

# Fundamentals of Electron Paramagnetic Resonance

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**Bioinorganic Workshop 2024**

# EPR in 5 sec

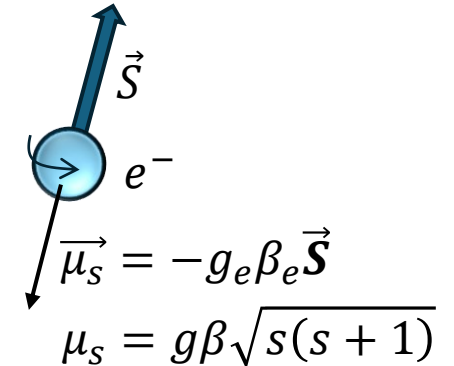
## Electron Paramagnetic Resonance

Electrons have spin angular momentum  $\vec{S}$   
The generated magnetic dipole moment  $\vec{\mu}_s$  interacts  
with external magnetic field

For  $s = \frac{1}{2}$  two spin states:  $m_s = +1/2$  and  $m_s = -1/2$

With magnetic field - degeneracy is lifted.

External magnetic oscillation @ microwave frequencies with energy matching  
that of the gap between  $m_s = +1/2$  and  $m_s = -1/2$ , will cause excitation of the system



$\beta_e = 9.27 \times 10^{-24} \text{ J/T}$  - Bohr magneton

$g$  - electron g-value

Free electron:  $g = g_e = 2.0023$

$s$  - spin quantum number

\* *ESR = Electron Spin Resonance (mostly used interchangeably with EPR)*

# Early history

- 1896. P. Zeeman and H. Lorentz:

**Splitting of sodium spectral lines in magnetic field. (Zeeman effect)**

- 1922. Stern, Gerlach:

**Electron magnetic moment in an atom can take only discrete orientations.**

- 1925. G.E. Uhlenbeck, S. Goudsmit:

**Link between magnetic moment and electron spin angular momentum.**

- 1938. I.I. Rabi.

**Transition between levels induced by an oscillating magnetic field.**

- 1944. E. Zavoisky (USSR):

**First experimental observation of an EPR (in  $\text{CuCl}_2 \cdot 2\text{H}_2\text{O}$ )**

- 1946. Cumberow, Halliday (USA).

**EPR in manganous salts.**

- 1947. Bagguley, Griffiths (UK)

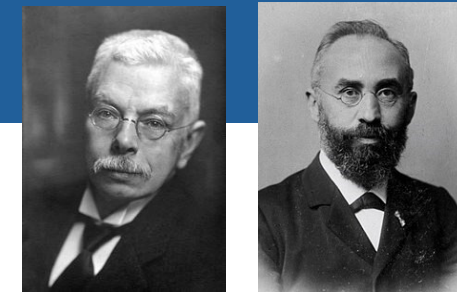
**EPR in chrome alum.**

# EPR spectroscopy

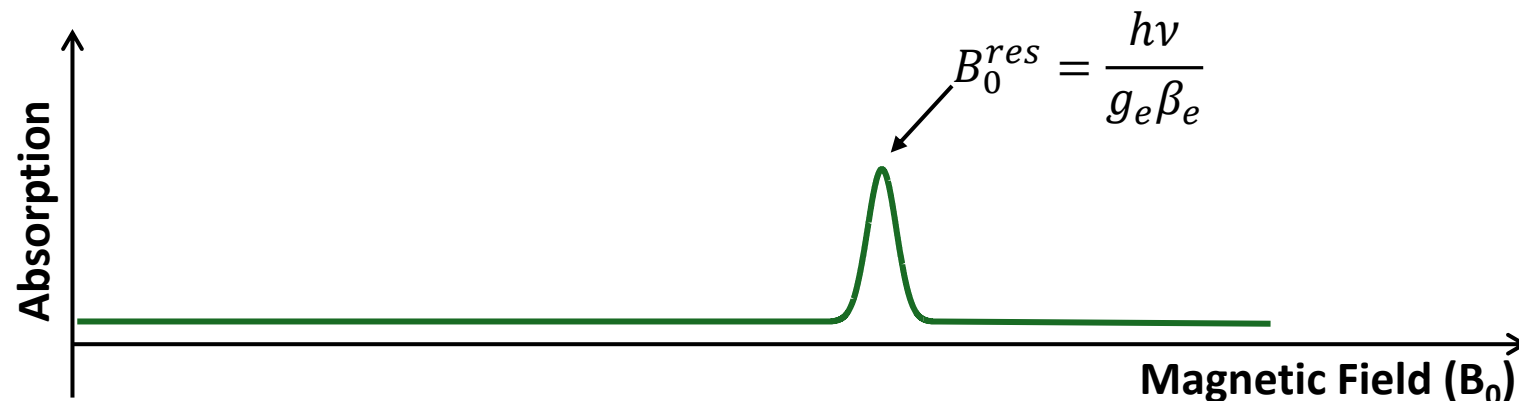
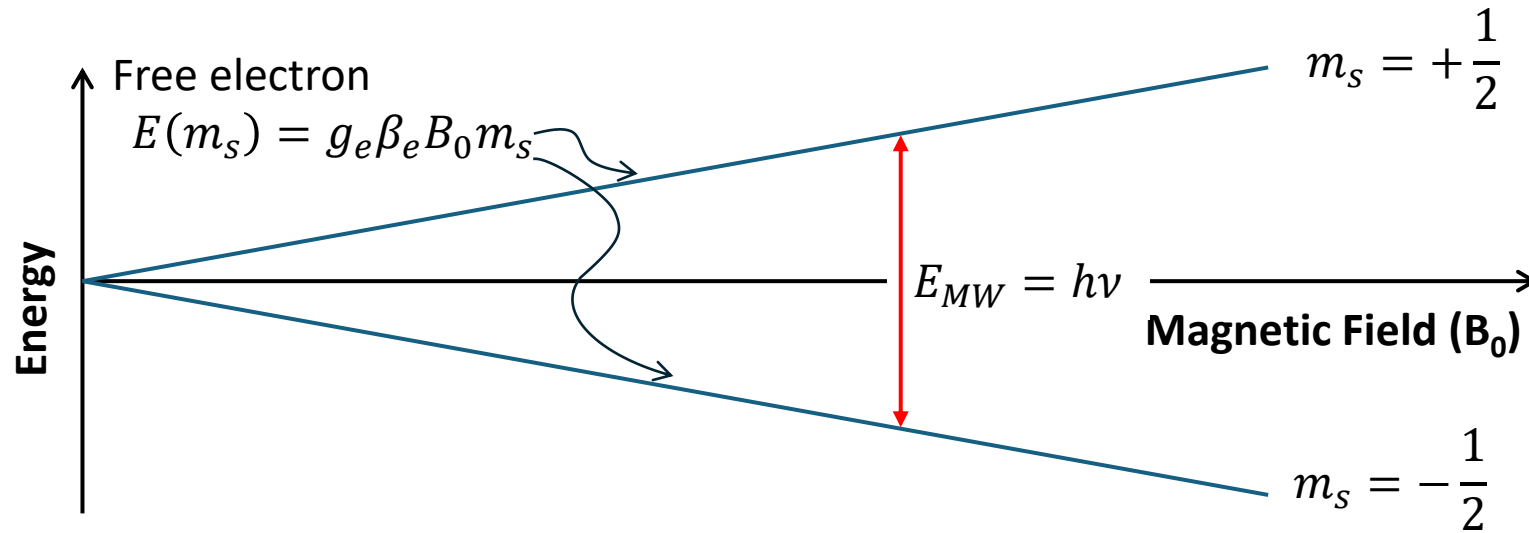
- Uses:
  - Organic free-radicals
  - Transition metal ions with an odd number of electrons or high spins
  - Excited states with  $S > 0$ , e.g., triplet states
  - Conduction electron spin resonance in nanoparticles
- Molecules with all electrons paired are *EPR-silent*
- Most integer spin systems are EPR-silent too (but not all)

- Zeeman effect and g-anisotropy
- Hyperfine coupling
- Zero-field splitting
- Practical aspects

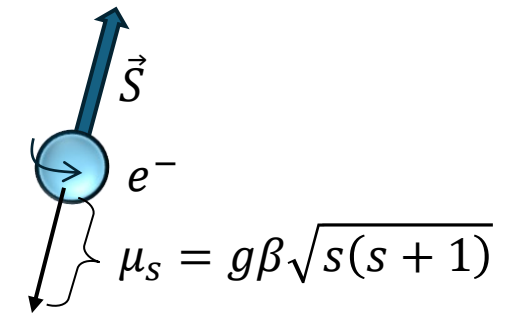
# Zeeman effect



Pieter Zeeman reported spectral line splitting of sodium in magnetic field in 1896  
Nobel prize in 1902, jointly with Hendrik Lorentz



$$\vec{\mu}_s = -g_e \beta_e \vec{S}$$



$\mu_s$  – magnetic dipole moment  
 $\beta_e = 9.27 \times 10^{-24} \text{ J/T}$  – Bohr magneton  
 $g$  – electron g-value  
 Free electron:  $g = g_e = 2.0023$   
 $s$  – spin quantum number  
 $h = 6.626 \times 10^{-34} \text{ J} \cdot \text{sec}$  – Planck constant  
 $\nu$  – microwave frequency

\* It is useful to keep in mind that the g-factor is a slope of  $E_{m_s}(B_0)$

# Where it gets interesting

In virtually any “real” case, the g-factor is not isotropic

- Movement of  $e^-$  on an orbit induces additional magnetic dipole moment

$$\vec{\mu}_L = -\beta_e \vec{L} \text{ and } \vec{\mu}_S = -\beta_e g_e \vec{S}$$

- Spin orbit coupling mixes ground and excited states. So, interaction of the system with an external magnetic field is described by a Hamiltonian:

$$\mathcal{H}_{Zeeman} = \beta_e \vec{B}_0 (\vec{L} + g_e \vec{S}) + \lambda \vec{L} \vec{S}$$

Because orbital momentum-derived magnetic moment is “locked” to the molecular frame (symmetry)

***Zeeman interaction will depend on how the molecule is oriented w.r.t. the external field.***

- the system is best described with a total angular momentum  $\vec{J} = \vec{L} + \vec{S}$ :

$$\mu = g\beta_e \sqrt{J(J+1)} \quad (\text{recall Russell-Saunders terms})$$

- Due to the SO-coupling (mixing) neither  $L$  or  $S$  are “good quantum numbers” (but  $J$  is)

# Spin Hamiltonian Formalism

- If SOC contribution is relatively small, we can describe system purely in term of “spin” (spin Hamiltonian formalism)
- BUT, we need to account for SOC and anisotropy of the Zeeman effect.
- So, g-factor MUST encode anisotropy associated with the existence of the spin-orbit coupling → **g-tensor**.

$$\mathcal{H}_{eff} = \beta_e \vec{B}_0 \mathbf{g} \vec{S} = \beta_e \begin{bmatrix} B_x & B_y & B_z \end{bmatrix} \begin{bmatrix} g_{xx} & g_{xy} & g_{xz} \\ g_{yx} & g_{yy} & g_{yz} \\ g_{zx} & g_{zy} & g_{zz} \end{bmatrix} \begin{bmatrix} S_x \\ S_y \\ S_z \end{bmatrix}$$

when SO coupling sufficiently small, we can use first-order perturbation to get:

$$g_{ij} = g_e + 2\lambda \sum_n \frac{\langle \phi_0 | L_i | \phi_n \rangle \langle \phi_n | L_j | \phi_0 \rangle}{E_n - E_0}$$

← mixing MOs  
← gap between electronic states

We can choose the coordinate system so that g-tensor is diagonal:

$$\mathcal{H}_{eff} = \beta_e \vec{B}_0 \mathbf{g} \vec{S} = \beta_e \begin{bmatrix} B_x & B_y & B_z \end{bmatrix} \begin{bmatrix} g_x & 0 & 0 \\ 0 & g_y & 0 \\ 0 & 0 & g_z \end{bmatrix} \begin{bmatrix} S_x \\ S_y \\ S_z \end{bmatrix}$$

$g_{x,y,z}$  are principal components of the g-tensor.

**This coordinate system is related to the molecular symmetry!**

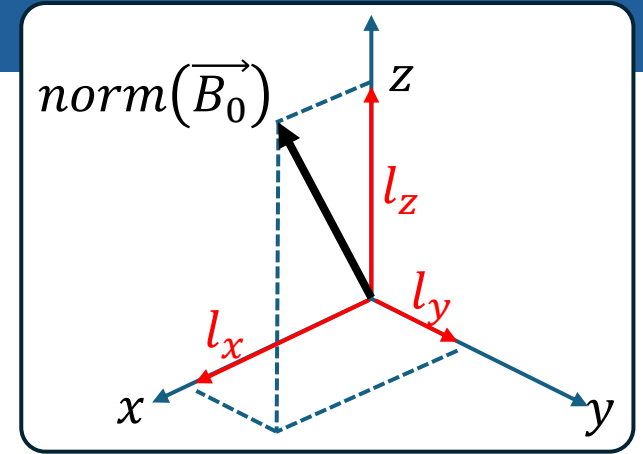
**Because g-tensor parameterizes SOC, it is a “fingerprint” of the system**



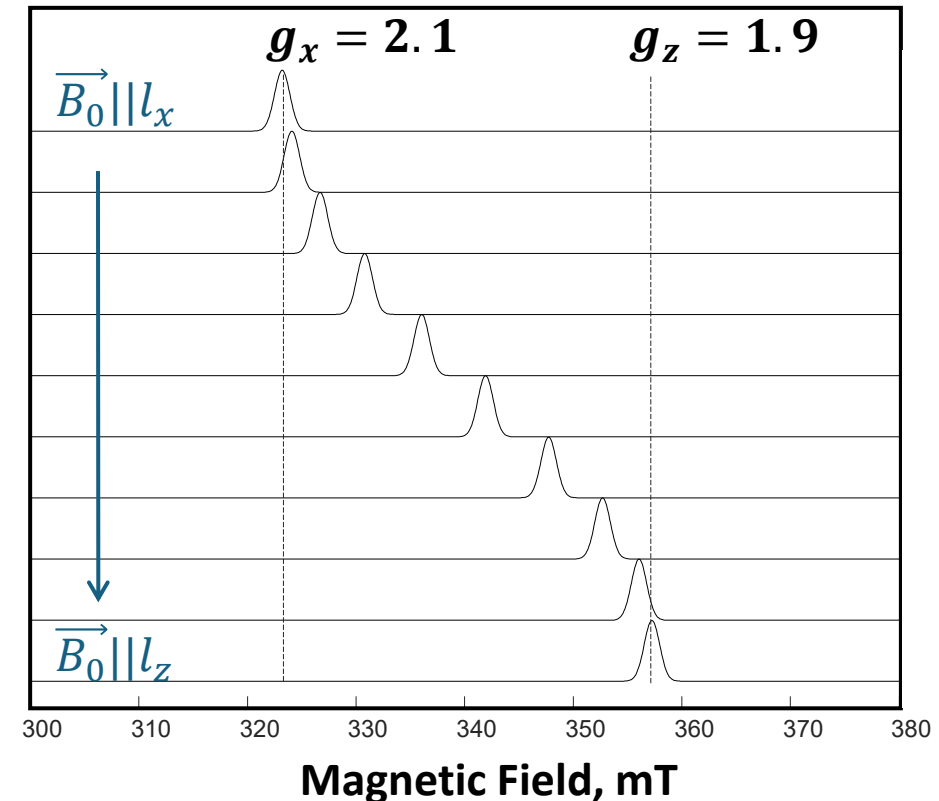
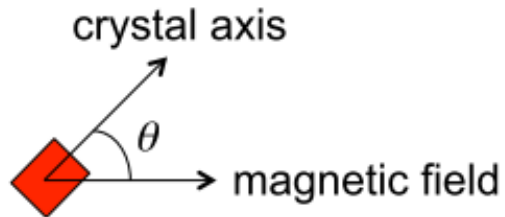
# g-anisotropy

In the absence of any other magnetic interactions (yes, there is more ... ) resonance position on an EPR spectrum is dictated by

$$g_{eff} = \sqrt{g_x^2 l_x^2 + g_y^2 l_y^2 + g_z^2 l_z^2} \quad \text{and} \quad B_0 = \frac{h\nu}{\beta_e g_{eff}}$$

