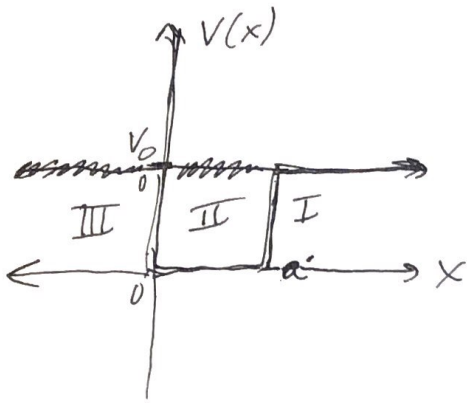


HW 5 Solutions

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 Quantum Mechanics I
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1. Calculate energy & wavefunction of:



$$V(x) = \begin{cases} \infty & x < 0 \\ 0 & 0 \leq x < a \\ V_0 & x \geq a \end{cases}$$

Start in section II, apply Schrödinger eqn:

$$\boxed{\text{II}} \quad \frac{d^2\psi}{dx^2} = -\frac{2mE}{\hbar^2} \psi \quad (V=0)$$

General solution $\psi_{\text{II}} = A \sinh(kx) + B \cos(kx)$

Apply BC $\psi_{\text{II}}(0) = \psi_{\text{III}}(0) = 0$ bc infinite wall

$$\Rightarrow \psi_{\text{II}}(0) = A \sinh(0) + B \cos(0) = B, \text{ so } B = 0$$

$$\Rightarrow \boxed{\psi_{\text{II}}(x) = A \sinh(kx)}$$

Now Section I:

$$\boxed{\text{I}}: \quad \frac{d^2\psi}{dx^2} = -\frac{2m(E-V_0)}{\hbar^2} \psi \quad (V=V_0)$$

$l^2 \rightarrow$ Note that l^2 includes the minus sign, because $E < V_0$ and we need $l^2 > 0$

General solution: $\psi_{\text{I}} = C e^{lx} + D e^{-lx}$

As $x \rightarrow \infty$, first term would blow up, so must have $C=0$; then $\boxed{\psi_{\text{I}}(x) = D e^{-lx}}$

And of course:

$$\boxed{\text{III}} \quad \boxed{\psi_{\text{III}}(x) = 0}$$

Now we can apply continuity @ $x=a$:

$$\psi_{\text{II}}(a) = \psi_{\text{I}}(a) \Rightarrow A \sinh(ka) = D e^{-la} \quad (1)$$

And continuity of derivative @ $x=a$:

$$\psi'_{\text{II}}(a) = \psi'_{\text{I}}(a) \Rightarrow A k \cosh(ka) = -D l e^{-la} \quad (2)$$

Let's divide $\frac{(2)}{(1)}$: $k \cot(ka) = -l \Rightarrow \boxed{\cot(ka) = -\frac{l}{k}}$

Now, l & k are both only dependent on E + constants, so this is a transcendental equation that we could solve graphically. We can rewrite it more nicely with the following change of variables:

see that $k^2 + l^2 = \frac{2mE}{\hbar^2} + \left(-\frac{2m(E-V_0)}{\hbar^2}\right) = \frac{2mV_0}{\hbar^2}$, so call $z_0 = \frac{a}{\hbar} \sqrt{2mV_0} = a \sqrt{k^2 + l^2}$
 and $z = ka = \frac{a}{\hbar} \sqrt{2mE}$

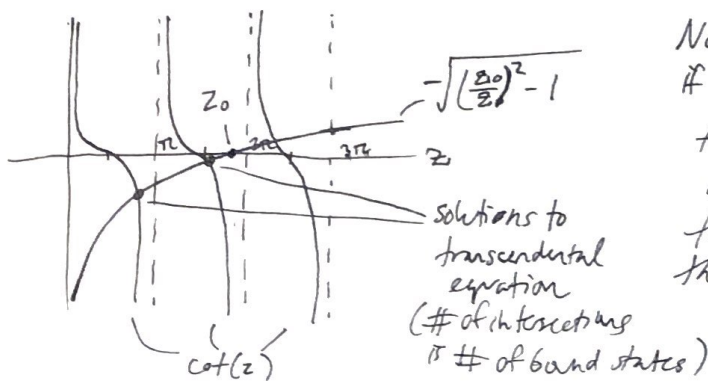
then we have $z_0^2 - z^2 = \frac{a^2}{\hbar^2} 2mV_0 - \frac{a^2}{\hbar^2} 2mE = \frac{a^2}{\hbar^2} 2m(V_0 - E) = a^2 l^2$

$\Rightarrow \sqrt{z_0^2 - z^2} = la$

Now rewrite $\cot(ka) = -\frac{l}{k}$ as:

$$\cot(z) = -\frac{l a}{k a} = -\frac{\sqrt{z_0^2 - z^2}}{z} = -\sqrt{\left(\frac{z_0}{z}\right)^2 - 1}$$

This is now clearly just a function of z . We can sketch the two sides:



Notice that there are only real solutions if $\sqrt{\left(\frac{z_0}{z}\right)^2 - 1} > 0$, so $z_0 > z$. Going back to the definitions, this is $\frac{a}{\hbar} \sqrt{2mV_0} > \frac{a}{\hbar} \sqrt{2mE}$, or $V_0 > E \Rightarrow E < V_0$; because of course the energies have to be within the depth of the well for bound states!

So our final solution is:

$$\psi(x) = \begin{cases} 0 & x < 0 \\ A \sinh(kx) & 0 \leq x < a \\ D e^{-lx} & x > a \end{cases}$$

where k & l are found by solving $\cot(ka) = -\frac{l}{k}$

$$E = \frac{\hbar^2 k^2}{2m} \quad (\text{by rearranging } k^2 \text{ definition})$$

We could relate A & D using the continuity BC at $x=a$, and then find the remaining constant via normalization, but the problem said we didn't have to worry about normalizing :)