \quad \Delta = \frac{R_{ML}^2}{R_{ML}^2 + 2D_{ML}^2} E_{LMCT}$$

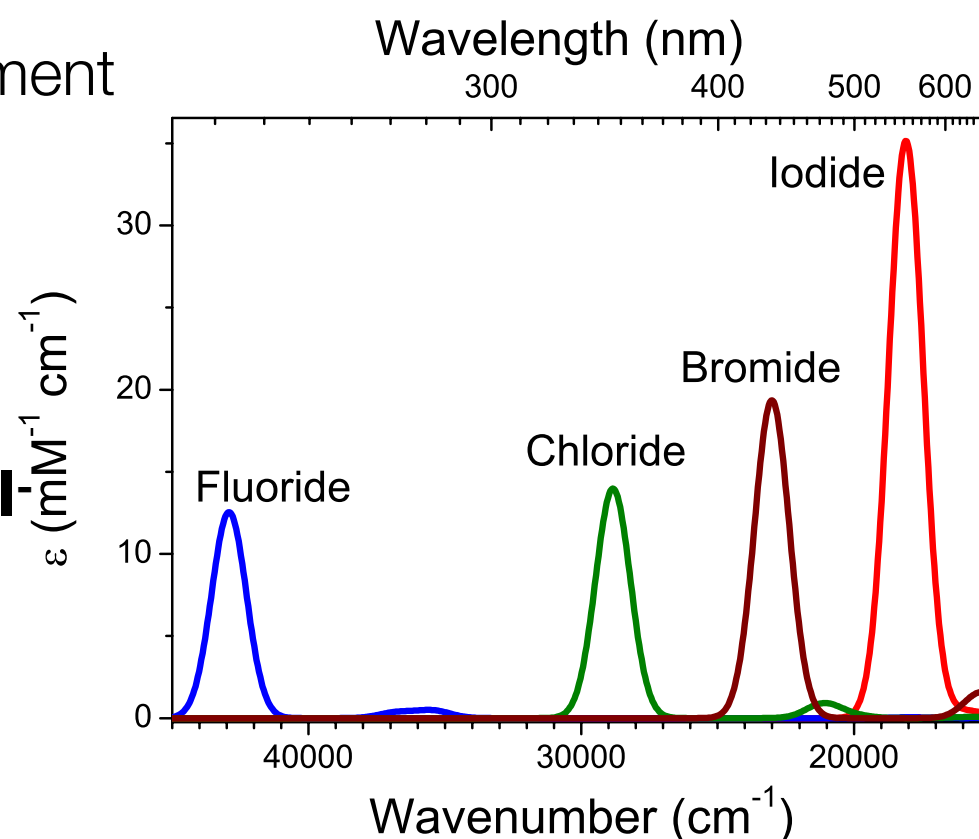


Let's Apply it (*in silico*)

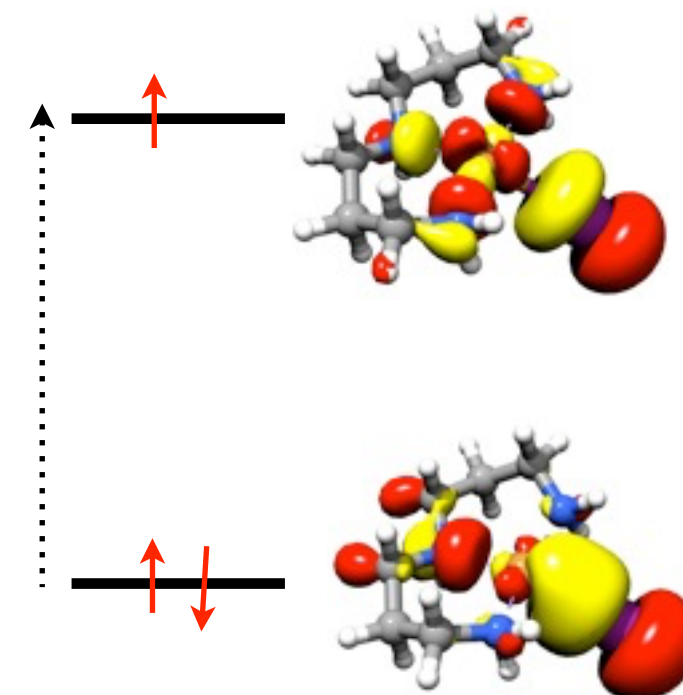
Hypothetical Cu(II) fragment



=F⁻, Cl⁻, Br⁻, I⁻



σ-LMCT Transition



	IP/eV	$R_{ML}/\text{\AA}$	D_{ML}^2/D^2	E_{LMCT}/cm^{-1}	Δ/eV	β/eV	β^2/Δ^2
F⁻	2,8	1,818	4,29	42918	4,79	1,12	0,05
Cl⁻	3,2	2,228	7,12	28340	3,13	0,77	0,06
Br⁻	3,4	2,373	12,35	23010	2,41	0,73	0,09
I⁻	3	2,644	28,53	18096	1,67	0,69	0,17