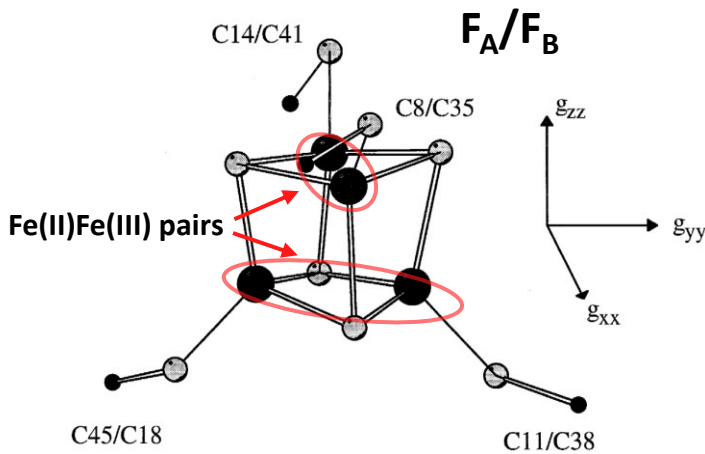
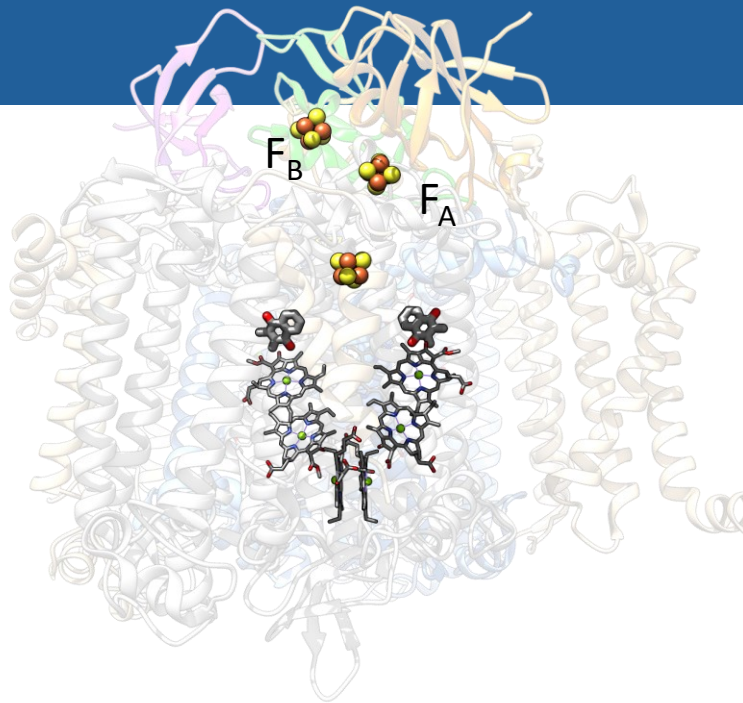
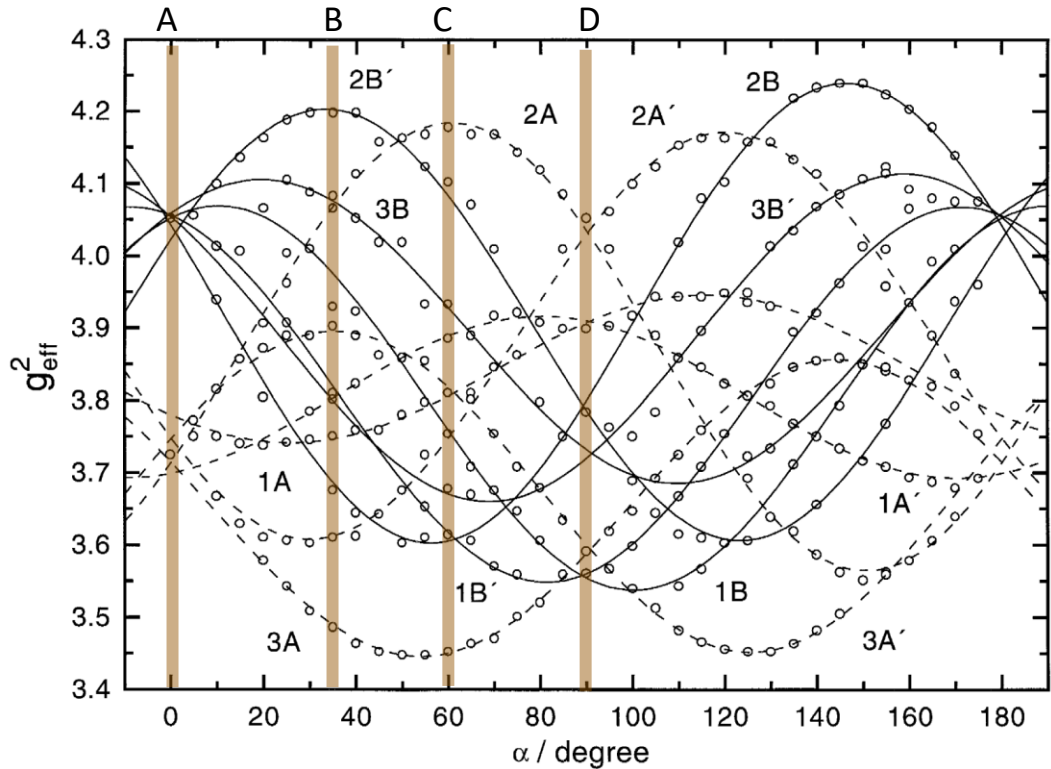
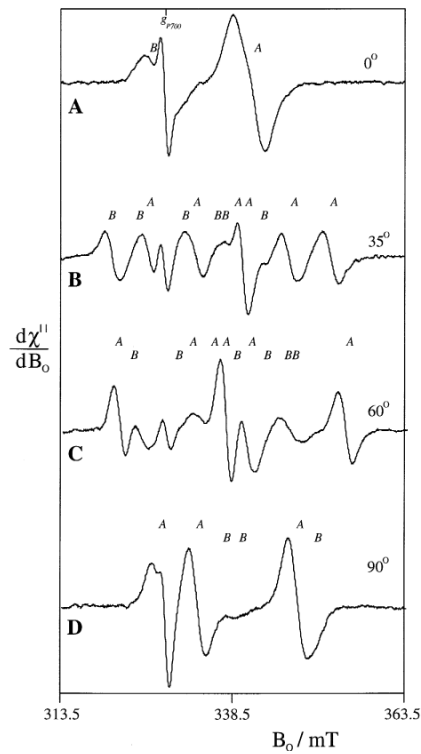


A single crystal “thought” experiment



$$* \frac{h}{\beta_e} = 71.448 [mT]/[GHz]$$

# Single crystal study

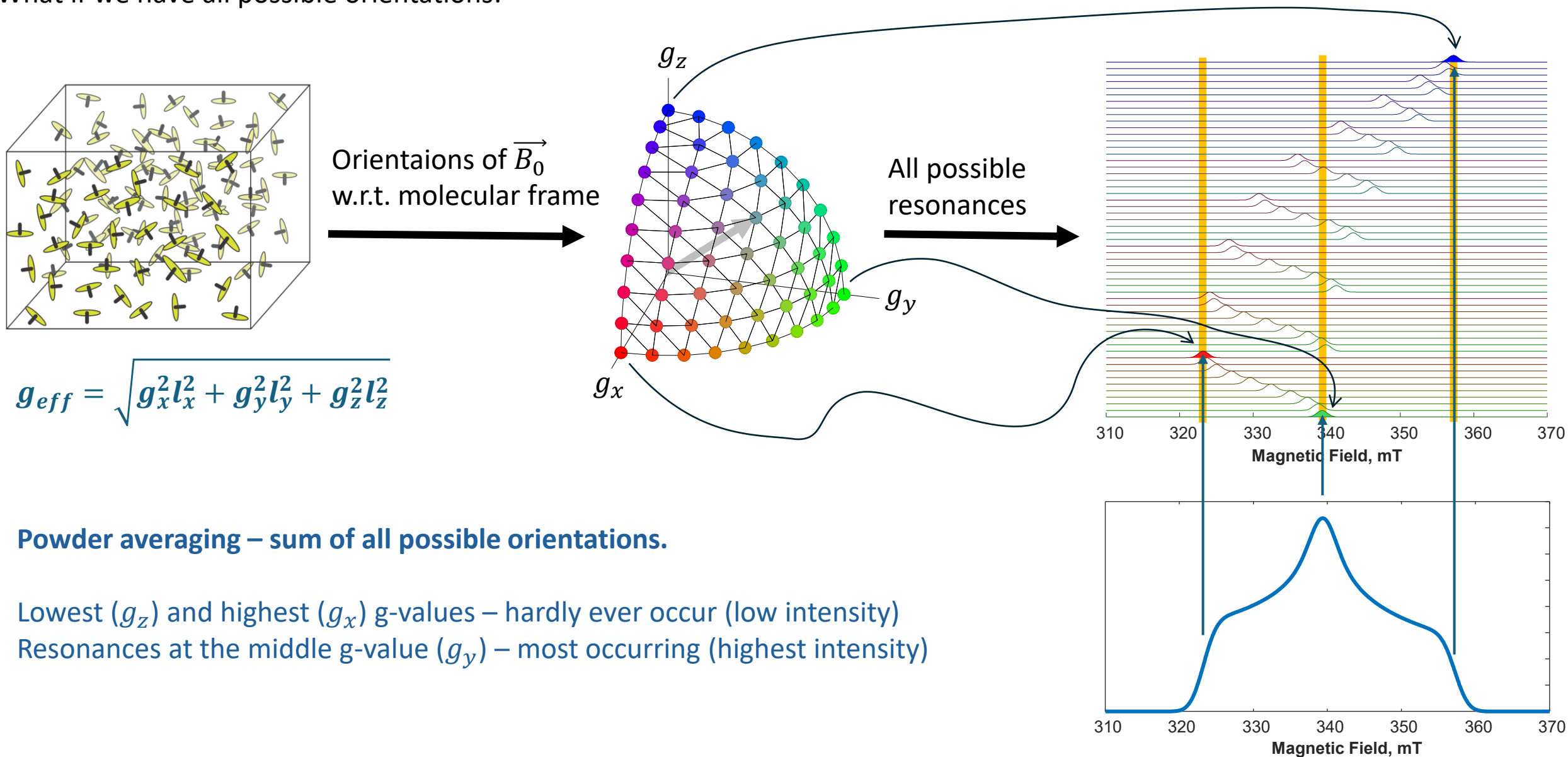


	g-Values of the FeS center			Angles with <i>c</i> -axis <sup>a</sup>		
	<i>g</i> <sub>xx</sub>	<i>g</i> <sub>yy</sub>	<i>g</i> <sub>zz</sub>	<i>β</i> <sub>x</sub>	<i>β</i> <sub>y</sub>	<i>β</i> <sub>z</sub>
F <sub>A</sub> <sup>-</sup>	1.856	1.941	2.051	53°	48°	64°
F <sub>B</sub> <sup>-</sup>	1.880	1.916	2.056	77°	60°	34°

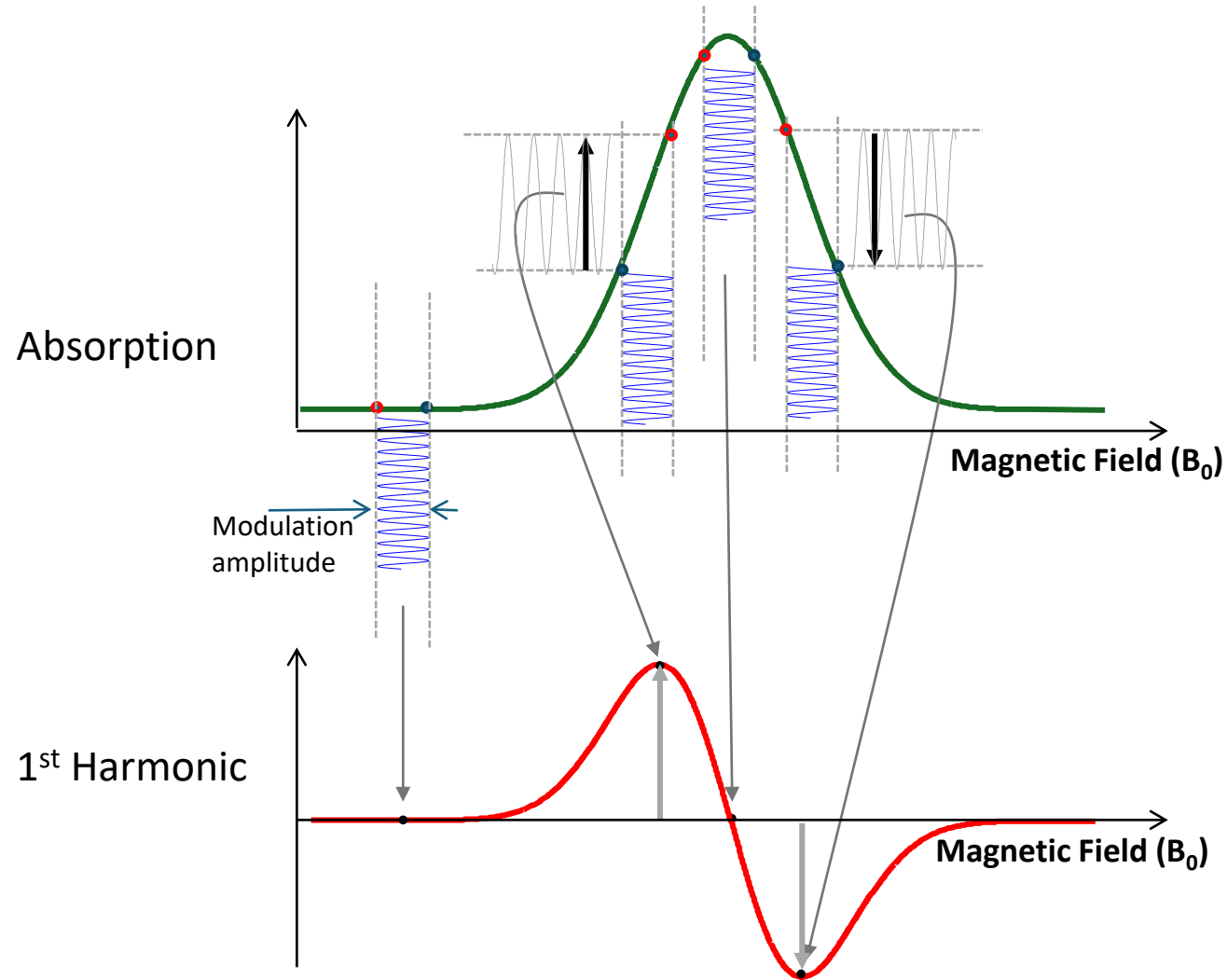
# Powder averaging

$$g_x = 2.1, g_y = 2.0, g_z = 1.9$$

What if we have all possible orientations:



# Field modulation, lock-in amplification



The amplitude and phase of the modulation is compared to a reference signal

**We effectively record the 1<sup>st</sup> derivative of the absorption line.**

# Types of shapes (g-anisotropy only)

Principal g-values can be obtained from the features of the spectra

## Isotropic case

$$g_x = g_y = g_z = g_{iso}$$

## Axial case

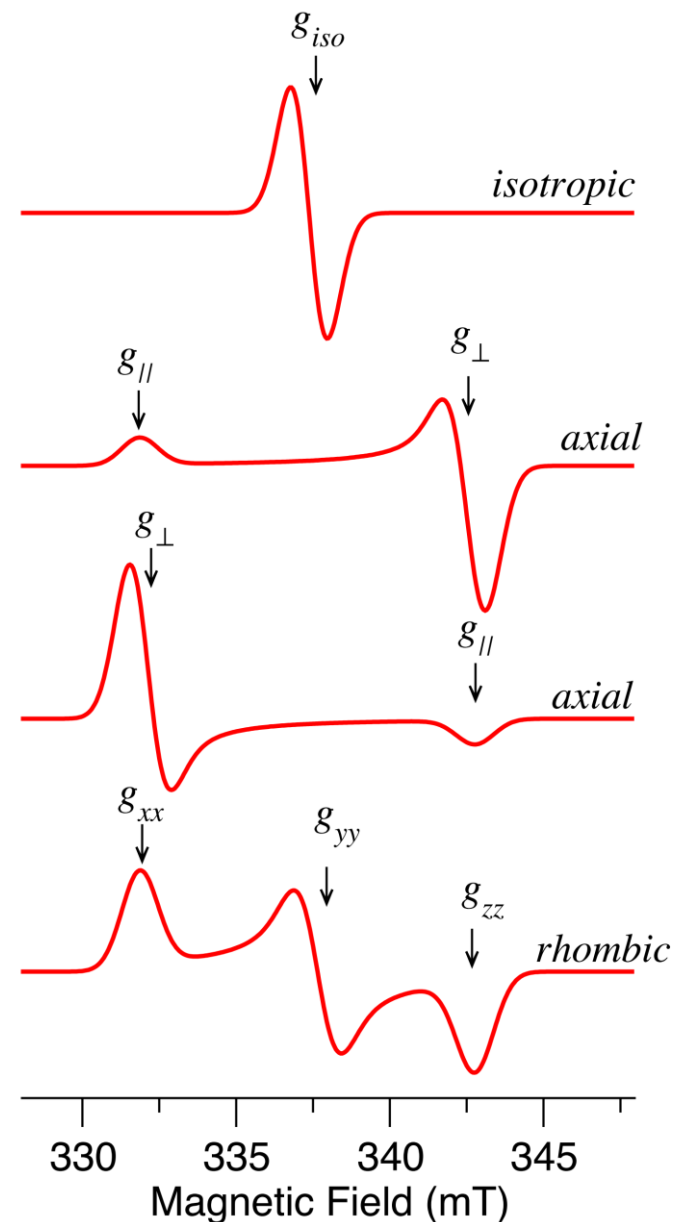
$$g_x = g_y = g_{\perp}, g_z = g_{\parallel}$$

it is common practice to assign  $g_{\parallel}$  to  $g_z$

## Rhombic case

$$g_x \neq g_y \neq g_z$$

exact assignment to the molecular frame depends on how much is known about the system (sometimes not much at all)



$$g_{ij} = g_e + 2\lambda \sum_n \frac{\langle \phi_0 | L_i | \phi_n \rangle \langle \phi_n | L_j | \phi_0 \rangle}{E_n - E_0}$$

