

QM I Midterm I Solutions

TA: Kate Stray-Fisher
 Prof: Javad Shebani
 Quantum Mechanics I
 Fall 2021

1. 3-level system: $\hat{H} = \epsilon \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 2i \\ 0 & -2i & 0 \end{pmatrix}$

a) Yes, \hat{H} is Hermitian. $\hat{H}^* = \epsilon \begin{pmatrix} 0 & i & 0 \\ -i & 0 & -2i \\ 0 & 2i & 0 \end{pmatrix}$, $(\hat{H}^*)^T = \epsilon \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 2i \\ 0 & -2i & 0 \end{pmatrix}$
 we can see that $(\hat{H}^*)^T = \hat{H}$, Hermitian definition ✓

b) In Dirac notation, we can translate: $H_{ij} = \langle e_i | \hat{H} | e_j \rangle$ where e_i are the 3 states, $|1\rangle, |2\rangle, |3\rangle$. Then we can read off the matrix elements, seeing that all the 0 elements will disappear:

$$\hat{H} = \epsilon [-i |1\rangle\langle 2| + i |2\rangle\langle 1| + 2i |2\rangle\langle 3| - 2i |3\rangle\langle 2|]$$

To check this, you could go the other direction and reconstruct \hat{H} by sandwiching this \hat{H} in Dirac form between all possible pairs of states.

c) Eigenvalues:

$$\det \begin{pmatrix} -\lambda & -i\epsilon & 0 \\ i\epsilon & -\lambda & 2i\epsilon \\ 0 & -2i\epsilon & -\lambda \end{pmatrix} = -\lambda(\lambda^2 + \underbrace{4i^2\epsilon^2}_{-4\epsilon^2}) + \underbrace{i\epsilon(-\lambda i\epsilon)}_{\lambda\epsilon^2} = -\lambda^3 + 5\lambda\epsilon^2 = \lambda(5\epsilon^2 - \lambda^2) = 0$$

$$\Rightarrow \boxed{\lambda = 0, \sqrt{5}\epsilon, -\sqrt{5}\epsilon}$$

d) Eigenvector for largest eigenvalue, $\lambda = \sqrt{5}\epsilon$:

$$\hat{H}\psi = \lambda\psi$$

$$\epsilon \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 2i \\ 0 & -2i & 0 \end{pmatrix} \begin{pmatrix} a \\ b \\ c \end{pmatrix} = \sqrt{5}\epsilon \begin{pmatrix} a \\ b \\ c \end{pmatrix}$$

System of eqns (cancel ϵ):

$$-ib = \sqrt{5}a \rightarrow b = i\sqrt{5}a$$

$$ca + 2ic = \sqrt{5}b$$

$$-2ib = \sqrt{5}c \rightarrow c = -\frac{2ib}{\sqrt{5}} = -\frac{2i(i\sqrt{5}a)}{\sqrt{5}} = 2a$$

Choose $a=1$, then: $v = \begin{pmatrix} 1 \\ i\sqrt{5} \\ 2 \end{pmatrix} A$ (A is normalization constant)

$$\text{Normalize: } 1 = (v^*)^T \cdot v = A^2 (1 -i\sqrt{5} \ 2) \begin{pmatrix} 1 \\ i\sqrt{5} \\ 2 \end{pmatrix} = A^2 [1 + 5 + 4] = 10A^2 = 1 \Rightarrow A = \frac{1}{\sqrt{10}}$$

$$\boxed{v = \frac{1}{\sqrt{10}} \begin{pmatrix} 1 \\ i\sqrt{5} \\ 2 \end{pmatrix}}$$

e) Normalize $\psi = \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$: $1 = |\psi|^2 = A^2 (1 \ 1 \ 1) \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} = A^2 (1^2 + 1^2 + 1^2) = 3A^2 = 1 \Rightarrow A = \frac{1}{\sqrt{3}}$

$$\psi = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}, \quad \psi = \frac{1}{\sqrt{3}} (|1\rangle + |2\rangle + |3\rangle) \text{ in Dirac notation}$$

