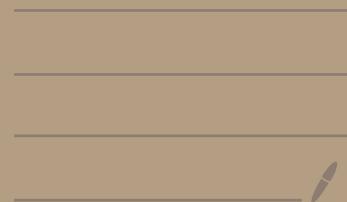


Long Time Behavior and Local Regularity for solutions to the Navier-Stokes equations

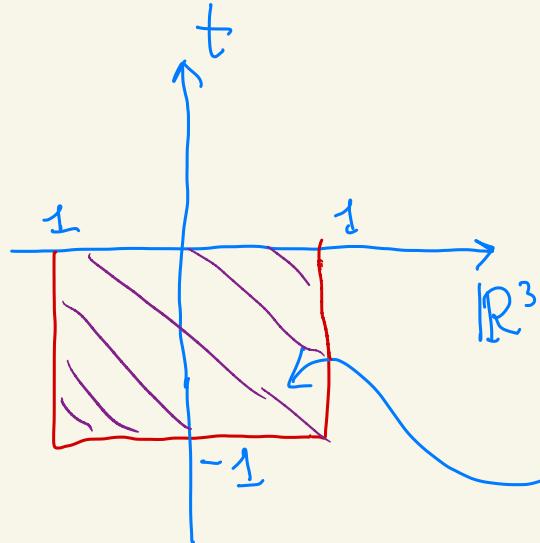
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Local Regularity Setting



$$Q = \Omega \times [-1, 0]$$

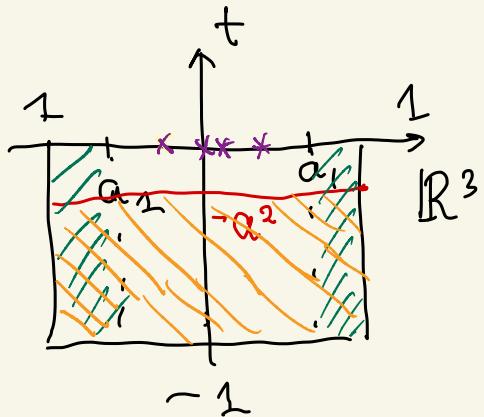
$$\begin{cases} \partial_t w + w \cdot \nabla w - \Delta w = -\nabla P \\ \operatorname{div} w = 0 \end{cases}$$

in Q

First time singularity scenario

1. Energy solution: $w \in L_{2,\infty}(Q) = L_\infty(-1,0; L_2(B))$,
 $\nabla w \in L_2(Q)$, $r \in L_{3/2}(Q)$

2. Mimic the first time singularity:



$$u \in L_\infty(\{x_1 < |x| < 1\} \times]-1, 1[)$$

$$u \in L_\infty(B \times]-1, -\alpha^2[)$$

for any $0 < \alpha < 1$

3.

Type I singularity

for example

$$|w(x, t)| \leq \frac{c_d}{|x| + \sqrt{-t}}$$

Potential singularity is $z = (x, t) = (0, 0)$

Question: Is $z = 0$ a regular point
of w ? In other words,
is it true $w \in L_\infty(Q)$?

The answer is Yes if
 c_s is sufficiently small

Question 2. How small?

Our way to attack the problem
is via duality.

Assume that $\underline{z} = 0$ is a singular point of w , then $w \notin L_\infty(Q(\delta))$, $\forall 0 < \delta < 1$

$$\textcircled{1} + \textcircled{2} \implies \exists \text{ non-trivial}$$

mild bounded ancient solution

$\bar{u} : \mathbb{R}^3 \times]-\infty, 0[$ such that

- (i) $\bar{u} \in L_\infty(\mathbb{R}^3 \times]-\infty, 0[)$ ($|\bar{u}| \leq 1$)
- (ii) $\exists \bar{p} \in L_\infty(-\infty, 0; \mathcal{BMO})$
- (iii) $\partial_t \bar{u} + \bar{u} \cdot \nabla \bar{u} - \Delta \bar{u} = -\nabla \bar{p}$ } $\mathbb{R}^3 \times]-\infty, 0[$
 $\operatorname{div} \bar{u} = 0$
- (iv) $|\bar{u}(0)| = 1.$!!!