- Organic free-radicals

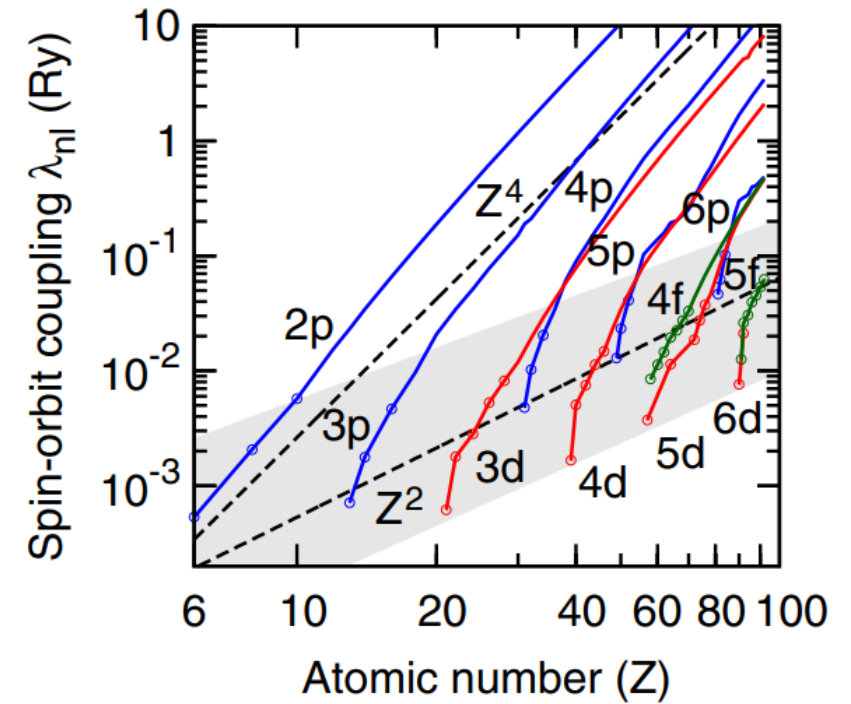
large  $\Delta E$ , small  $\lambda$

- hydrocarbon radicals:  $g \approx 2.002 - 2.003$
- N/O-based radicals:  $g \approx 2.003 - 2.006$
- S-based radicals :  $g \approx 2.007 - 2.010$

- Transition metal ions

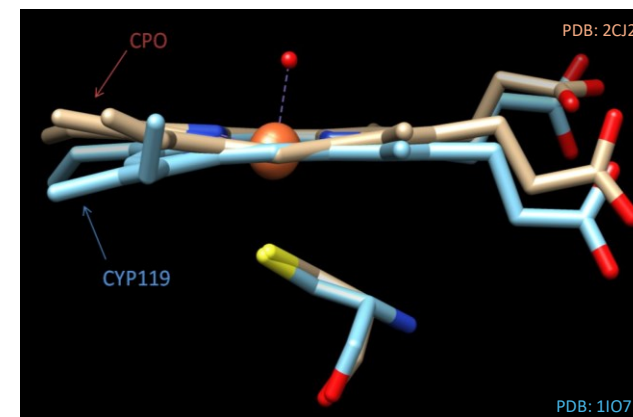
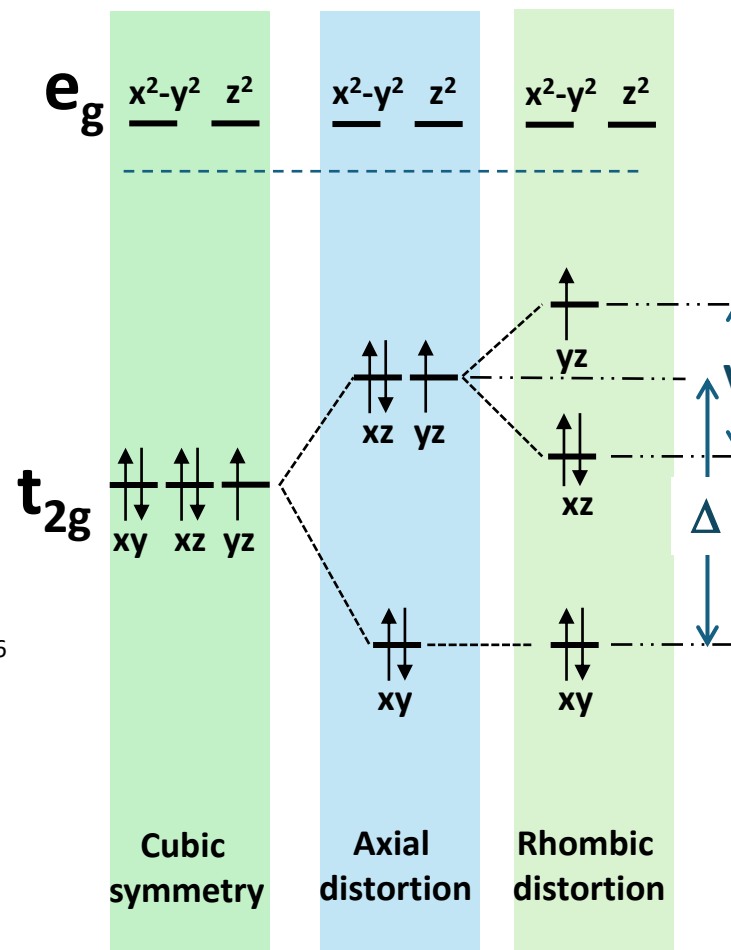
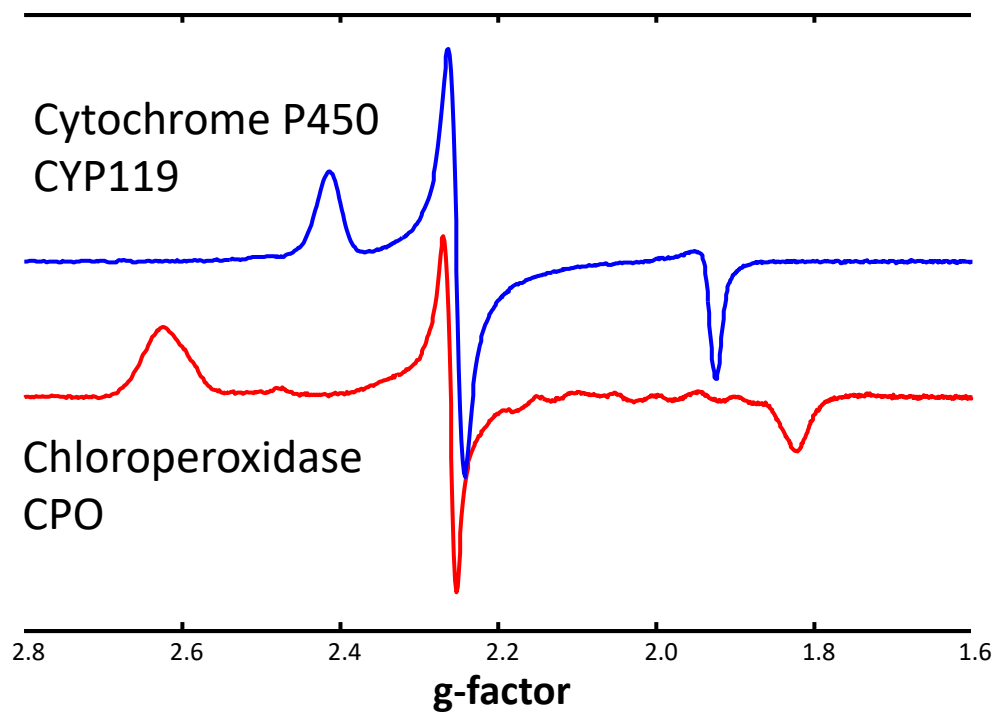
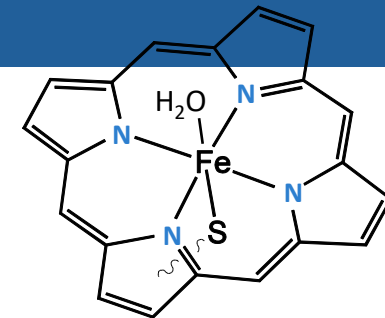
(relatively) small  $\Delta E$ , large  $\lambda$

- varies dramatically from system to system.
- can be below or above  $g_e = 2.0023$ 
  - V(IV)O,  $d^1, g \approx 1.9 - 1.99$
  - Cu(II),  $d^9, g \approx 2.1 - 2.2$



<https://journals.aps.org/prb/pdf/10.1103/PhysRevB.92.035135>

# Example: low spin Ferric heme

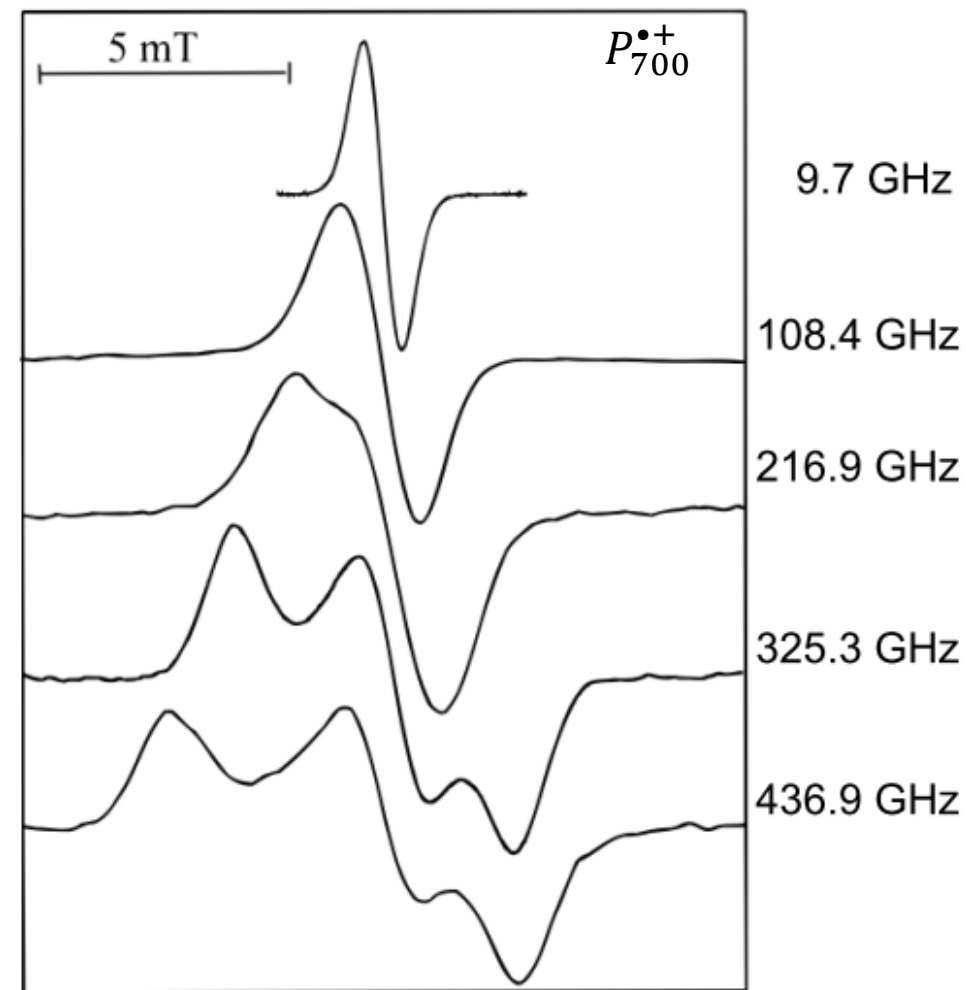
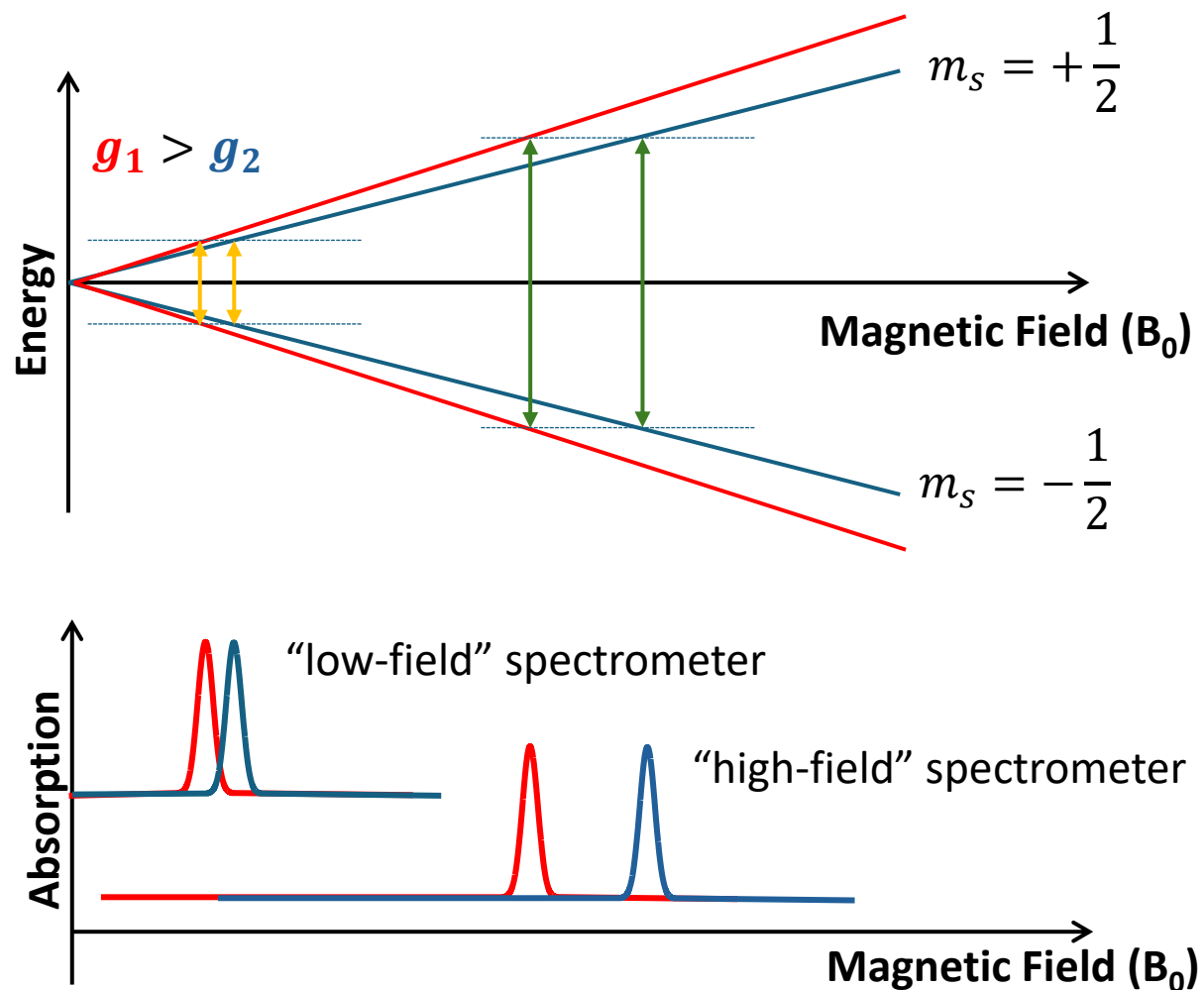


$$\frac{V}{\lambda} = \frac{g_x}{g_z + g_y} + \frac{g_y}{g_z - g_x}$$

$$\frac{\Delta}{\lambda} = \frac{g_x}{g_z + g_y} + \frac{g_z}{g_z - g_x} - \frac{V}{\lambda}$$

# Radicals

Some radicals are so isotropic, they require a high field to resolve g-anisotropy



$g$ : [ 2.00317, 2.00264, 2.00226 ]

Bratt et al. J.Phys.Chem.B (1997) 101 9689-9689



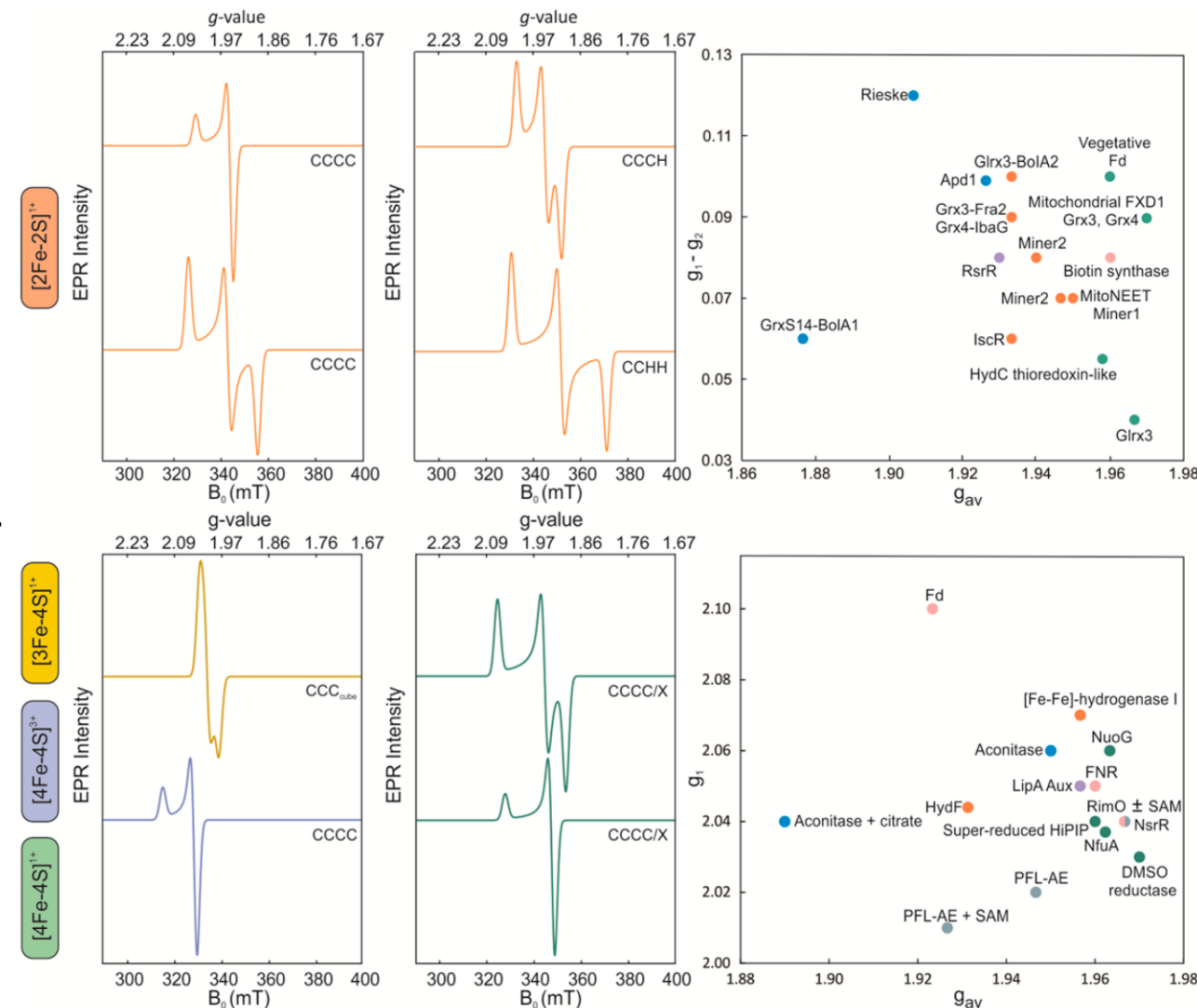
# FeS clusters etc.