f) $\langle H \rangle = \langle \psi | H | \psi \rangle = \frac{\epsilon^2}{3} (1 \ 1 \ 1) \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 2i \\ 0 & -2i & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} = \frac{\epsilon^2}{3} (1 \ 1 \ 1) \begin{pmatrix} -i \\ 3i \\ -2i \end{pmatrix} = \frac{\epsilon^2}{3} (-i + 3i - 2i) = 0 = \langle H \rangle$

g) $\sigma_H = \sqrt{\langle H^2 \rangle - \langle H \rangle^2}$; $\langle H^2 \rangle = \langle \psi | H H | \psi \rangle = \frac{\epsilon^2}{3} (1 \ 1 \ 1) \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 2i \\ 0 & -2i & 0 \end{pmatrix} \begin{pmatrix} -i \\ 3i \\ -2i \end{pmatrix} = \frac{\epsilon^2}{3} (1 \ 1 \ 1) \begin{pmatrix} 3 \\ 5 \\ 6 \end{pmatrix} = \frac{14\epsilon^2}{3}$
 $\Rightarrow \sigma_H = \sqrt{\frac{14\epsilon^2}{3} - 0} = \boxed{\sqrt{\frac{14}{3}} \epsilon = \sigma_H}$

2) Particle in infinite quantum well b/w $0 \leq x \leq a$:

$$\Psi(x, t=0) = \sqrt{\frac{8}{5a}} \left[1 + \cos\left(\frac{\pi x}{a}\right) \right] \sinh\left(\frac{\pi x}{a}\right)$$

a) $\Psi(x, t=t_0) = ?$

First notice that we can rewrite Ψ in terms of just sines, using $\sin(2\theta) = 2 \sin\theta \cos\theta$:

$$\Psi(x, t=0) = \sqrt{\frac{8}{5a}} \left[\sinh\left(\frac{\pi x}{a}\right) + \sinh\left(\frac{\pi x}{a}\right) \cos\left(\frac{\pi x}{a}\right) \right] = \sqrt{\frac{8}{5a}} \left[\sinh\left(\frac{\pi x}{a}\right) + \frac{1}{2} \sin\left(\frac{2\pi x}{a}\right) \right]$$

these two terms are just the $n=1$ & $n=2$ states of the infinite well:

$$\text{inf. well: } \Psi_n(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right), \quad \text{and } E_n = \frac{\hbar^2 \pi^2 n^2}{2ma^2}$$

$$\text{So } \Psi(x, t=0) = \sqrt{\frac{8}{5a}} \left[\sqrt{\frac{a}{2}} \Psi_1 + \sqrt{\frac{a}{2}} \frac{1}{2} \Psi_2 \right] = \frac{2}{\sqrt{5}} \left[\Psi_1 + \frac{1}{2} \Psi_2 \right]$$

To include time dependence, use: $\Psi(x, t) = \sum c_n \Psi_n e^{-i\hbar E_n t / \hbar}$ where E_n 's are these and we can read off that $c_1 = \frac{2}{\sqrt{5}}$ & $c_2 = \frac{1}{\sqrt{5}}$:

$$\boxed{\Psi(x, t) = \frac{2}{\sqrt{5}} \left(\Psi_1 e^{-\frac{i\hbar E_1 t}{\hbar}} + \frac{1}{2} \Psi_2 e^{-\frac{i\hbar E_2 t}{\hbar}} \right)}$$

b) average energy $\langle H \rangle = \langle \Psi | H | \Psi \rangle = \sum |c_n|^2 E_n$ (very useful eqn, see Griffiths!)

$$\Rightarrow \langle H \rangle = \frac{4}{5} \left[\frac{\pi^2 \hbar^2}{2ma^2} + \left(\frac{1}{2}\right)^2 \frac{4\pi^2 \hbar^2}{2ma^2} \right] = \frac{4\pi^2 \hbar^2}{5ma^2}$$

The energy is independent of time so $\langle H \rangle$ at $t=0$ is the same as $\langle H \rangle$ at $t=t_0$.