Properties of mild bounded ancient solutions :

\bar{u} is of class C^∞ , all derivatives of \bar{u} and $\nabla \bar{p}$ are bounded.

Decay ?

$$\textcircled{3} \Rightarrow |\bar{u}(x, t)| \leq \frac{C_d}{|x| + \sqrt{-t}}$$

$$\forall x \in \mathbb{R}^3, \forall t \leq 0$$

It is more convenient
to work with

$$u(x, t) = \bar{u}(x, -t) \quad \text{for } t > 0.$$

$$p(x, t) = \bar{p}(x, -t)$$

$$(i) \quad \left. \begin{array}{l} -\partial_t u + u \cdot \nabla u - \Delta u = -\nabla p \\ \operatorname{div} u = 0 \end{array} \right\} \text{in } Q_+ = [R^3 \times]0, \infty[$$

$$(ii) \quad |u(x, t)| \leq 1, \quad p \in L_\infty(0, \infty; BMO)$$

$$(iii) \quad |u(x, t)| \leq \frac{C_d}{\sqrt{1+t}}$$

$$(iv) \quad |u(0)| = 1.$$

Question: $(i) + (ii) + (iii) \Rightarrow u \equiv 0$?
 Q_+

Duality approach

I. $\partial_t v - u \cdot \nabla v - \Delta v - \nabla q = -\operatorname{div} F$
(or $\partial_t v - u \cdot \nabla v + v \cdot \nabla u - \nabla q = -\operatorname{div} F$),
 $\operatorname{div} v = 0 \quad \text{in } Q_+ = \mathbb{R}^3 \times]0, \infty[$

$$v \Big|_{t=0} = 0$$

$$F \in (F_{ij}) \in C^\infty(\mathbb{R}^3), \quad F = -F^\top$$

$$\int_{Q_+} u \cdot \operatorname{div} F dx = - \lim_{T \rightarrow \infty} \int_{\mathbb{R}^3} u(x, T) \cdot v(x, T) dx$$

$$|u(x, t)| \leq \frac{c_d}{|x| + \sqrt{t}} \Rightarrow$$

$$\left| \int_{\mathbb{R}^3} u(x, t) \cdot v(x, t) dx \right| \leq \frac{c_d}{\sqrt{t}} \|v(\cdot, t)\|_{L_1(\mathbb{R}^3)}$$

→ 0 ?

Theorem. $\forall m = 0, 1, 2, \dots$

$$\|v(\cdot, t)\|_1 \leq c(m, C_d, F) \frac{\sqrt{t}^{3/2}}{\ln^m(t+\ell)}$$

and

$$\|v(\cdot, t)\|_2 \leq c(m, C_d, F) \frac{1}{\ln^m(t+\ell)}$$

$$\forall t \geq 0$$

$$\text{II. } \partial_t v - u \cdot \nabla v - \Delta v - \nabla q = 0$$

(or $\partial_t v - u \cdot \nabla v + v \cdot \nabla u - \Delta v - \nabla q = 0$)

$$\operatorname{div} v = 0 \quad \text{in } Q_+$$

$$v \Big|_{t=0} = v_0 \in C_{0,0}^\infty(\mathbb{R}^3)$$

$$C_{0,0}^\infty(\mathbb{R}^3) = \{ v \in C_0^\infty(\mathbb{R}^3) : \operatorname{div} v = 0 \}$$

$$\int_{\mathbb{R}^3} u(x, t) \cdot v(x, t) dx = \int_{\mathbb{R}^3} u(x, 0) \cdot v_0(x) dx$$

$\forall t \geq 0$

as $t \rightarrow \infty$

?