## Coupled clusters are a more complicated case.

While the EPR spectrum may be relatively simple...  
Note that we (typically) only observe the ground state multiplet.

*Apparent g-anisotropy:*

- + SOC parameters of individual ions (intrinsic  $g, D$ )
- + valence delocalization (“double exchange”)
- + spin-spin interactions (exchange)



Quist et al. Inorganics 2023, 11(12), 475

# Hyperfine interaction

Interaction between unpaired electron and a magnetic nucleus leads to splitting of energy levels

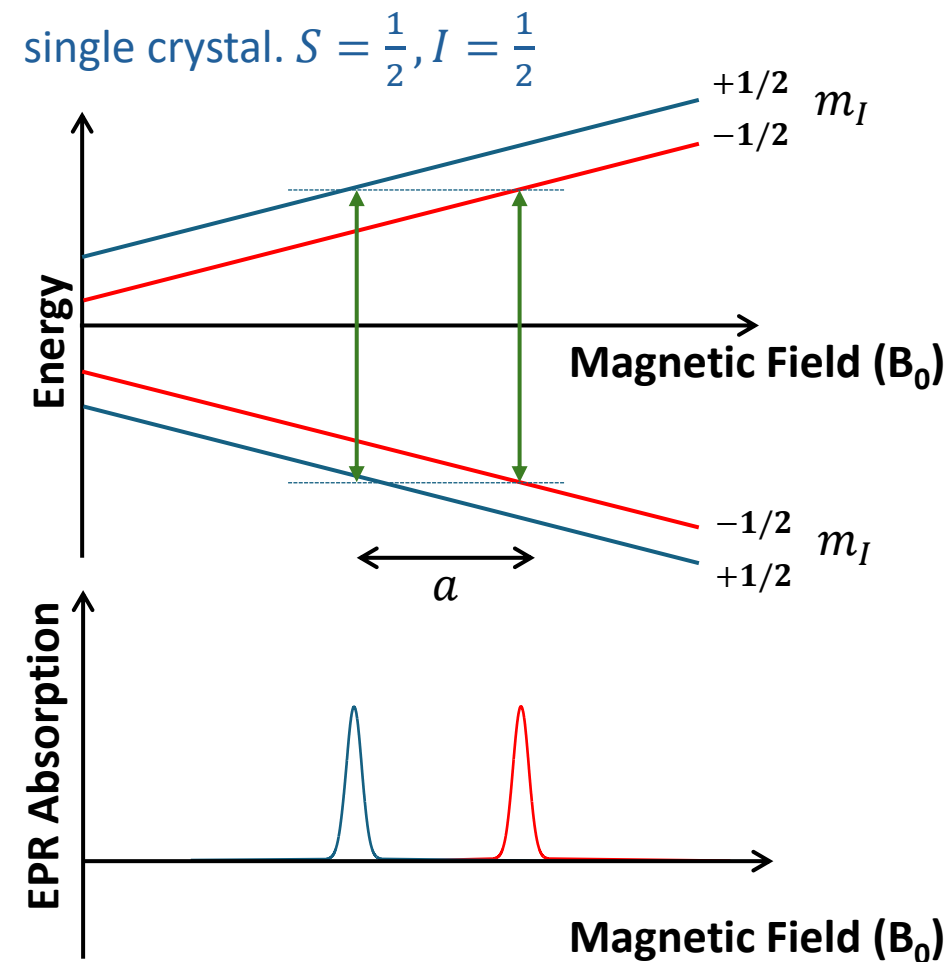
## Isotropic contribution:

Direct spin density at the nucleus

$$a_{iso} = \frac{2}{3} \frac{\mu_0 \beta_e \beta_n}{h} g_e g_n |\psi(0)|^2$$

## Anisotropic contribution (dipolar):

$$a_{dip} = \frac{\mu_0}{4\pi} \frac{\beta_e \beta_n}{h} g_e g_n \left\langle \frac{3 \cos^2 \Theta - 1}{r^3} \right\rangle$$

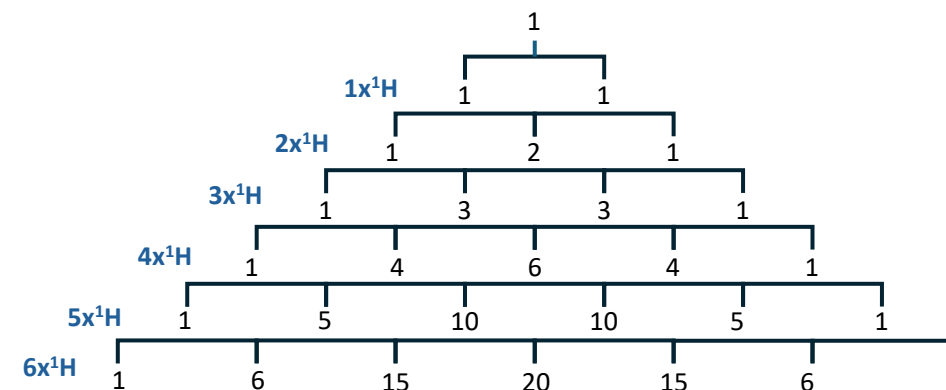
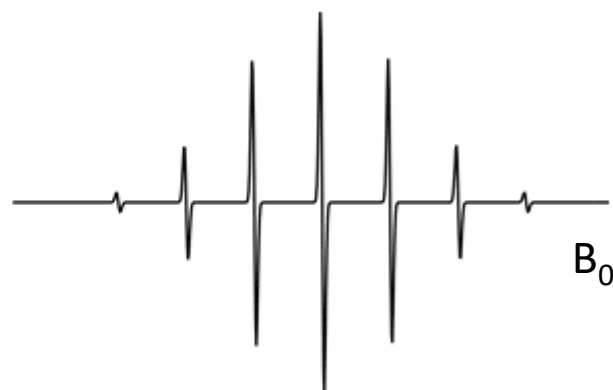
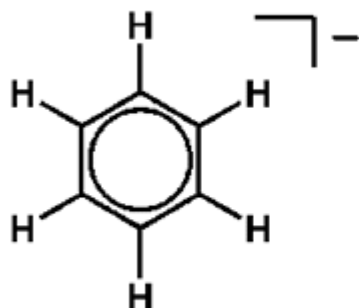


# Isotropic HF coupling

Each nucleus adds to the splitting

Extreme case of high symmetry:

Benzene Radical Anion

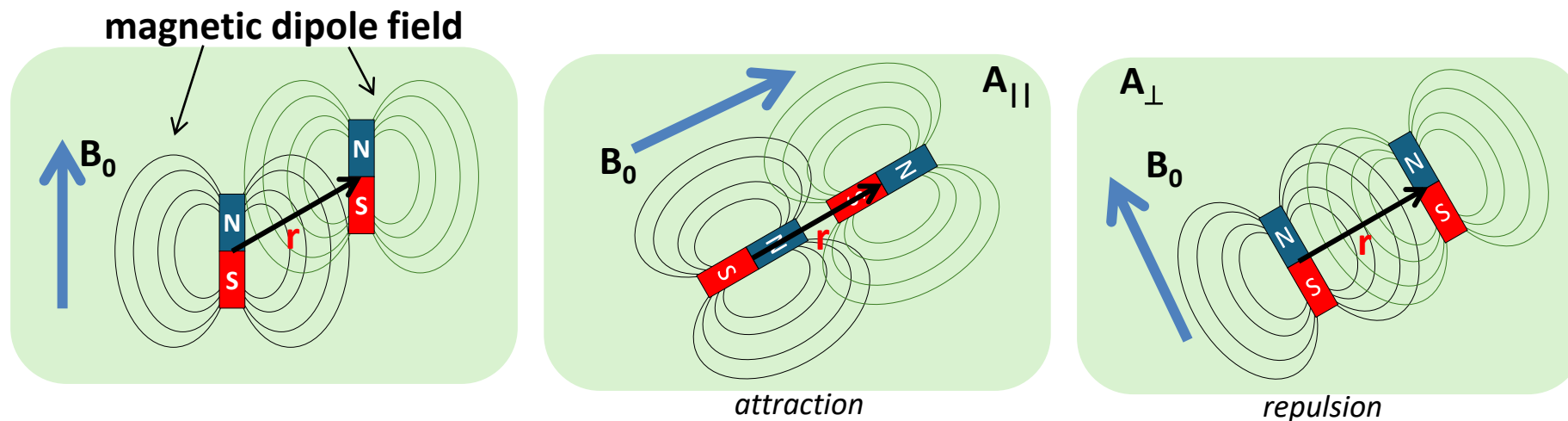


McConnel rule for <sup>1</sup>H HF couplings on C<sub>α</sub>:

$$A_{iso}^{\alpha} = \rho_{\pi}^{\alpha} Q_{\alpha}, \rho_{\pi}^{\alpha} \text{-spin density on } C_{\alpha}, Q_{\alpha} \text{-proportionality constant}$$

# Dipolar HF coupling

“classical” interpretation: interaction of two spins (nuclear and/or electron) through space:



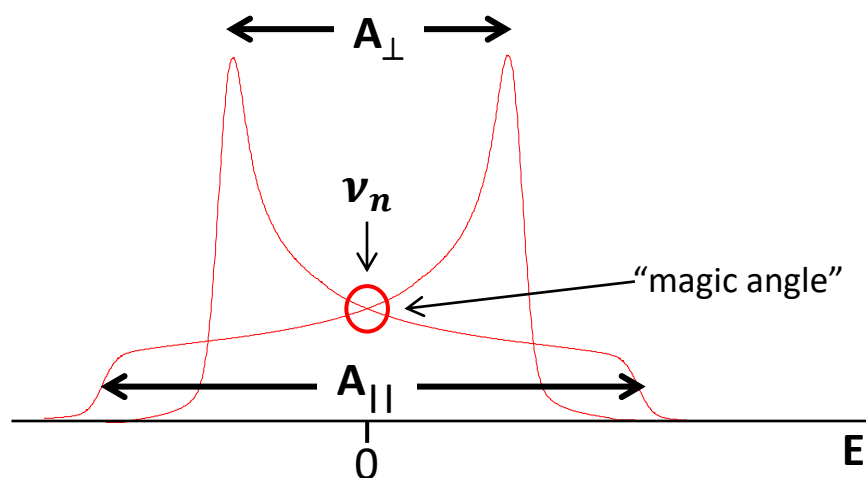
Energy of the interaction change sign as we rotate the sample in external magnetic field

Pure Hyperfine dipole-dipole interaction:

$$A_{x,y} = A_{\perp} = T$$

$$A_z = A_{||} = -2T$$

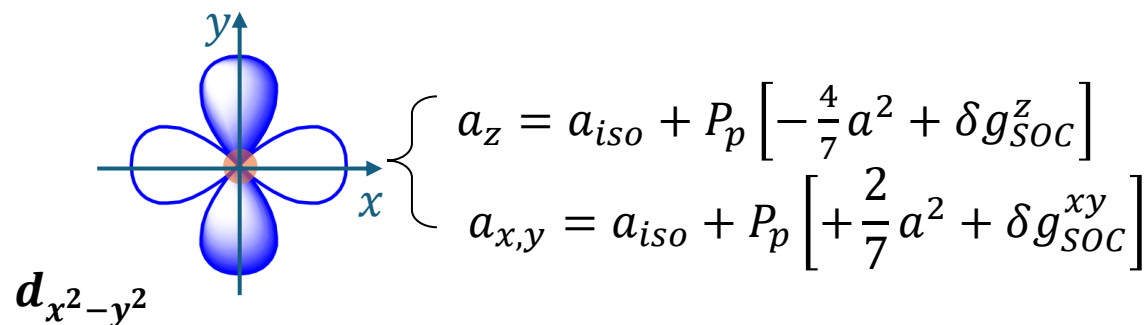
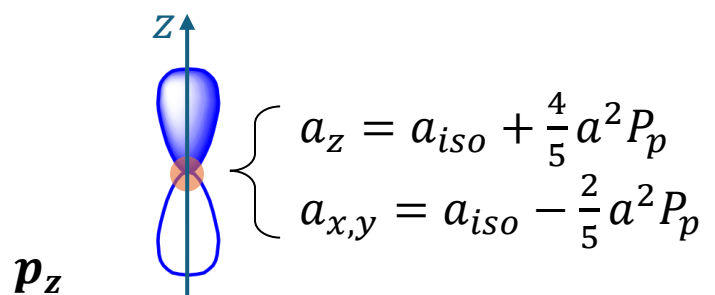
$$T = \underbrace{\frac{\mu_0}{4\pi h} \beta_e \beta_n}_{\text{just constants}} \cdot \frac{g_e g_n}{r^3}$$



# Anisotropic HF coupling

Nucleus at the center of anisotropic electron spin density (e.g. metal centers  $^{57}\text{Fe}$ ,  $^{61/63}\text{Cu}$  etc.) will exhibit anisotropic HF interaction

For example:



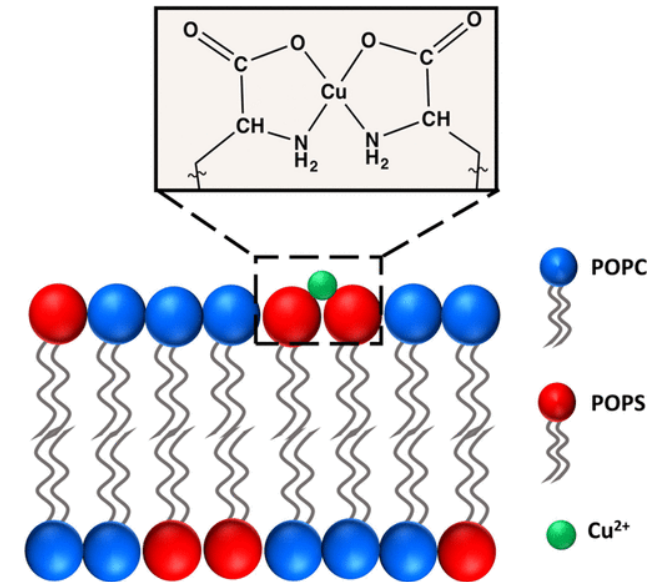
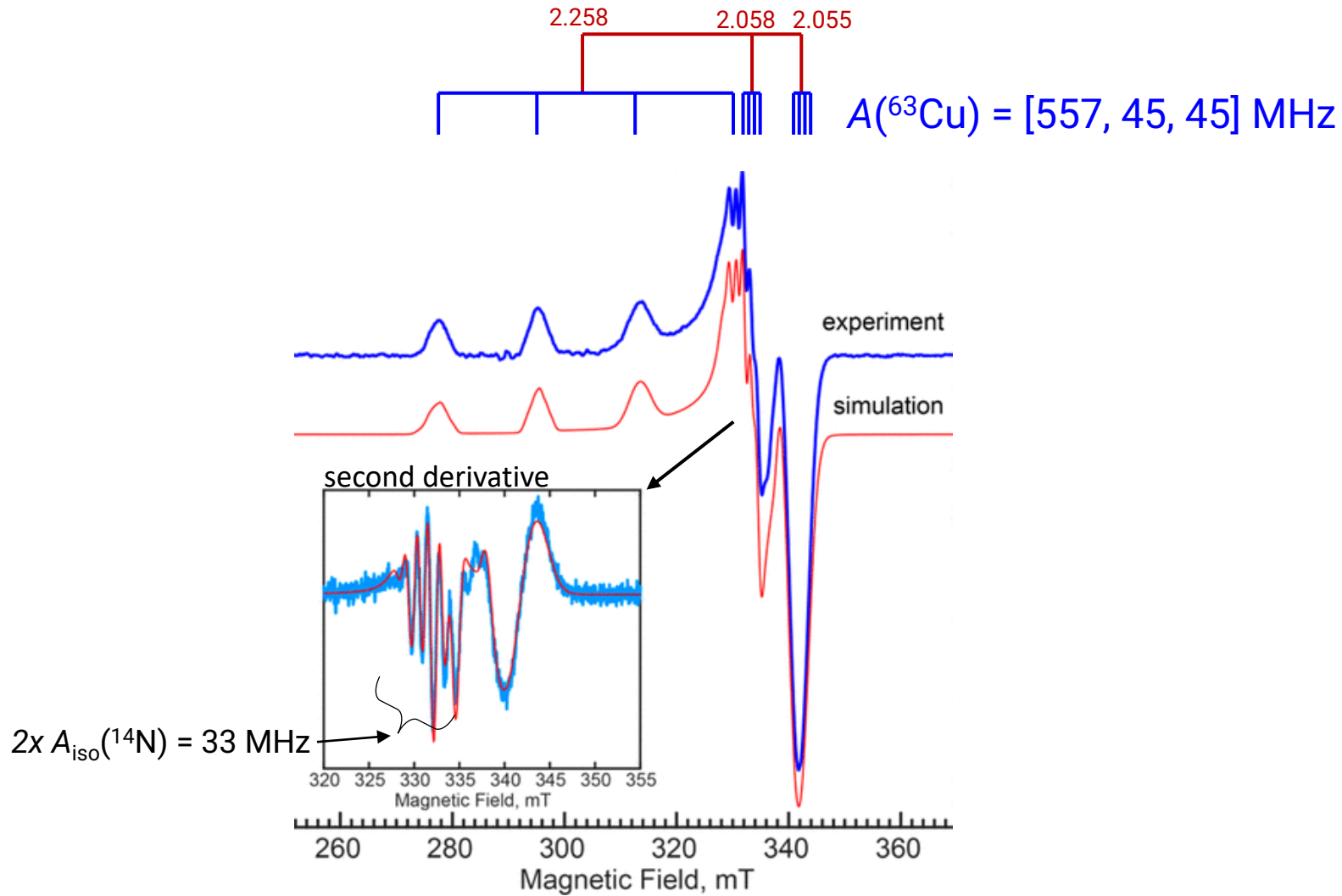
$a$  - covalency, e.g., amount of p (or d) character in MO

$P_p, P_d$  - proportionality constants

$\delta g_{SOC}$  - contributions from spin-orbit coupling (SOC)

$a_{iso}$  - originating from a pseudo-fermi contact (think of it as  $s$ - $p$ ,  $s$ - $d$  MO mixing)

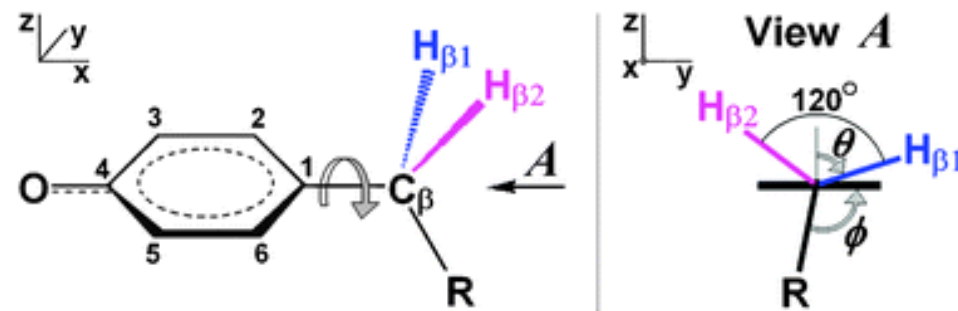
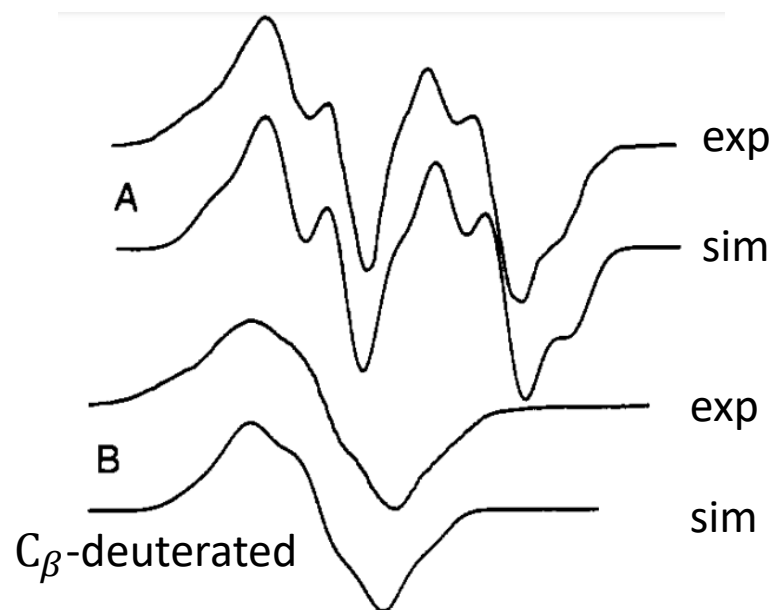
# Example: $\text{Cu}^{2+}$ EPR



\*note that  $^{63}\text{Cu}$  and  $^{61}\text{Cu}$  isotopes have similar abundance.  
Second simulation using  $g_n$ -scaled HFCs for the latter added into total simulation.

# Example: Tyrosyl radical

*E.coli* RNR



$$A_{iso}^{\beta_1} = \rho_{C1}(B' + B'' \cos^2 \Theta)$$

$$A_{iso}^{\beta_2} = \rho_{C1}(B' + B'' \cos^2 \{\Theta - 120^\circ\})$$

$$g_{xyz} = [2.0077, 2.0054, 2.0023]$$

$$A_{xyz}^{3,5} = [-9.6, -2.8, -7] \text{ MHz}$$

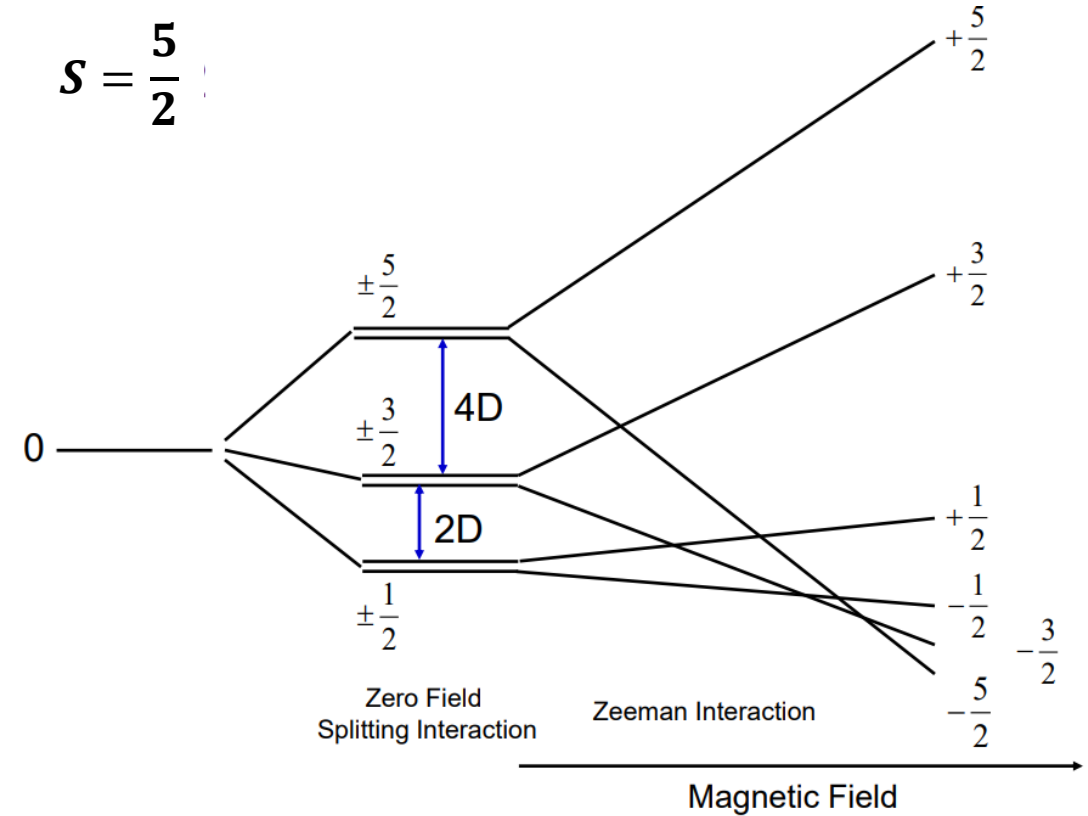
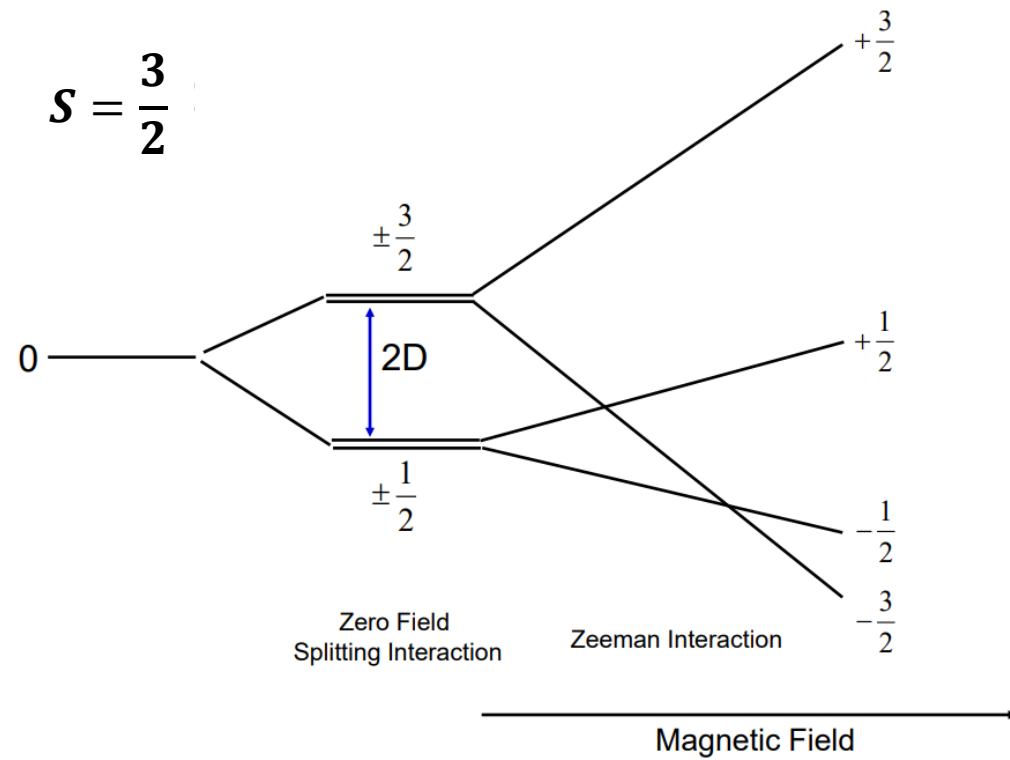
$$A_{xyz}^{\beta_1} = [19.6, 19.6, 21.2] \text{ MHz}$$

$$A_{xyz}^{\beta_2} = [-0.17, -0.17, 1.75] \text{ MHz}$$

# Zero Field Splitting

$S > 1/2$ : SOC and spin-spin coupling splits the spin states according to  $|m_s|$

This effect has a profound effect on the EPR spectra





# Zero Field Splitting

Addition to spin Hamiltonian:  $\mathcal{H}_{ZFS} = \vec{S} \mathbf{D} \vec{S}$

In its principal axes:

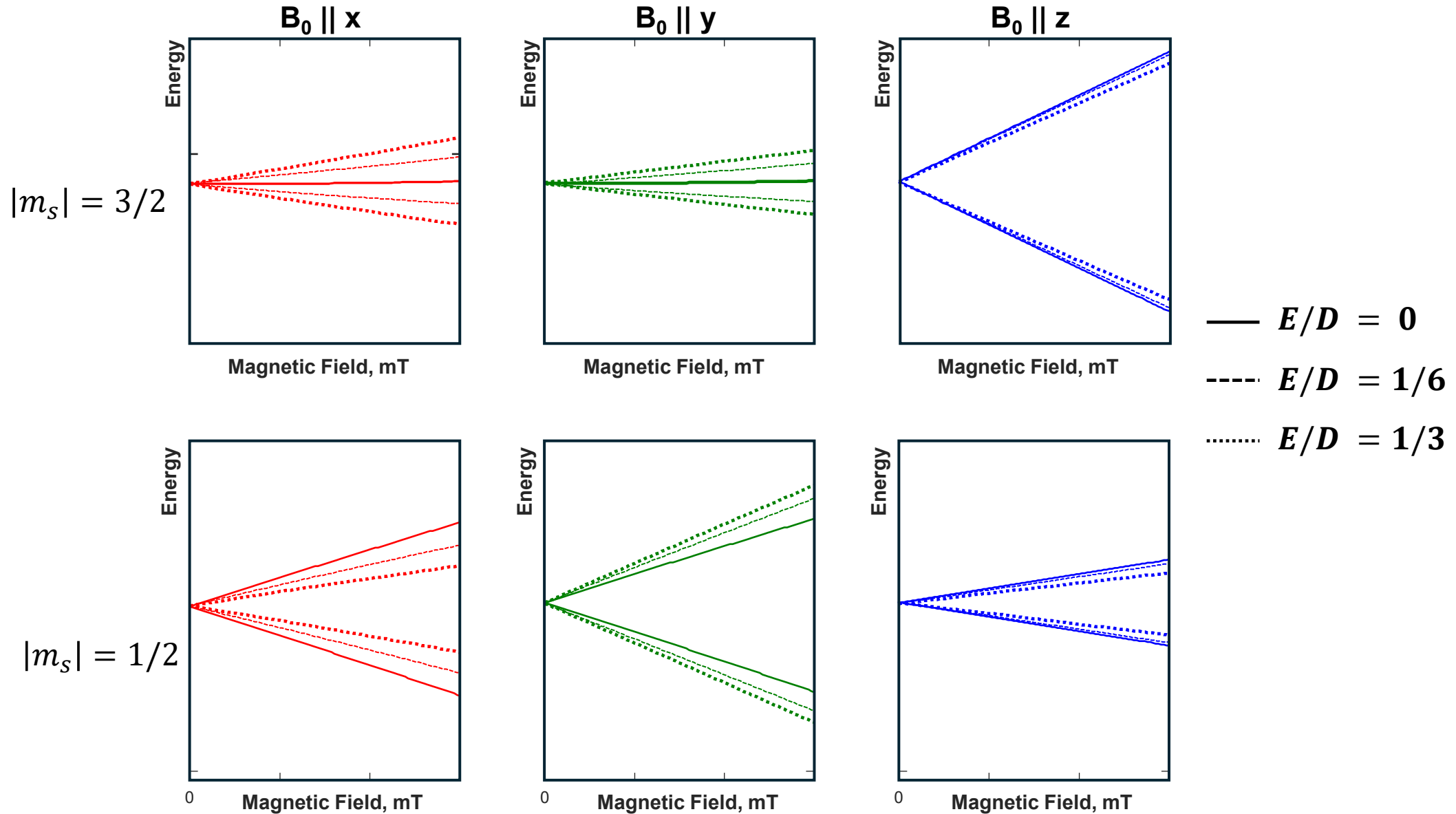
$$\mathbf{D} = \begin{bmatrix} -\frac{1}{3}D + E & 0 & 0 \\ 0 & -\frac{1}{3}D - E & 0 \\ 0 & 0 & \frac{2}{3}D \end{bmatrix}$$