$$c) P(0 \leq x \leq \frac{a}{2}, t=t_0) = \int_0^{a/2} |\Psi(x, t=t_0)|^2 dx$$

$$= \frac{4}{5} \int_0^{a/2} \left(\Psi_1^* e^{i\hbar E_1 t / \hbar} + \frac{1}{2} \Psi_2^* e^{i\hbar E_2 t / \hbar} \right) \left(\Psi_1 e^{-i\hbar E_1 t / \hbar} + \frac{1}{2} \Psi_2 e^{-i\hbar E_2 t / \hbar} \right) dx$$

$$= \frac{4}{5} \int_0^{a/2} \left(\underbrace{\Psi_1^* \Psi_1}_{I_1} + \underbrace{\frac{1}{4} \Psi_2^* \Psi_2}_{I_2} + \underbrace{\frac{1}{2} \Psi_1^* \Psi_2 e^{-i\hbar(E_2-E_1)t/\hbar}}_{I_3} + \underbrace{\frac{1}{2} \Psi_2^* \Psi_1 e^{i\hbar(E_2-E_1)t/\hbar}}_{I_3} \right) dx$$

$$I_1 = \left(\frac{2}{a}\right) \int_0^{a/2} \sin^2\left(\frac{\pi x}{a}\right) dx = \left(\frac{2}{a}\right) \int_0^{a/2} \left[\frac{1}{2} - \frac{1}{2} \cos\left(\frac{2\pi x}{a}\right) \right] dx = \left(\frac{2}{a}\right) \left[\frac{1}{2}x - \frac{1}{2} \left(\frac{a}{2\pi}\right) \sin\left(\frac{2\pi x}{a}\right) \right]_0^{a/2} = \frac{1}{2}$$

$$I_2 = \left(\frac{2}{a}\right) \frac{1}{4} \int_0^{a/2} \sin^2\left(\frac{2\pi x}{a}\right) dx = \frac{1}{2a} \int_0^{a/2} \left[\frac{1}{2} - \frac{1}{2} \cos\left(\frac{4\pi x}{a}\right) \right] dx = \frac{1}{2a} \left[\frac{1}{2}x - \frac{1}{2} \left(\frac{a}{4\pi}\right) \sin\left(\frac{4\pi x}{a}\right) \right]_0^{a/2} = \frac{1}{8}$$

$$I_3 = \frac{1}{2} \left(e^{i\hbar(E_2-E_1)t/\hbar} + e^{-i\hbar(E_2-E_1)t/\hbar} \right) \left(\frac{2}{a}\right) \int_0^{a/2} \sinh\left(\frac{\pi x}{a}\right) \sin\left(\frac{2\pi x}{a}\right) dx$$

$\cos\left(\frac{\hbar(E_2-E_1)t}{\hbar}\right)$, by Euler's identity (using $\sinh x \sin y = \frac{1}{2} [\cos(x-y) - \cos(x+y)]$):

$$= \left(\frac{2}{a}\right) \cos\left(\frac{\hbar(E_2-E_1)t}{\hbar}\right) \int_0^{a/2} \frac{1}{2} \left[\cos\left(\frac{\pi x}{a}\right) - \cos\left(\frac{3\pi x}{a}\right) \right] dx = \frac{1}{a} \cos\left(\frac{\hbar(E_2-E_1)t}{\hbar}\right) \left[\left(\frac{a}{\pi}\right) \sin\left(\frac{\pi x}{a}\right) - \left(\frac{a}{3\pi}\right) \sin\left(\frac{3\pi x}{a}\right) \right]_0^{a/2}$$

$$= \frac{2a}{3\pi} \cos\left(\frac{\hbar(E_2-E_1)t}{\hbar}\right) \frac{3\pi^2 \hbar^2}{2ma^2} \left(\frac{a}{\pi} + \frac{a}{3\pi} \right) = \frac{4a}{3\pi}$$

$$\boxed{P = \frac{4}{5} \left(\frac{1}{2} + \frac{1}{8} + \frac{4a}{3\pi} \cos\left(\frac{\hbar(E_2-E_1)t}{\hbar}\right) \right) = \frac{1}{2} + \frac{16}{15\pi} \cos\left(\frac{\hbar(E_2-E_1)t}{\hbar} \frac{3\pi^2 \hbar^2}{2ma^2}\right)}$$