↓

0

yes
 $\implies u(x, 0) = 0$ (but $|u(0)|=1$)

Theorem. $\forall m = 0, 1, 2, \dots$

$$\|v(\cdot, t)\|_1 \leq c(m, C_d, v_0) \frac{\sqrt{t}^{3/2}}{\ln^m(t+\ell)}$$

and

$$\|v(\cdot, t)\|_2 \leq c(m, C_d, v_0) \frac{1}{\ln^m(t+\ell)}$$

$$\forall t \geq 1$$

So, we need to study the decay of $\|v(\cdot, t)\|_{L_p}$ if $t \gg 1$
 and $p = 1$ or $p = 2$

What would we have for
 heat equation : $\partial_t v - \Delta v = 0, v|_{t=0} = u_0$

$$\Rightarrow \|v(\cdot, t)\|_{L_p} \leq \|u_0\|_{L_p} \quad \text{if } p \geq 1.$$

$$\|u_0\|_p$$

Why it is problem to estimate $\|v(\cdot, t)\|_1$? (Heuristic only)

Heat equation:

$$\partial_t v - \Delta v = 0 \quad * \quad \frac{v}{|v|}$$

$$\Rightarrow \partial_t \int_{\mathbb{R}^3} |v|^2 dx \leq 0$$

Stokes equations:

$$\partial_t v - \Delta v = u \cdot \nabla v + \nabla q$$

$$\Delta q = \operatorname{div}(u \cdot \nabla v)$$

If $u(\cdot, t) \cdot \nabla v(\cdot, t) \in L^1$ (Hardy space)

$$\Rightarrow \nabla q(\cdot, t) \in L^1$$

$$\Rightarrow \partial_t \int_{\mathbb{R}^3} |v| dx \leq c \|u(\cdot, t) \cdot \nabla v(\cdot, t)\|_1$$

However,

$$\|u \cdot \nabla v\|_{L^1} \leq c \left(\|u\|_2 \|\nabla v\|_2 + \|\nabla u\|_2 \|v\|_2 \right)$$

$$\sup_{0 \leq t < \infty} \|u(\cdot, t)\|_2 - ?$$

Perturbation method (unfortunately) ↓
 Later on

$$\partial_t v - \cancel{u \cdot \nabla v}$$

$$-\Delta v - \nabla q = 0$$

This equation can be rewritten

$$\partial_t v - \operatorname{div}(v \otimes u) - \Delta v - \nabla q = 0$$

with

$$\Delta q = -\operatorname{div} \operatorname{div}(v \otimes u)$$

energy estimate

$$\left\| v(\cdot, t) \right\|_{L_2^2}^2 + 2 \int_0^t \int_{\mathbb{R}^3} |\nabla v|^2 dx dt$$

$$= \| u_0 \|^2_{L_2^2}$$

So, we have perturbed heat equation

$$\partial_t v - \Delta v = \operatorname{div} F$$

where

$$F = \mathcal{X}(u, v) + q \mathbb{I}$$

and

$$\mathcal{X}(u, v) = v \otimes u$$

and

$$v \Big|_{t=0} = u_0 \quad \text{in } \mathbb{R}^3$$

Weak point (perturbation)

$$\partial_t v + \mathbf{b} \cdot \nabla v - \Delta v = 0 \text{ with } \operatorname{div} \mathbf{b} = 0, v|_{t=0} = u_0$$

$$\Rightarrow \int_{\mathbb{R}^3} v(x, t) dx = \int_{\mathbb{R}^3} u_0(x, \epsilon) dx$$

However, if we use representation formula

$$v(x, t) = \int_{\mathbb{R}^3} \Gamma(x-y, t) u_0(y) dy + \int_0^t \int_{\mathbb{R}^3} \nabla \Gamma(x-y, t-\tau) \cdot \mathbf{b}(y, \tau) v(y, \tau) dy d\tau$$

Positive results are coming from smallness
of \mathbf{b} !!

$$v_i(x, t) = v_i^1(x, t) + v_i^2(x, t)$$

where

$$v_i^1(x, t) = \int_{\mathbb{R}^3} \Gamma(x-y, t) u_0(y, t) dy \Leftrightarrow \begin{aligned} \partial_t v^1 - \Delta v^1 &= 0 \\ v^1|_{t=0} &= u_0 \end{aligned}$$

and

$$v_i^2(x, t) = \int_0^t \int_{\mathbb{R}^3} K_{ijl}(x-y, t-\tau) \chi_{lj}(y, \tau) dy d\tau$$

$$\chi = V \otimes U$$

$$K_{i|e}(x, t) = \frac{\partial}{\partial x_j} K_{i|e}(x, t), \quad \leftarrow \text{ (Ozeen)}$$