High symmetry systems with little-to-no SOC,  $D, E \cong 0$

Systems with an axial symmetry,  $E \cong 0$

Note that “rhombicity”  $E$  takes values from 0 to  $\frac{1}{3}D$

# Zero Field Splitting ( $S=3/2$ )

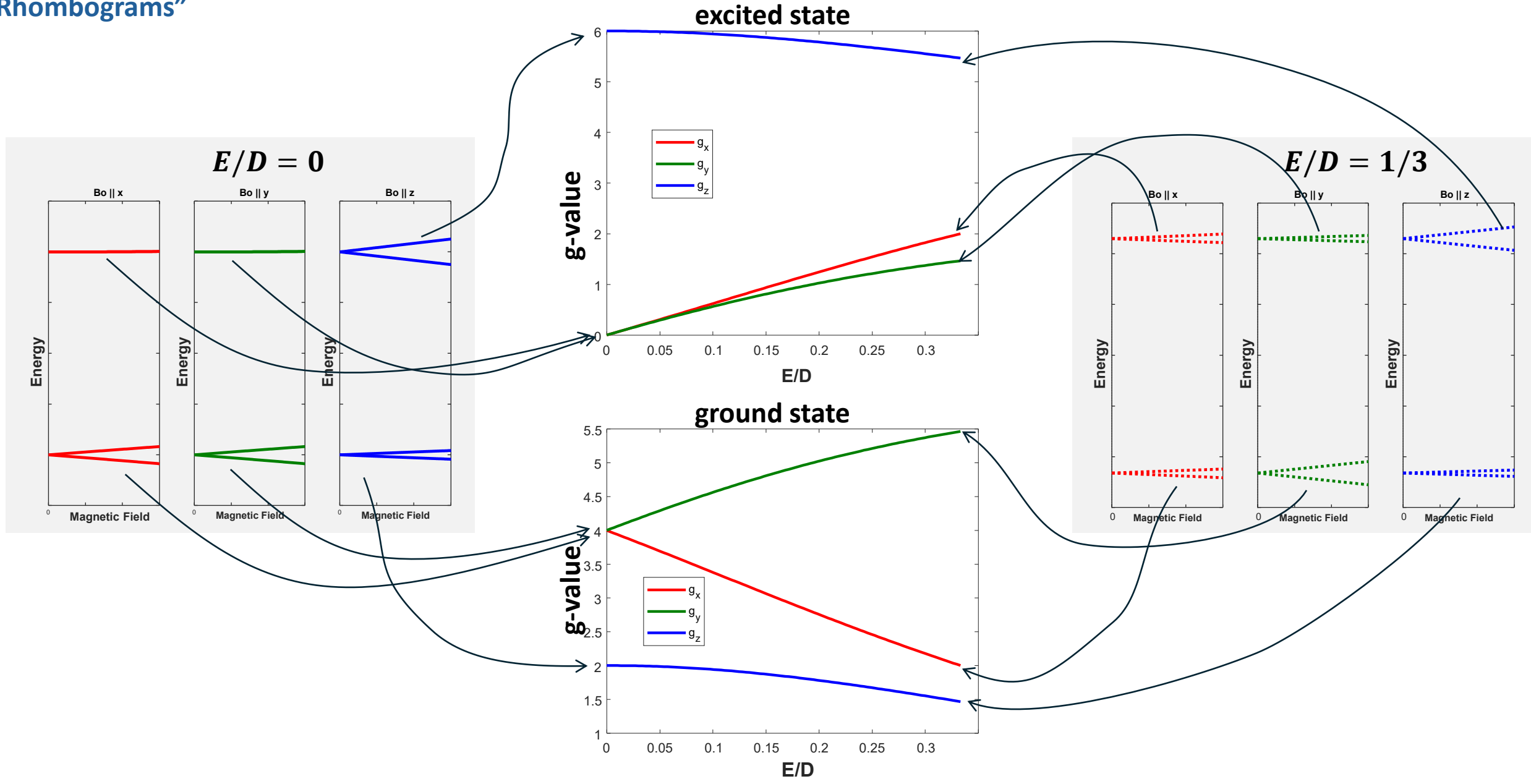


Wouldn't it be nice to streamline prediction of EPR spectra in terms of effective g-values?

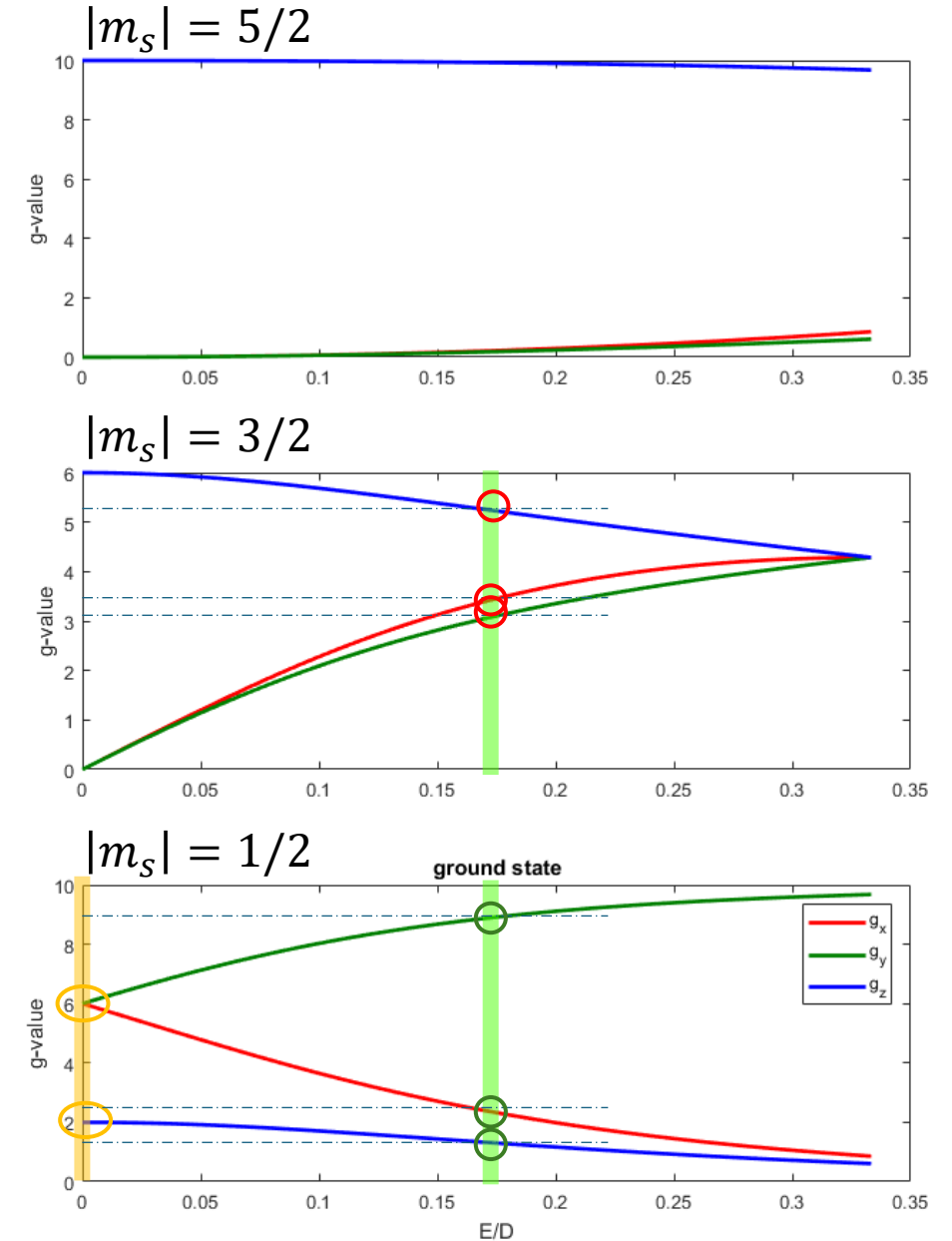
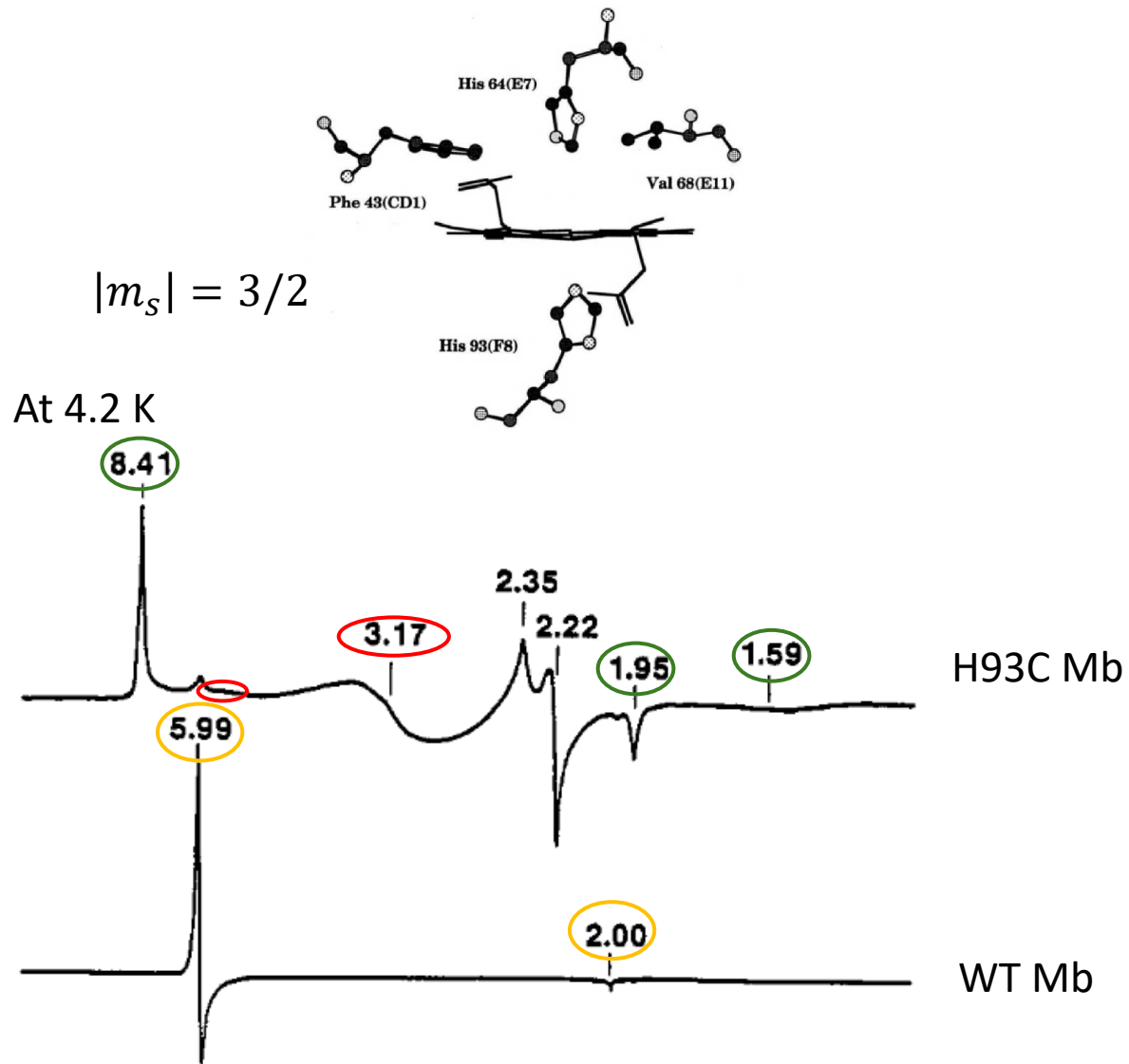
# Zero Field Splitting ( $S=3/2$ )

Let's treat every doublet as an effective  $S = 1/2$   
The slope of  $E(B_0)$  represents effective  $g$  value

“Rhombograms”



# Example: High spin Human Myoglobin ( $S=5/2$ )



# Let's talk about instrumentation



\* “EPR spectrometer” according to AI

# Spectrometer design

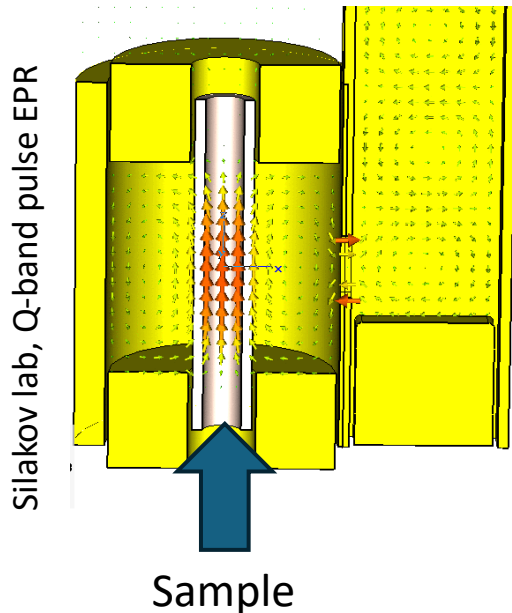
**It is much easier to change magnetic field than do a broad-band sweep @ MW frequency**

- Frequency is fixed within a range
- Magnetic field is adjusted

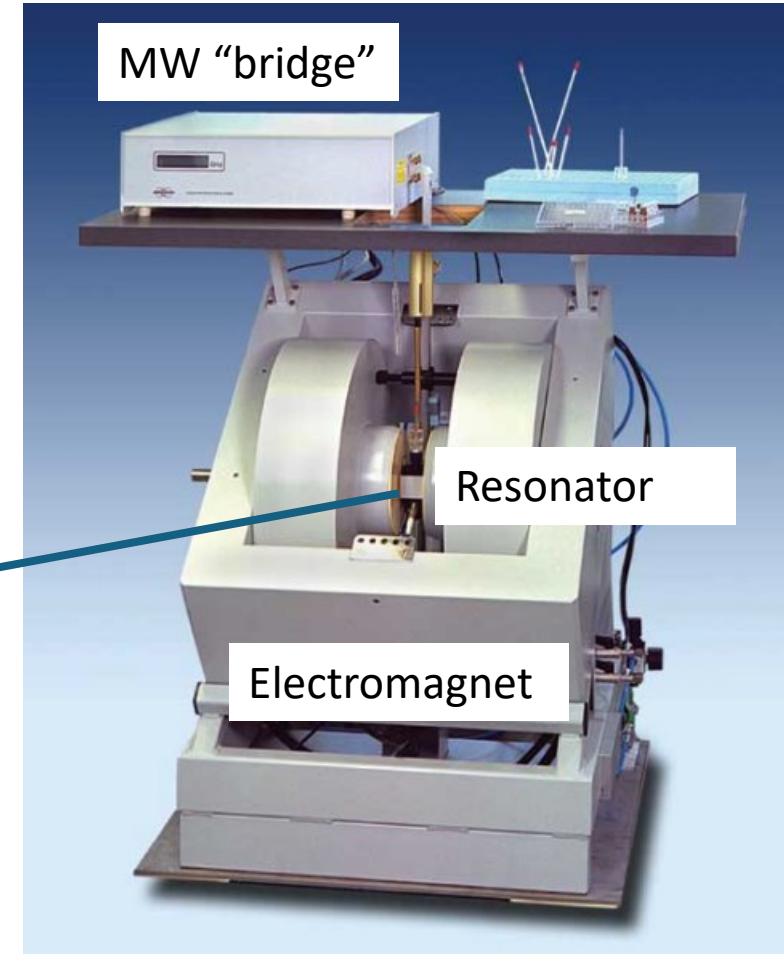
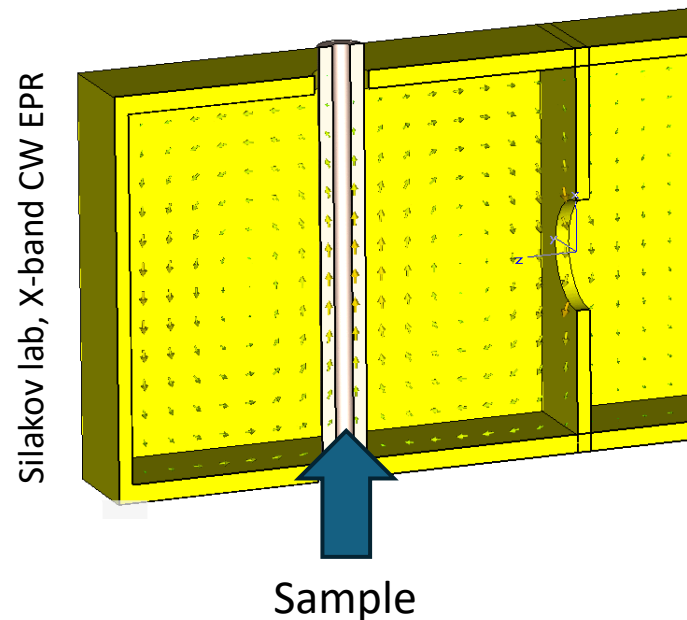
## Resonators

- Really narrows the band, but substantially improves sensitivity

**Cylindrical cavity**



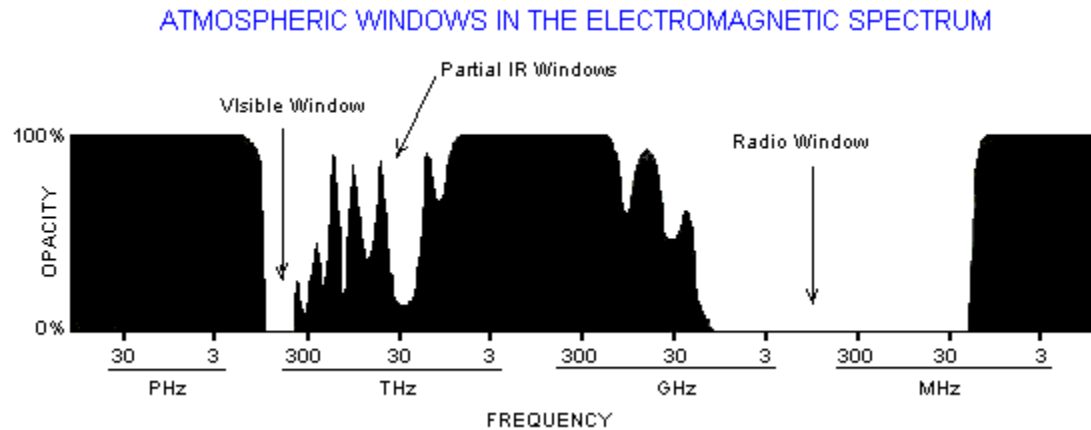
**Rectangular cavity**



# Spectrometer design

Historically, only microwave components at specific MW “bands” were available due to gaps in atmospheric EM absorption.