where

$$K_{i|e} = \Phi_{g|e} - \delta_{ie} \Delta \bar{\Phi}$$

and

$$\Delta \bar{\Phi} = \square$$

Elementary estimates

$$|K(x, t)| \leq \frac{c_1}{(|x|^2 + t)^2}$$

and

$$\iint_{\mathbb{R}^3} |K(x, t)| dx dr \leq \frac{c_*}{\sqrt{t}}$$

Recall

Theorem. $\forall m = 0, 1, 2, \dots$

$$\|v(\cdot, t)\|_1 \leq c(m, c_d, v_0) \frac{\sqrt{t}}{\ln^m(t+\ell)^{3/2}}$$

and

$$\|v(\cdot, t)\|_2 \leq c(m, c_d, v_0) \frac{1}{\ln^m(t+\ell)}$$

$$\forall t \geq 1$$

Remark

Weak Leray - Hopf solution

$$\begin{cases} \partial_t v + v \cdot \nabla v - \Delta v = -\nabla q \\ \operatorname{div} v = 0 \end{cases} \quad \text{in } \mathbb{R}^3 \times [0, \infty[$$

$$v|_{t=0} = u_0 \in C^\infty_{0,0}(\mathbb{R}^3)$$

$$\|v(\cdot, t)\|_{L_2} \leq \frac{c}{\sqrt{t}} \quad t \geq 1$$

(Solv?)

Proof of Theorem A

1° Estimate of $\|v(\cdot, t)\|_{L_1}$ via $\|v(\cdot, t)\|_{L_2}$

Boundedness of $\|v(\cdot, \epsilon)\|_{L_2} \Rightarrow \|v(\cdot, t)\|_{L_1} \leq c(\sqrt{t})^{3/2}$

2° Modification of Schonbek and Wigner
technique \Rightarrow

$$\|v(\cdot, t)\|_{L_2} \leq \frac{c}{\ln t} \quad t > 1$$

3° Iteration

Little bit about Step 1.

$$(i) \quad \frac{6}{5} < p \leq \frac{3}{2}$$

$$A_p(t) := \int_0^t \frac{ds}{\sqrt{t-s}} \sqrt{s}^{-\frac{5p-6}{2p}} c_d \|v^2(\cdot, s)\|_2^d$$

$\| \cdot \|_p$ via $\| \cdot \|_2$

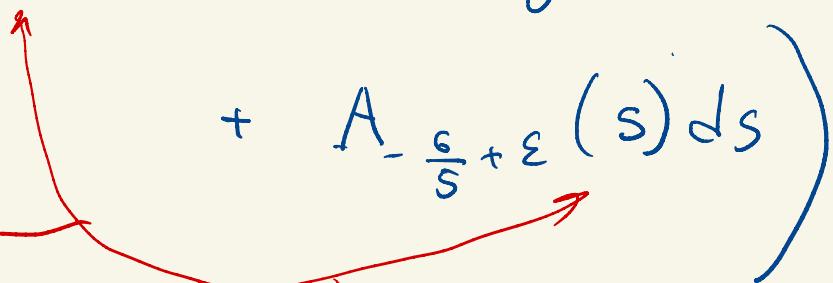
$$\|v^2(\cdot, +)\|_p \leq C(p) \left(c(c_d, v_0, p) \int_0^{\frac{3(1-p)}{p}} (t) + A_p(t) \right)$$

$$f(t) = \max \{ 1, \sqrt{t} \}$$

$$0 < \varepsilon < \frac{3}{10}$$

(ii)

$$\|v^2(\cdot, t)\|_1 \leq C(\varepsilon, c_d) \int_0^t \frac{ds}{\sqrt{t-s}} \sqrt{s}^{-\frac{-3+10\varepsilon}{6+5\varepsilon}} \left(C(c_d, v_0, \varepsilon) \sqrt{s}^{-\frac{3(1+5\varepsilon)}{6+5\varepsilon}} + \right.$$

$$\left. + A_{-\frac{6}{5}+\varepsilon}(s) ds \right)$$


$\|v\|_1$ since $\|v\|_{\frac{6}{5}+\varepsilon}$

$$\Rightarrow \|v^2(\cdot, t)\|_1 \leq C \sqrt{t}^{\frac{3}{2}}$$

Energy estimate Comments to Step 2

$$\partial_t y + \frac{1}{2} \|\nabla v(\cdot, \tilde{z})\|_{L_2}^2 d\tilde{z} = 0$$

where

$$y = \|\nabla(\cdot, t)\|_{L_2}^2$$

Fourier Transform

$$\partial_t y = - \int_{\mathbb{R}^3} |\xi|^2 |\hat{v}(\xi, t)|^2 d\xi = - \int_{|\xi| \leq g(t)} - \int_{|\xi| > g(t)}$$

$$\leq -g^2(t) \int_{\mathbb{R}^3} |\hat{v}(\xi, t)|^2 d\xi + \underbrace{\int (g^2(t) - |\xi|^2) |\hat{v}(\xi, t)|^2 d\xi}_{|\xi| \leq g(t)} + \underbrace{\int (g^2(t) - |\xi|^2) |\hat{v}(\xi, t)|^2 d\xi}_{|\xi| > g(t)}$$

From the equation

$$\partial_t \hat{v} + |\beta|^2 \hat{v} = -\hat{H},$$

where

$$H = -\operatorname{div}(v \otimes u) - \nabla q = -\operatorname{div}(v \otimes u + q \mathbb{I})$$



$$\hat{v}(\beta, t) = - \int_0^t e^{-|\beta|^2(t-\tau)} \hat{H}(\beta, \tau) d\tau + e^{-|\beta|^2 t} \left| \hat{v}_0(\beta) \right|$$

$$\Rightarrow |\hat{H}(\beta, \tau)| \leq c |\beta| \|v(\cdot, \tau) \otimes u(\cdot, \tau)\|_1$$

$$\partial_t y + g^2(t) y = c g^5(t) \left[g(t) \mathbb{E} \int_0^t a^2(s) ds + c(v_0) \right] := K(s)$$

$$a(s) = \| |u(\cdot, s)| |v(\cdot, s)| \|_1$$

$$y(t) \leq c \int_0^t e^{- \int_s^t g^2(z) dz} K(s) ds + y_0 e^{- \int_0^t g^2(z) dz}$$

$$g^2(t) = \frac{h'(t)}{h(t)}, \quad h(t) = \ln^k (t+s) \stackrel{m=0}{\Rightarrow} \| v^2(\cdot, t) \|_2 \leq \frac{c}{\ln(t+\ell)}$$

$k > 2m+2$

Upper estimate for c_d
(1st one)

$$\partial_t v - u \cdot \nabla v - \nabla \cdot v - \nabla q = -\operatorname{div} F, \quad \operatorname{div} v = 0$$
$$v|_{t=0} = 0$$

Theorem. Assume that

$$4c_* c_d \leq 1,$$

then

$$\int_{\mathbb{R}^3} u(x, T) \cdot v(x, T) dx \rightarrow 0 \quad T \rightarrow \infty$$

$$\boxed{\int_{\mathbb{R}^3} |K(x, y)| dy \leq \frac{C_*}{T}}$$

$$f(t) \leq c_* \int_0^t \frac{1}{\sqrt{t-s}} \left(\frac{c_d}{\sqrt{s}} f(s) + \|F(\cdot, s)\|_1 \right) ds, \quad (5.1)$$

where $f(t) := \|v(\cdot, t)\|_1$. Since F is compactly supported, (5.1) can be reduced to the following form:

$$f(t) \leq A + c_* c_d \int_0^t \frac{1}{\sqrt{t-s}\sqrt{s}} f(s) ds.$$

Now, fix an arbitrary $T > 0$. Then, for any $t \in]0, T]$, we have

$$f(t) \leq A + 4c_* c_d M(T),$$

where $M(T) = \sup_{0 < t \leq T} f(t)$. Hence,

$$M(T) \leq A + 4c_* c_d M(T)$$

for any $T > 0$. Finally, we see that

$$\|v(\cdot, t)\|_1 \leq c = \frac{A}{1 - 4c_* c_d}$$

for all $t > 0$. Therefore,

$$\left| \int_{\mathbb{R}^3} u(x, t) \cdot v(x, t) dx \right| \leq \frac{c}{\sqrt{t}} \rightarrow 0$$

as $t \rightarrow \infty$.