- EPR Spectrometers are operated at specific “bands”

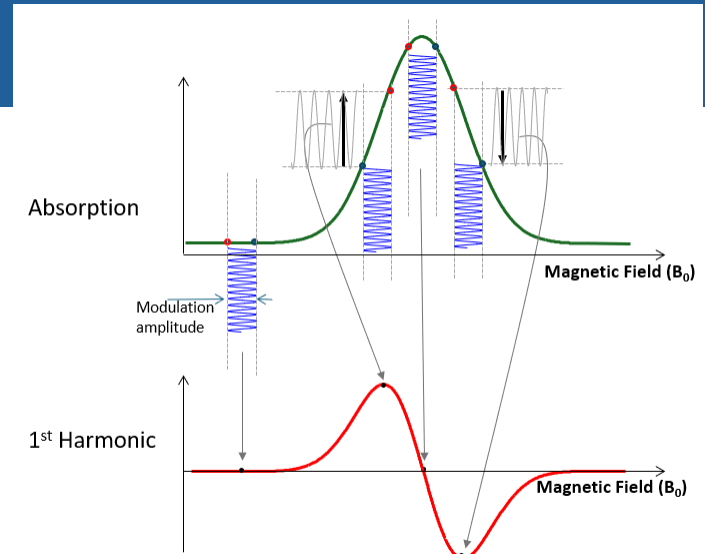
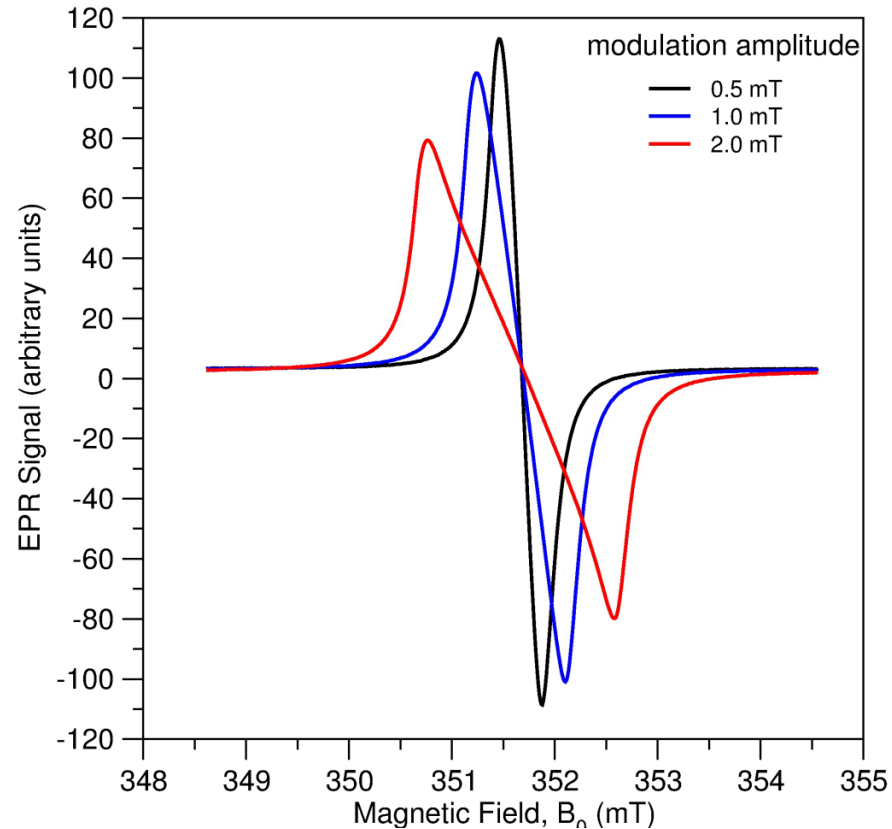
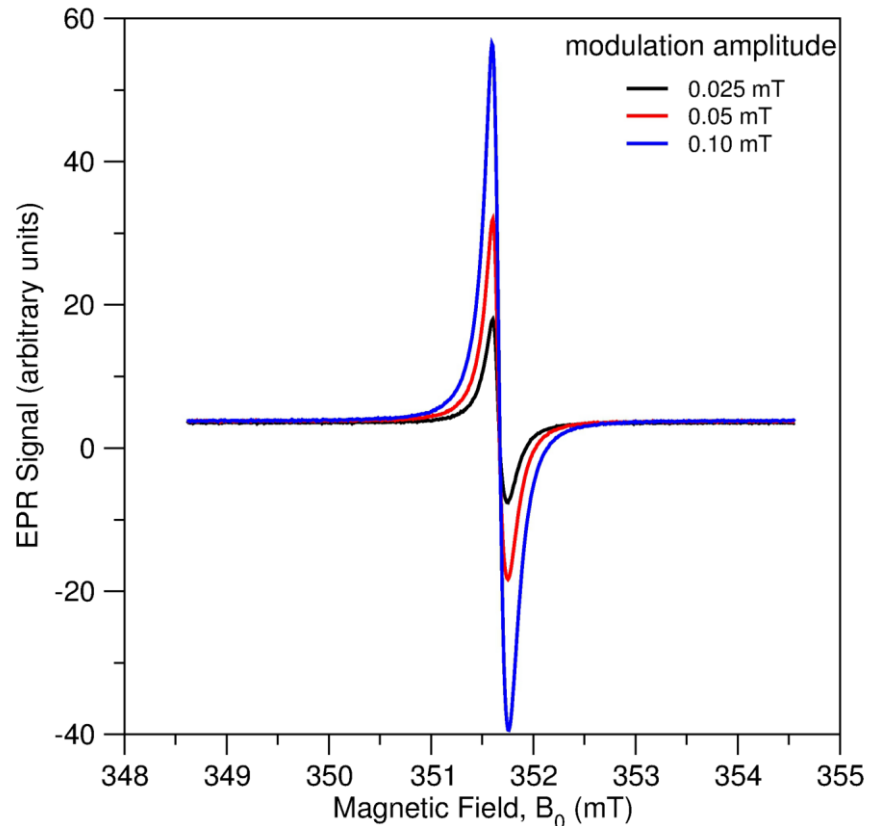


Frequency (GHz)	Frequency Band	Field for $g=2.0023$ (T)
1.2	L	0.043
2.4	S	0.086
9.5	X	0.34
34	Q	1.2
95	W	3.4
263	mm-band	9.4

# Field Modulation

**Direct detection of low-power microwave frequencies is hard due to noise.**

- Virtually all commercial continuous wave EPR instruments use field modulation



more  $\neq$  better

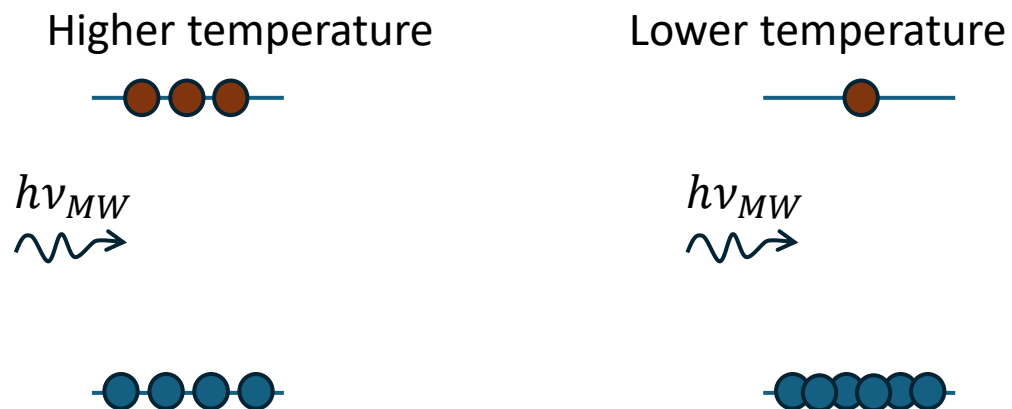


# Why cryogenic temperatures?

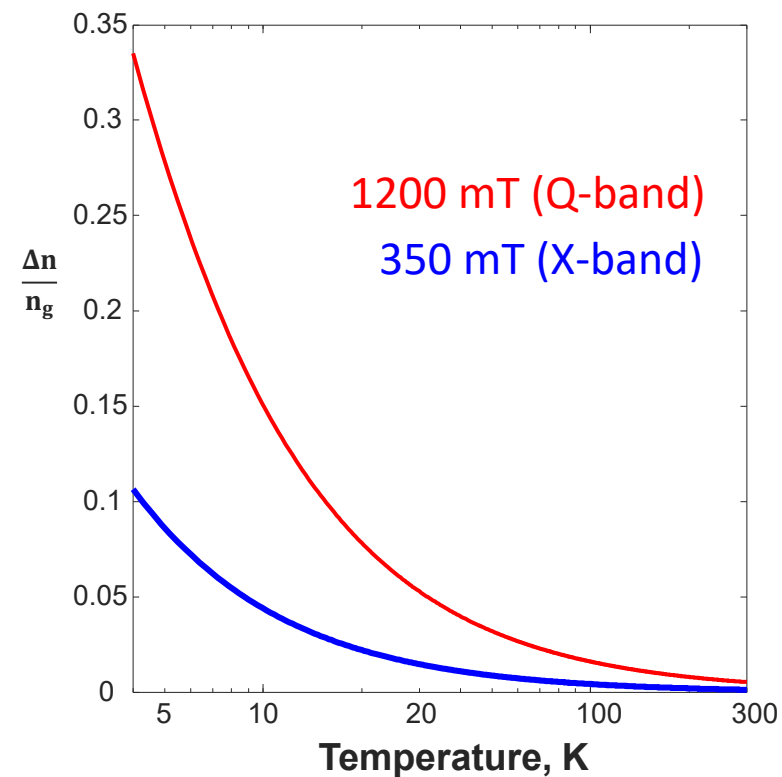
## 1. Combat Boltzmann population distribution

EPR signal depends on population difference between spin states:

$$\frac{n_{excited}}{n_{ground}} = \exp\left(-\frac{g\beta_e B_0}{kT}\right)$$



Lower temperature = more signal (not the end of the story, though)

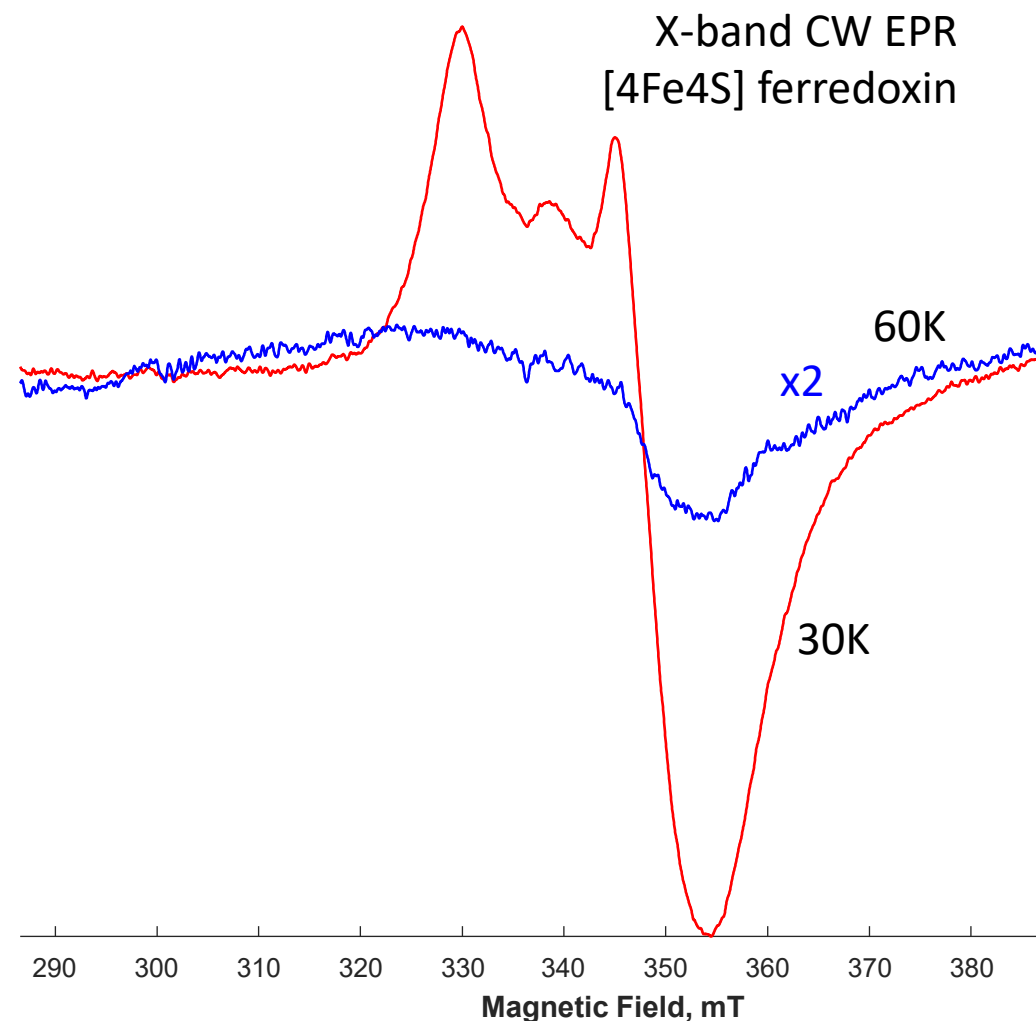


# Why cryogenic temperatures?

## 2. Combat fast relaxation (line broadening)

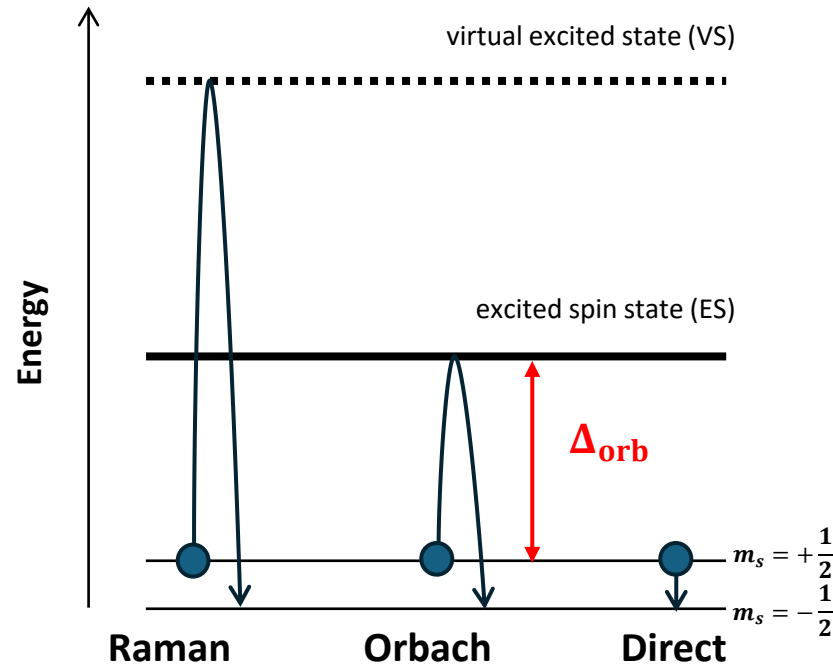
$$\Delta B \propto \frac{1}{T_1} + \frac{1}{T_2}$$

$T_2$  normally has a moderate temperature behavior  
 $T_1$  has strong temperature behavior