□

Second estimate

$$K : \mathcal{L}_2 \rightarrow \mathcal{J}_2$$

$$\mathcal{L}_2 = \left\{ \mathbf{F} = (\mathbf{F}_{ij}) \in L_2(\mathbb{R}^3) : \operatorname{div} \operatorname{div} \mathbf{F} = 0 \right\}$$

$$\mathcal{J}_2 = \left\{ \mathbf{v} \in L_2(\mathbb{R}^3) : \operatorname{div} \mathbf{v} = 0 \right\}$$

$$\Rightarrow A_F = KF \Leftrightarrow \operatorname{rot} A_F = -\operatorname{div} F$$

$$\downarrow$$
$$-\Delta A_F = -\operatorname{rot} \operatorname{div} F$$

$$M : L_2(\mathbb{R}^3; M^{3 \times 3}) \rightarrow L_2(\mathbb{R}^3)$$

$$q_r = M F \Leftrightarrow \Delta q_r = - \operatorname{div} \operatorname{div} F$$

Theorem C If $c_d \leq \frac{\sqrt{3}}{2 \|K\| (1 + \sqrt{3} \|M\|)}$



$$\lim_{T \rightarrow \infty}$$

$$\int_{\mathbb{R}^3} v(x, T) \cdot u(x, T) \, dx = 0.$$

The last theorem (Schrödinger's idea)

Translation from semigroup language

In the proof of the theorem we assume that A is skew-symmetric and therefore satisfies condition (1.3).

Equation (1.1) can be written as follows:

$$\partial_t v - \Delta v = \operatorname{div} F_0, \quad (5.2)$$

where

$$F_0 = v \otimes u + \nabla q \mathbb{I} - F.$$

We know from previous results that

$$F_0 \in L_{2,\infty}(Q_+), \quad \operatorname{div} F_0 \in L_2(Q_+). \quad (5.3)$$

Since $\operatorname{div} \operatorname{div} F_0 = 0$, we can apply the elliptic theory and conclude that there exists a divergence free field $A(\cdot, t)$ such that

$$\operatorname{rot} A(\cdot, t) = \operatorname{div} F_0(\cdot, t) \quad (5.4)$$

in \mathbb{R}^3 and the following estimate holds

$$\|A(\cdot, t)\|_2 \leq \|K\| \|F_0(\cdot, t)\|_2 \quad (5.5)$$

for all $t \in]0, \infty[$.

Taking into account the definition of the operator M , one can go further and derive from (5.5)

$$\begin{aligned}\|A(\cdot, t)\|_2 &\leq \|K\|(\|v(\cdot, t) \otimes u(\cdot, t)\|_2 + \|q_{v \otimes u}(\cdot, t)\mathbb{I}\|_2 + \|F(\cdot, t)\|_2) \leq \\ &\leq \|K\|(1 + \sqrt{3}\|M\|)\|v(\cdot, t) \otimes u(\cdot, t)\|_2 + h(t),\end{aligned}$$

where $h(t) = \|K\|\|F(\cdot, t)\|_2$ and thus

$$\|A(\cdot, t)\|_2 \leq \|K\|(1 + \sqrt{3}\|M\|)\frac{c_d}{\sqrt{t}}\|v(\cdot, t)\|_2 + h(t) \quad (5.6)$$

With the above A , let us consider the Cauchy problem

$$\partial_t B - \Delta B = A \quad (5.7)$$

$$B(\cdot, 0) = 0. \quad (5.8)$$

Problem (5.7), (5.8) has a unique solution defined for all positive t and $B \in W_2^{2,1}(Q_T)$ for all $T > 0$. Since $A(\cdot, t)$ is divergence free, so is $B(\cdot, t)$. Now, let $w = \text{rot } B$. Then we can see that w is a solution to equation (5.2) and since it vanishes at $t = 0$, we can state that $w = v$.

$$\Rightarrow v = \text{rot } B$$

Now, let us analyse the Cauchy problem for B . It is easy to see that B satisfies the energy identity

$$\frac{1}{2}\partial_t\|B(\cdot, t)\|_2^2 + \|\nabla B(\cdot, t)\|_2^2 = \int_{\mathbb{R}^3} A(x, t) \cdot B(x, t) dx. \quad (5.9)$$

Taking into account the simple identity

$$\|v(\cdot, t)\|_2 = \|\nabla B(\cdot, t)\|_2,$$

one can derive from (5.6) the following estimate

$$\begin{aligned} \frac{1}{2}\partial_t\|B(\cdot, t)\|_2^2 + \|v(\cdot, t)\|_2^2 &\leq \|K\|(1 + \sqrt{3}\|M\|) \frac{c_d}{\sqrt{t}} \|v(\cdot, t)\|_2 \|B(\cdot, t)\|_2 + \\ &+ h(t) \|B(\cdot, t)\|_2. \end{aligned}$$

Applying the Young inequality, we find

$$\frac{1}{2}\partial_t\|B(\cdot, t)\|_2^2 \leq \|K\|^2(1 + \sqrt{3}\|M\|)^2 \frac{c_d^2}{4t} \|B(\cdot, t)\|_2^2 + \frac{1}{2}h(t)(\|B(\cdot, t)\|_2^2 + 1)$$

Let us introduce the important constant

$$l = \|K\|^2(1 + \sqrt{3}\|M\|)^2 \frac{c_d^2}{2}.$$

Then the previous inequality leads to

$$\|B(\cdot, t)\|_2^2 \leq t^l \int_0^t \frac{h(\tau)}{\tau^l} \exp\left(-\int_\tau^t h(s)ds\right) d\tau.$$

$$h(s) = \|K\| \|F(\cdot, t)\|_2$$

Taking into account that F is compactly supported in Q_+ , we have

$$\|B(\cdot, t)\|_2^2 \leq c_F t^l.$$

From here, it is easy to derive the following:

$$\int_0^t \|v(\cdot, s)\|_2^2 ds \leq c_F t^l. \quad (5.10)$$

We denote all the constant depending of F and its support by c_F .

Having estimate (5.10) in mind, let us go back to equation (5.2) multiplying it by tv and integrating result over $\mathbb{R}^3 \times]0, t[$, as a result, we find the following differential inequality

$$\begin{aligned} \frac{1}{2}t\|v(\cdot, t)\|_2^2 + \int_0^t \|\nabla v(\cdot, s)\|_2^2 ds &= \frac{1}{2}\|v(\cdot, t)\|_2^2 + \int_0^t \int_{\mathbb{R}^3} sF(x, s) \cdot v(x, s) ds \leq \\ &\leq c_F \left(\int_0^t \|v(\cdot, s)\|_2^2 ds + 1 \right). \end{aligned}$$

The latter, together with boundedness of $\|v(\cdot, t)\|_2$, implies the bound

$$\|v(\cdot, t)\|_2^2 \leq c_F(t+1)^{l-1},$$

(Energy Inequality (5.9)
and
 $\|v(\cdot, t)\|_2 =$
 $= \|\nabla B(\cdot, t)\|_2$)

which, in turn, allows to improve the decay of $\|v(\cdot, t)\|_1$. To this end, we are going back to (2.4) and (2.5). Indeed, by the assumption of the theorem $l < 3/4$,

$$\begin{aligned} A_p(t) &\leq c \int_0^t \frac{1}{\sqrt{t-s}} s^{-\frac{5p-6}{4p}} (s+1)^{l-1} ds \leq c \int_0^t \frac{1}{\sqrt{t-s}} s^{-\frac{5p-6}{4p} + l - 1} ds \leq \\ &\leq ct^{\frac{6-3p}{4p} + l - 1}. \end{aligned}$$

Letting $p = 6/5 + \varepsilon$, for sufficiently small positive ε , we find

$$\|v(\cdot, t)\|_1 \leq c(\sqrt{t})^{\frac{3}{2} + 2(l-1)}.$$

This shows

$$\left| \int_{\mathbb{R}^3} v(\cdot, t) \cdot u(\cdot, t) dx \right| \leq c(\sqrt{t})^{\frac{1}{2} + 2(l-1)} \rightarrow 0$$

as $t \rightarrow \infty$ provided $l < \frac{3}{4}$. \square