# Relaxation of coupled systems

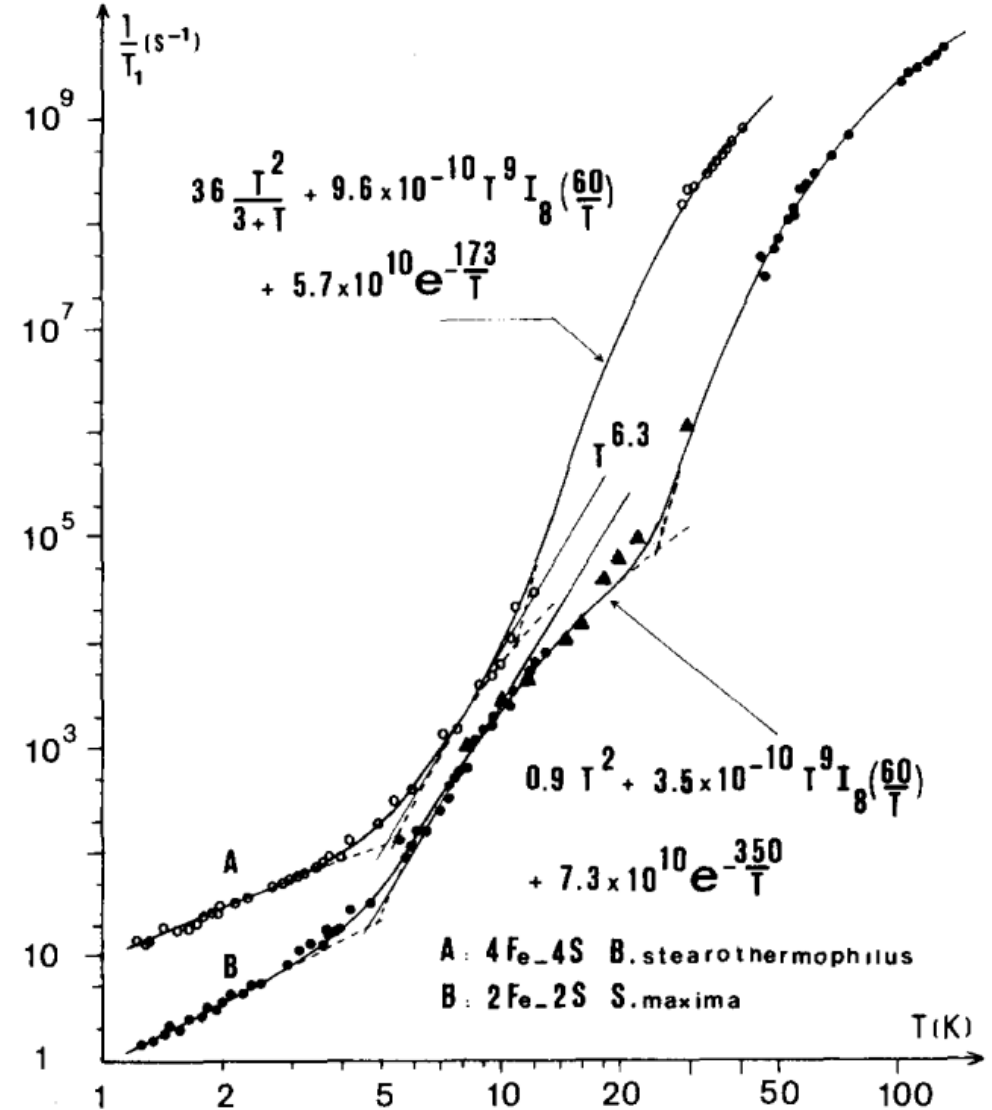
## Multiple pathways for spin-lattice relaxation ( $T_1$ )



$$\left(\frac{1}{T_1}\right)_{\text{raman}} \propto T^9$$

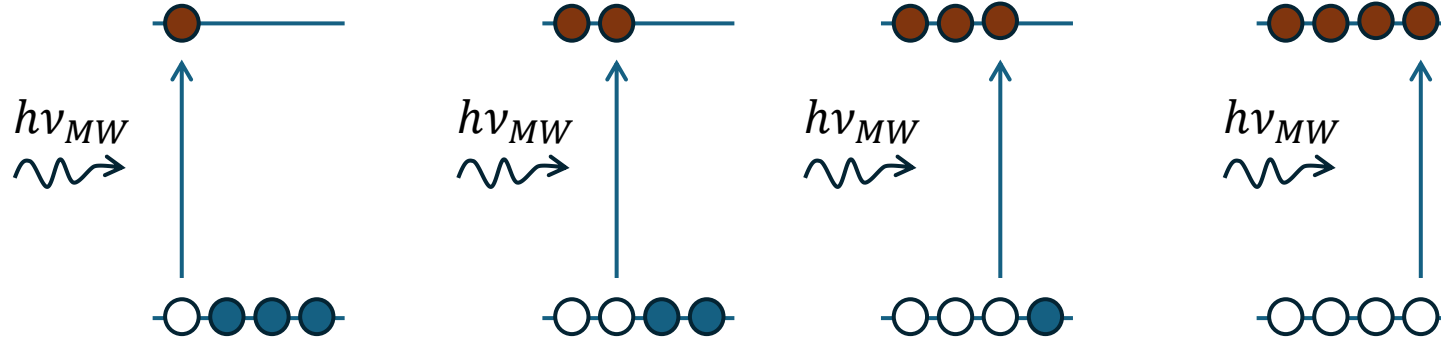
$$\left(\frac{1}{T_1}\right)_{\text{orbach}} \propto \frac{1}{\exp\left(\frac{\Delta_{orb}}{T}\right) - 1}$$

$$\left(\frac{1}{T_1}\right)_{\text{direct}} \propto T^1$$



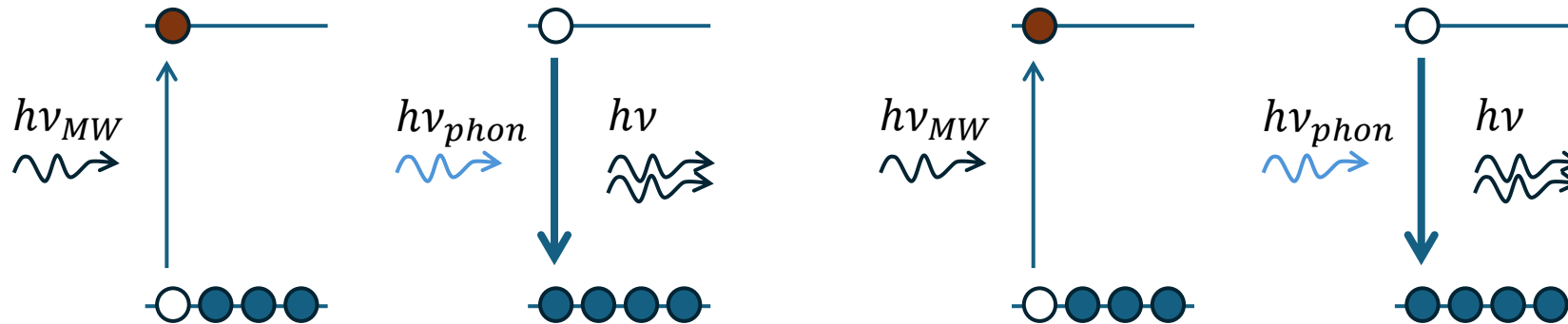
# Power saturation

## Extreme case #1. Insignificantly slow relaxation



Nothing left to do  
No more quants of energy can  
be absorbed  
**SATURATION**

## Extreme case #2. FAST (direct) relaxation



We will never run out  
of things to do

# Power Saturation

MW amplitude ( $\sqrt{P_{MW}}$ )  $\approx$  # of quants of energy, i.e., how quickly we can “pump” energy into the system

If relaxation is **very fast**

more power  $\rightarrow$  more absorption events  $\rightarrow$  **more** EPR signal

If relaxation is **very slow**

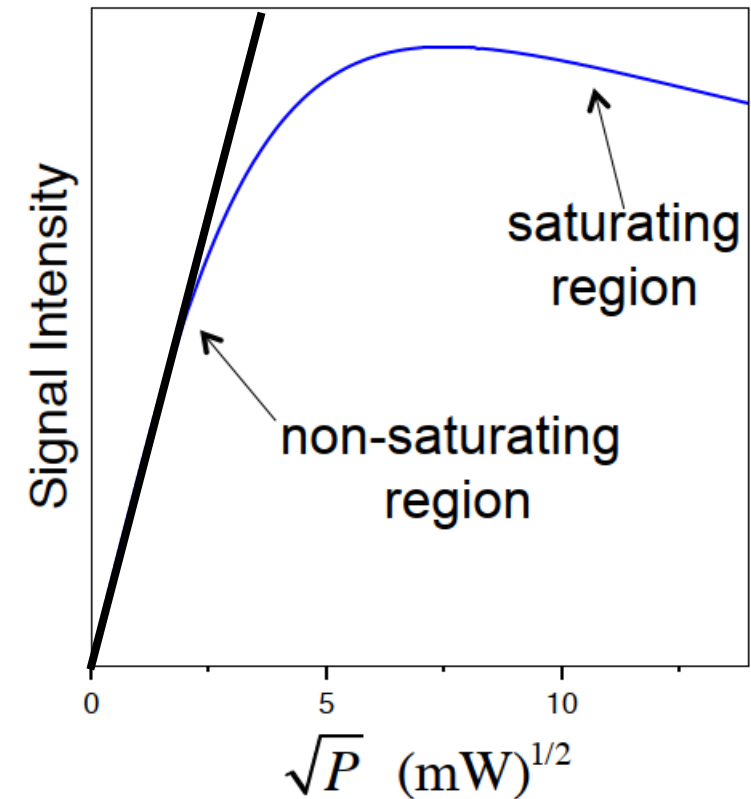
more power  $\rightarrow$  faster transition saturates  $\rightarrow$  **less** EPR signal

$$A = a \frac{\sqrt{P} T_2^2}{\left(1 + \frac{P}{P_{0.5}}\right)^{1/2}}$$

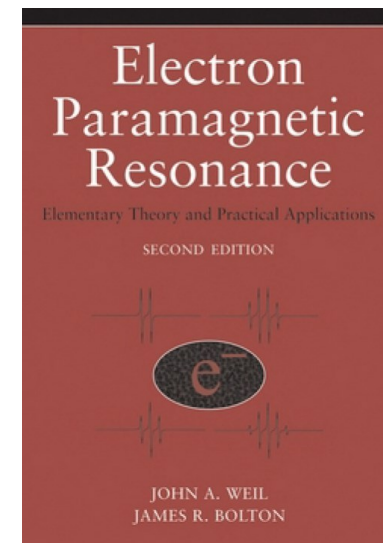
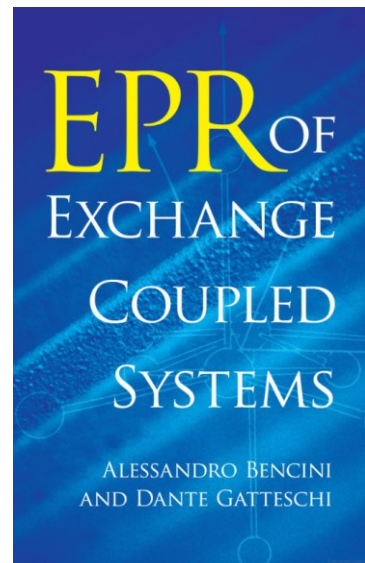
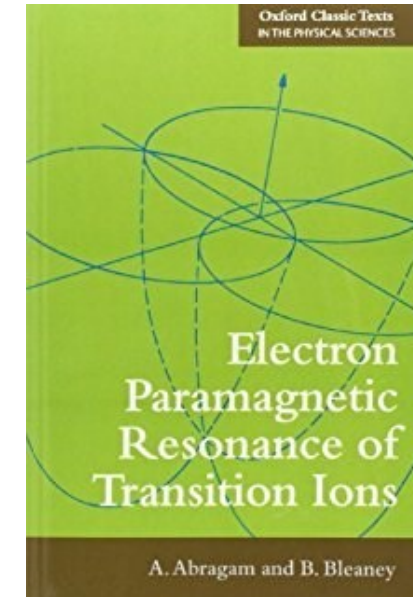
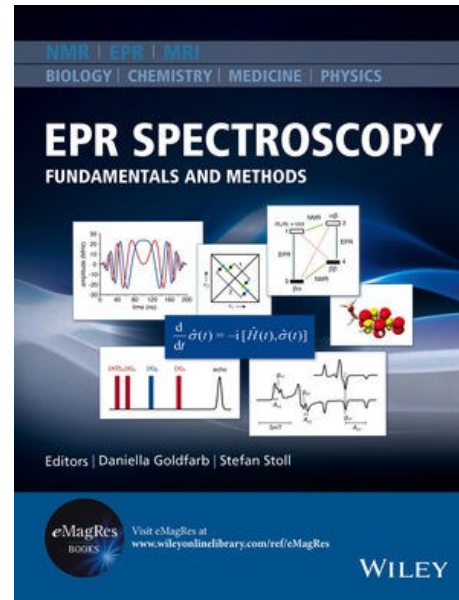
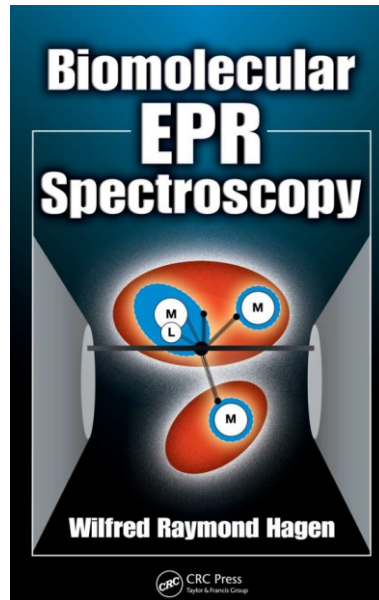
$a$  – instrument  
dependent constant

$$P_{0.5} = (\gamma_e^2 T_1 T_2)^{-1}$$

Lower temperature  $\rightarrow$  Slower relaxation  $\rightarrow$  Need less power to saturate



# Literature







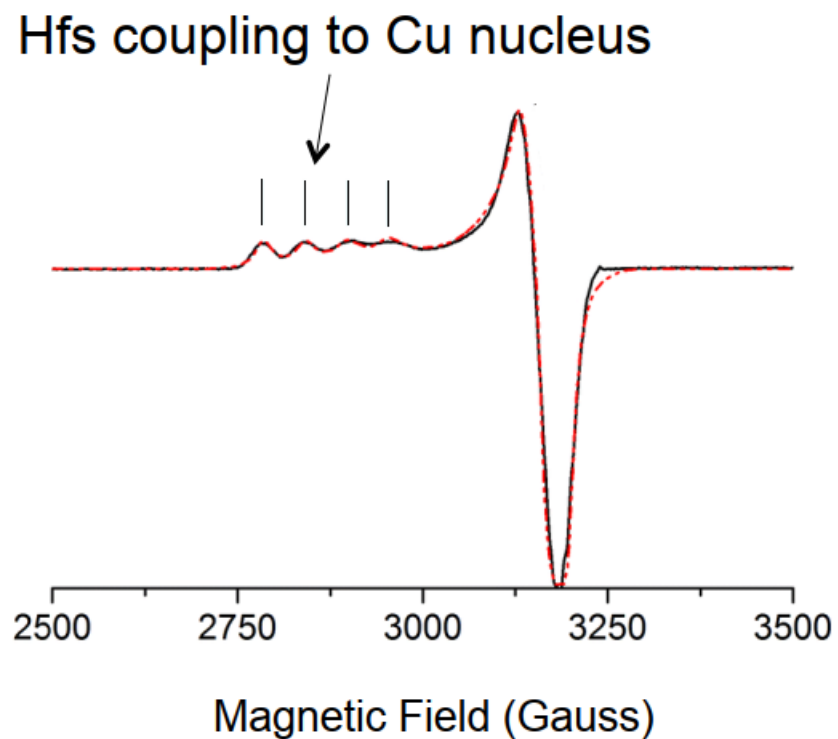


## Resolved Hyperfine Coupling

Azurin:

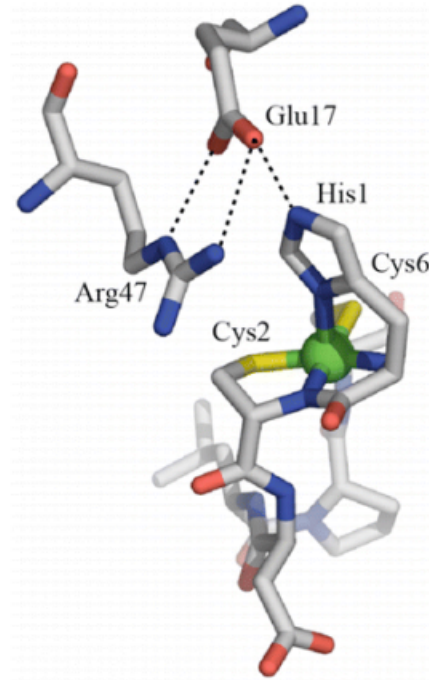
In proteins containing Cu(II) hyperfine coupling to the  $I=3/2$  copper nucleus is observed.

Usually it is only resolved on the  $g_{\perp}$  component

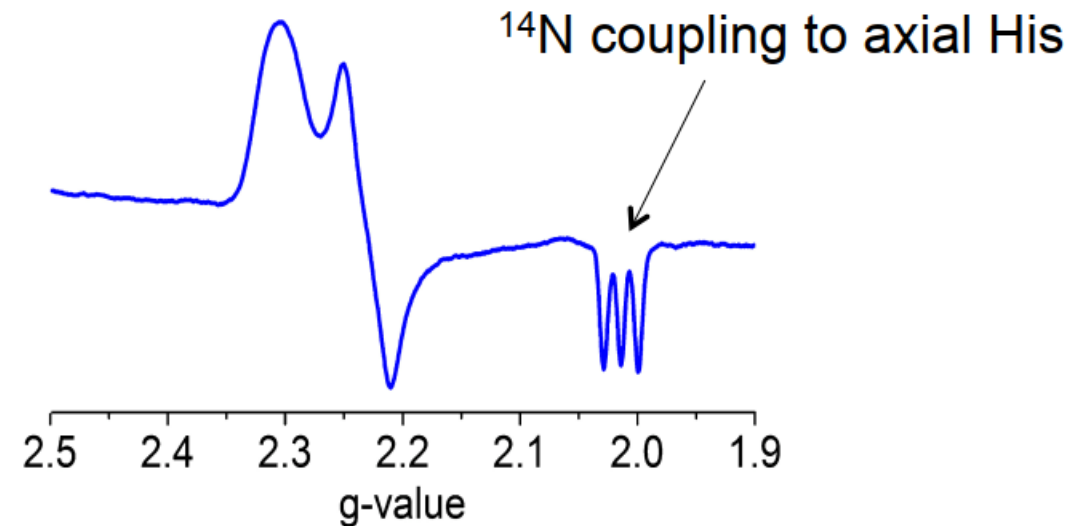


## Resolved Hyperfine Coupling

Ni Superoxide dismutase:



Hyperfine coupling to ligands is also sometimes resolved on one or more of the g-tensor components



# Let's talk about instrumentation



“EPR spectrometer” according to AI